

## RESEARCH ARTICLE

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## Key Points:

- Temporal variations of Jupiter's northern aurora were detected using Hisaki/EXCEED
- Auroral power enhancement are accompanied by a slight far ultraviolet color ratio (CR) increase
- CR-brightness distribution and longitude dependence are consistent with other observations

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## Variation of Jupiter's aurora observed by Hisaki/EXCEED: 1. Observed characteristics of the auroral electron energies compared with observations performed using HST/STIS

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**Abstract** Temporal variation of Jupiter's northern aurora is detected using the Extreme Ultraviolet Spectroscope for Exospheric Dynamics (EXCEED) on board JAXA's Earth-orbiting planetary space telescope Hisaki. The wavelength coverage of EXCEED includes the H<sub>2</sub> Lyman and Werner bands at 80–148 nm from the entire northern polar region. The prominent periodic modulation of the observed emission corresponds to the rotation of Jupiter's main auroral oval through the aperture, with additional superposed ~50%–100% temporal variations. The hydrocarbon color ratio (CR) adopted for the wavelength range of EXCEED is defined as the ratio of the emission intensity in the long wavelength range of 138.5–144.8 nm to that in the short wavelength range of 126.3–130 nm. This CR varies with the planetary rotation phase. Short- (within one planetary rotation) and long-term (> one planetary rotation) enhancements of the auroral power are observed in both wavelength ranges and result in a small CR variation. The occurrence timing of the auroral power enhancement does not clearly depend on the central meridian longitude. Despite the limitations of the wavelength coverage and the large field of view of the observation, the auroral spectra and CR-brightness distribution measured using EXCEED are consistent with other observations.

### 1. Introduction

Aurorae represent the environment and dynamics of a coupled magnetosphere-ionosphere-thermosphere system. Jupiter's auroral emission is often categorized into three regions: low-latitude moon-footprint emission, main aurora emission, and high-latitude polar emission (see the reviews of *Clarke et al.* [2004], *Badman et al.* [2014], *Grodent* [2014], and references therein). The moon-footprint aurorae are caused by the relative motion of electrically conductive moons and plasma carried by the surrounding planetary magnetic field. The main aurora emission is associated with the plasma corotation enforcement current during the transport of the angular momentum from the planetary neutral atmosphere through the ionosphere to the magnetosphere [e.g., *Hill*, 2001; *Cowley and Bunce*, 2001; *Cowley et al.*, 2007]. Jupiter's polar region, enclosed by the main aurora, corresponds to both open and closed magnetic field lines [e.g., *Vogt et al.*, 2011]. Several auroral features in the polar region have been related to magnetospheric reconnection events [e.g., *Grodent et al.*, 2004], emissions at the open-closed field line boundary [*Pallier and Prangé*, 2004], and short-term bursts at the dayside cusp [e.g., *Waite et al.*, 2001].

The ultraviolet (UV) emissions of Jupiter's aurora are radiated from atmospheric molecular (H<sub>2</sub>) and atomic hydrogen (H) excited by precipitating electrons. Jupiter's UV emission spectra show significant absorption by hydrocarbons [e.g., *Yung et al.*, 1982]. This absorption effect is measured as the color ratio (CR), which is defined as the ratio of the intensity of the wavelength bands unabsorbed by hydrocarbons to that of the absorbed wavelength bands. The CR represents the column density of hydrocarbons above auroral emissions because the hydrocarbons are located in the deep atmosphere. Assuming it to be related to the penetration depth of auroral electrons, the CR is used to estimate the electron energy, while the altitude profile of the hydrocarbons also modifies the CR [e.g., *Livengood and Moos*, 1990; *Livengood et al.*, 1993; *Gérard et al.*, 2003]. The auroral

spectra observed using the Hubble Space Telescope (HST) Space Telescope Imaging Spectrograph (STIS) revealed that the CR varies between spatial structures [Gustin *et al.*, 2002, 2004; Gérard *et al.*, 2014] and shows short-term variations over a few tens of seconds [Gérard *et al.*, 2003]. The high-latitude emissions sometimes show electron energies and energy fluxes similar to those of the main oval emissions; however, high electron energies (large CR) with low fluxes are also present [Gustin *et al.*, 2004]. The CR related to the brightness of H Lyman  $\alpha$  and H<sub>2</sub> has also been applied [e.g., Harris *et al.*, 1996], which is suggested to be sensitive to lower energy electrons because the CR refers to hydrogen atoms at a higher altitude [Tao *et al.*, 2014].

Continuous observations reveal several timescales of auroral variations. The auroral area visible to an observer needs to be accounted for when considering variations of the integrated auroral intensity [e.g., Prangé *et al.*, 2001; Pryor *et al.*, 2005] and far UV (FUV) CR [Livengood and Moos, 1990]. Prangé *et al.* [2001] reported a dominant intensity variation over 5–10 days associated with magnetospheric fluctuations in addition to small variations over a few hours and a longer (> 6 weeks) trend. A sporadic large intensity enhancement during one planetary rotation was detected by Cassini [Pryor *et al.*, 2005]. Auroral monitoring using the HST reveals some auroral intensity enhancements when a solar wind compression region foreshock was estimated to arrive at Jupiter, while no clear correlations with reverse shocks occurred [Clarke *et al.*, 2009]. Different auroral responses to solar wind variations were observed in sets of HST observations separated by several months [Nichols *et al.*, 2009]. The timescale of the auroral intensity variation in the polar region is as short as a few tens of seconds, which is considered to reflect the localized solar wind variation. The short-term auroral bursts at the dayside cusp are found when the solar wind dynamic pressure is enhanced [Waite *et al.*, 2001]. The periodic intensity variation with a timescale of 2–3 min is suggested to be related to the pulsed magnetic reconnections at the dayside by analogy with similar phenomena observed on Earth [Bonfond *et al.*, 2011].

Temporal variations and occurrence properties under various outer (i.e., solar wind) and inner (e.g., Io volcanic activity) conditions are crucial for understanding these auroral phenomena and related magnetospheric dynamics. Our new tool for monitoring the Jovian aurora is the Extreme Ultraviolet Spectroscope for Exospheric Dynamics (EXCEED) [Yoshioka *et al.*, 2013; Yoshikawa *et al.*, 2014; Yamazaki *et al.*, 2014] on board Japan Aerospace eXploration Agency's (JAXA's) Earth-orbiting planetary telescope Hisaki. EXCEED observed extreme UV (EUV) emission from Jupiter's northern polar region, which was our main interest here, and the Io plasma torus continuously over 40 min of every 106 min of the Hisaki orbit from December 2013 to April 2014. In addition, the HST observations were also carried out during the first half of January 2014. EXCEED could detect sporadic, large auroral intensity enhancement lasting less than one planetary rotation, which was associated with auroral low-latitude intensifications observed in the HST images [Kimura *et al.*, 2015]. In this study, we investigate the time variation of the CR with auroral brightness variations using the EXCEED. Because the CR-energy relationship is model dependent, this study refers to the CR instead of converting it to electron energy. The parameters obtained in this study are estimated using the auroral emission from the entire northern polar region. Spatially resolved spectra measured through the STIS on board the HST are referred to check the spatial variations, as described in section 2. EXCEED covers different spectral ranges compared to STIS, such that we define a new CR for the EXCEED spectral analysis. The observation details of EXCEED and analyzed results, including the relationship to modeled solar wind conditions, are described in section 3. Section 4 concludes this paper.

## 2. Spatial Variation of Aurora Detected Using HST/STIS

Following the HST spectral analysis by Gustin *et al.* [2004], we check the CR-brightness relation using the spectra measured by HST/STIS over the first 2 weeks of January 2014.

### 2.1. Imaging and Spectral Observations

In the HST observations program (ID: GO13035), the FUV-MAMA detector of the STIS obtained FUV images and spectra of Jupiter's northern aurora. The auroral images were obtained using a SrF<sub>2</sub> long-pass filter to detect H<sub>2</sub> emission in the wavelength range of 125–170 nm with a plate scale of  $\sim 0.0224$  arc sec pixel<sup>-1</sup>. We used geometric distortion-corrected imaging data in "x2d" files with a unit of *c/s*. The long slit with a size of  $52 \times 0.5$  arc sec<sup>2</sup> with G140L grating provides spatially resolved spectra at wavelengths of 110–170 nm with a resolution of  $\sim 1.2$  nm. We used the flux- and wavelength-calibrated two-dimensional spectra in x2d files with a unit of  $\text{erg s}^{-1} \text{cm}^{-2} \text{\AA}^{-1} \text{arc sec}^{-2}$ . On each HST orbit, the observations were made in the following

**Table 1.** Summary of the Imaging and Spectral Observations Conducted Using the STIS: Date, Start Time in UT, Corresponding CML, Brightness Over 155–162 nm, CR<sub>STIS</sub>, Brightness Over 138.5–144.8 nm, and CR<sub>EXCEED</sub> Derived From the STIS Spectra at the Main Aurora on the Disk, Emitted Power Over 138.5–144.8 nm, and CR<sub>EXCEED</sub> Derived From the EXCEED Observation at the Timing Closest to the STIS Spectral Observation of Each Event

#	Date	STIS Image Observation 1 (Exposure 700 s)		STIS Spectral Observation (Exposure 200 s)		STIS Image Observation 2 (Exposure 736 s)		STIS Spectra Values of the Main Aurora on the Disk			EXCEED Values Closest to the STIS Spectral Observation		
		Start Time (UT)	CML (deg)	Start Time (UT)	CML (deg)	Start Time (UT)	CML (deg)	Brightness (155–162 nm) <sup>a</sup>	CR <sub>STIS</sub>	Brightness (138.5–144.8 nm) <sup>a</sup>	CR <sub>EXCEED</sub>	Emitted Power (138.5–144.8 nm) (GW) <sup>c</sup>	CR <sub>EXCEED</sub>
1	Jan 1	03:02	173	03:23	186	03:32	191	2.4	2.0	0.60	1.1	24.	1.1
2	Jan 2	09:19	192	09:40	204	09:49	210	6.8	4.3	1.3	1.8	27.	1.6
3	Jan 3	04:27	166	04:47	178	04:57	184	9.9	2.8	2.0	1.2	19.	1.4
4	Jan 4	01:10	197	01:31	210	01:40	215	2.1	1.7	0.58	1.0	32.	1.2
5	Jan 5	05:52	158	06:12	171	06:21	176	7.9	3.2	1.5	1.3	26.	1.4
6	Jan 6	02:35	190	02:55	202	03:05	208	2.1	2.1	0.44	0.92	18.	1.1
7	Jan 7	07:16	151	07:37	163	07:46	169	9.3	3.6	1.5	1.3	24.	1.6
8	Jan 10	03:48	117	04:09	129	04:18	135	– <sup>b</sup>	–	–	–	22.	1.4
9	Jan 11	00:31	149	00:52	161	01:01	167	19.9	3.1	4.0	1.3	40.	1.4
10	Jan 11	19:39	122	20:00	134	20:09	140	–	–	–	–	22.	1.2
11	Jan 13	01:56	141	02:16	153	02:26	159	5.2	2.6	0.94	1.0	22.	1.5
12	Jan 13	21:03	115	21:24	127	21:33	133	– <sup>b</sup>	– <sup>b</sup>	– <sup>b</sup>	– <sup>b</sup>	19.	1.4
13	Jan 13	22:39	172	22:59	185	23:09	191	2.9	2.1	0.62	1.1	17.	1.3
14	Jan 16	00:03	165	00:24	177	00:33	183	2.8	2.5	0.57	1.1	11.	1.1

<sup>a</sup>Unit is 10<sup>-13</sup> mW m<sup>-2</sup> arcsec<sup>-2</sup>.

<sup>b</sup>The STIS slit only crosses the aurora at the limb.

<sup>c</sup>The auroral power after the background subtraction, without scaling for visibility of the auroral region.

sequence, image (700 s), spectrum (200 s), and image (736 s), using a time-tag mode, such that the exposure time could be divided into shorter integration times. We use the time-integrated spectra and images over each interval in this analysis. This sequence was repeated for 14 HST orbits spaced over 2 weeks. The intensity profiles from the filtered images were compared to the slit spectra across the full wavelength range by accounting for the filter throughput the function. The slit position in the north-south direction is determined by the position of the limb, and that in the east-west direction was determined by matching the intensity profile along the slit with the profiles from the images obtained before and after the spectral observations. The date, time, and system III central meridional longitude (CML) of the spectral and image observations used in this study are summarized in Table 1.

## 2.2. Results

Figure 1a shows the HST image and slit position of the spectral observation (white vertical line) obtained on the same HST orbit on 2 January 2014. The slit crosses from the top to the bottom, the main auroral oval at the limb, faintly enhanced high-latitude emission, and the main auroral oval on the disk. These three regions are seen in the spatial profile of the auroral brightness along the slit (Figure 1b, left), as shown by light blue, red, and blue lines, respectively. Figures 1c and 1d show the image and profile obtained on 7 January in the same format. The emission intensity of the main aurora on 7 January was lower except for localized enhancement around noon at the limb. The high-latitude emission is faint in the region under the slit (Figure 1d). The localized bright high-latitude emission is detected in the auroral image (Figure 1c). The typical FUV CR used for the STIS spectra is defined as the intensity ratio between the wavelength bands unabsorbed by the hydrocarbons  $I_{(155-162\text{ nm})}$  and the wavelength bands absorbed by the hydrocarbons  $I_{(123-130\text{ nm})}$ , i.e.,

$$CR_{\text{STIS}} = I_{(155-162\text{ nm})} / I_{(123-130\text{ nm})}, \quad (1)$$

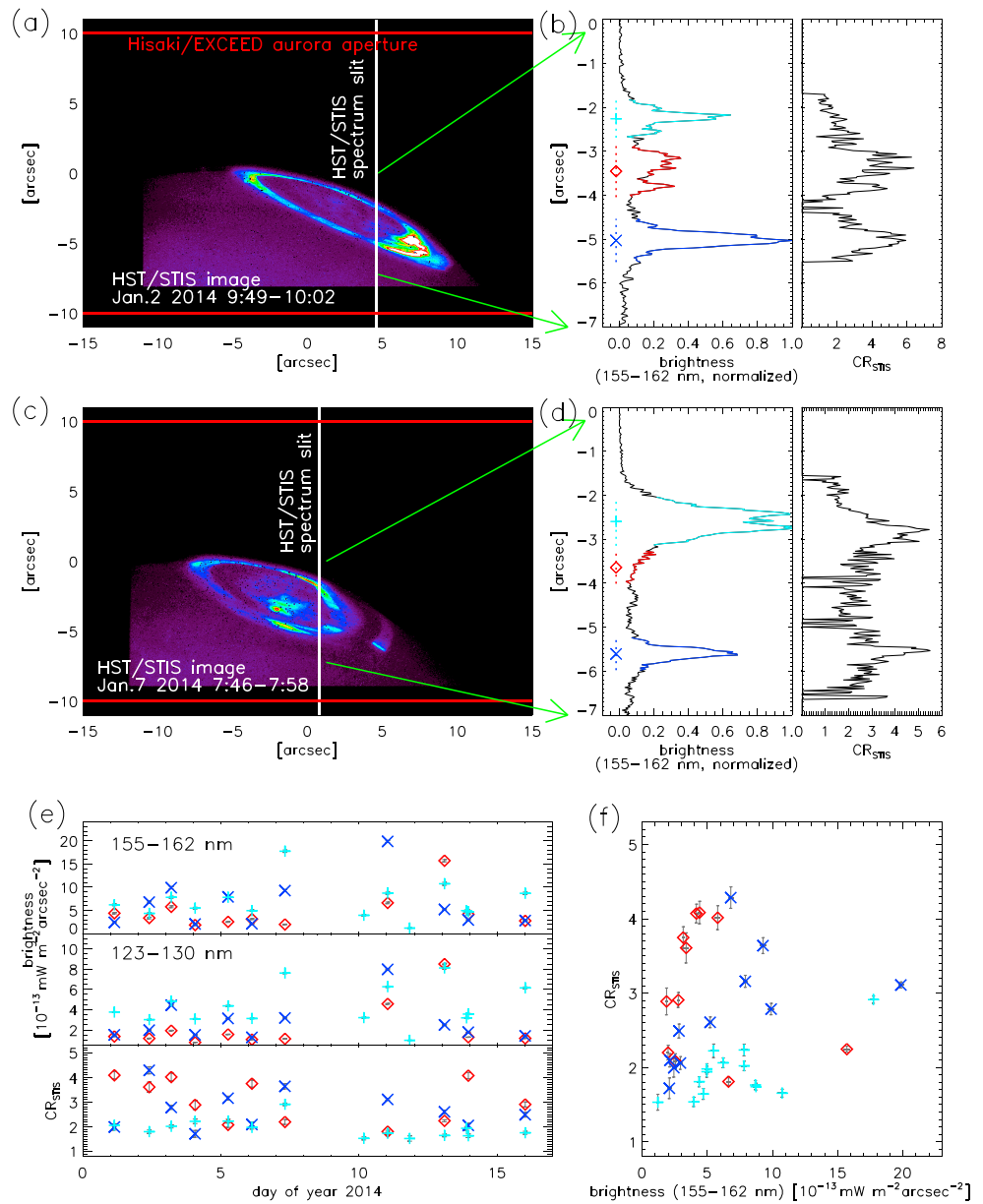
where  $I$  is the height-integrated intensity of the emission, e.g., in units of kilo-Rayleigh (kR) or photons  $\text{s}^{-1}$  integrated over the wavelength of the subscript. The absorption cross section of methane is shown using the blue line in Figure 2c. The STIS spectra are provided in energy flux units, while the CR is provided in photon flux units. A ratio of (155–162 nm)/(123–130 nm) obtained from the original STIS spectra is multiplied by 1.25 to account for the unit change. For simplicity, the viewing angle is not considered in our procedure; i.e.,  $CR_{\text{STIS}}$  is derived using the emission in each pixel. The  $CR_{\text{STIS}}$  values 2 and 8 are related to electron mean energies of 60 keV and 170 keV, respectively, for an eddy diffusion coefficient of  $1.4 \times 10^6 \text{ cm}^2 \text{ s}^{-1}$  at the homopause (see Gérard *et al.* [2003] for details). The  $CR_{\text{STIS}}$  profiles along the slit (Figures 1b (right) and 1d (right)) vary up to  $\sim 6$ , with larger values in the high-latitude region and the main oval on the disk. The time variations of the emitted power at these wavelengths and  $CR_{\text{STIS}}$  for the entire HST campaign are shown in Figure 1e and Table 1. The intensities are integrated over 200 s. The signal-to-noise ratio is high as indicated by the short error bars. The relationship between the unabsorbed brightness at wavelengths of 155–162 nm and  $CR_{\text{STIS}}$  varies between the spatial regions (Figure 1f).  $CR_{\text{STIS}}$  increases with the emitted power at the main auroral oval (light blue pluses and blue crosses).  $CR_{\text{STIS}}$  in the high-latitude region (red diamonds) has two components: one component is similar to the main aurora and the other component shows a lower intensity and higher  $CR_{\text{STIS}}$ .

This relationship is comparable with the relationship between the electron energy and electron energy flux derived by Gustin *et al.* [2004]. Their main auroral oval and high-latitude regions correspond to the main aurora at the disk and high-latitude emission here, respectively. The regional  $CR_{\text{STIS}}$ -brightness relation obtained here is consistent with their results from a different data set. The  $CR_{\text{STIS}}$  at the limb takes small values, as also observed by Gustin *et al.* [2002] and Gérard *et al.* [2014].

## 3. Temporal Variation of the Aurora Detected Using Hisaki/EXCEED

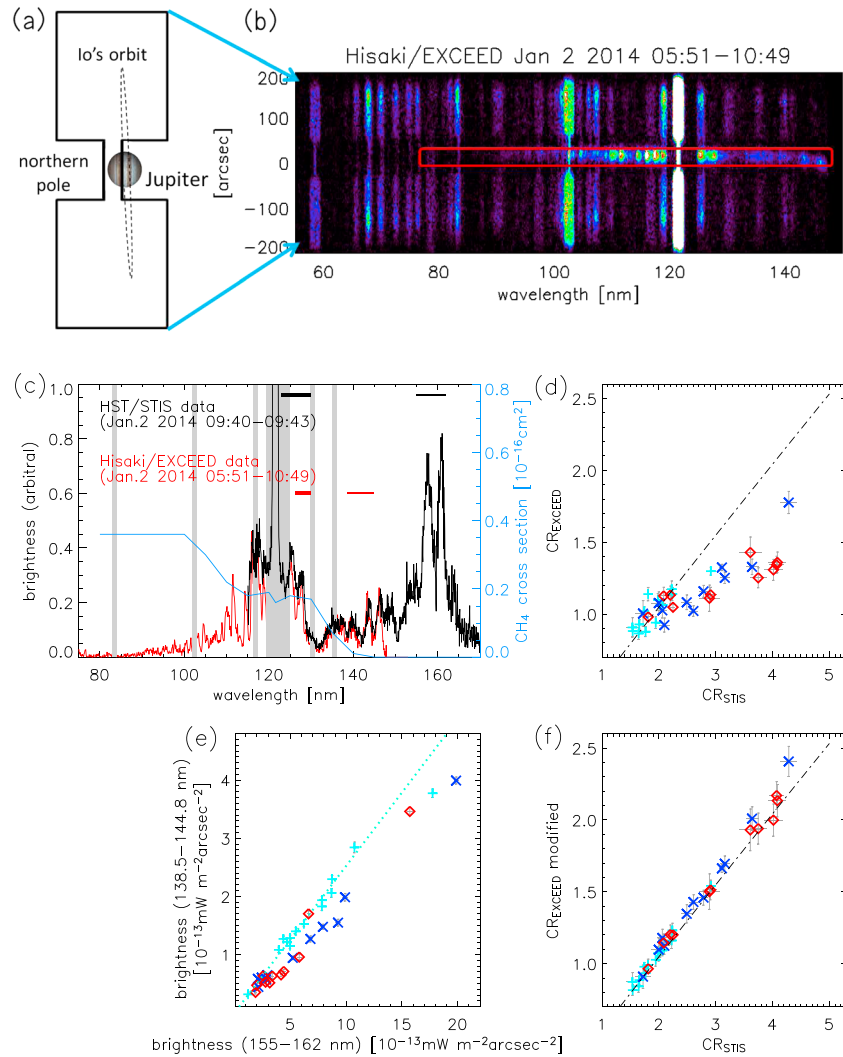
### 3.1. Observations

EXCEED counts EUV photons as a function of the position along the slit and wavelength dispersion. We used the data set obtained using the dumbbell-shaped slit, which detected emissions from the Io plasma torus and Jupiter's northern polar region simultaneously. The slit shapes and a sample imaging spectrum are shown in Figures 2a and 2b, respectively. The slit width in the polar region is 20 arc sec in the north-south direction (along Jupiter's rotation axis) with an effective spatial resolution along the slit (dawn-dusk direction) of 17 arc sec [Yoshikawa *et al.*, 2014] and a pointing accuracy of  $\pm 2$  arc sec. The red solid lines in Figure 1a show



**Figure 1.** (a) Image of Jupiter’s northern polar region, (b) the spatial profiles of the observed brightness integrated over wavelengths of (left) 155–162 nm and the (right) CR from the spectral observation along the slit on 2 January 2014, (c) the image and (d) the spatial profiles of the brightness and CR observed on 7 January 2014, (e) the time variations of the auroral brightness emitted at wavelengths (top) of 155–162 nm and (middle) of 123–130 nm and (bottom) the time variation of the CR, and (f) the relationship between the brightness at wavelengths 155–162 nm and the color ratio. The values in Figures 1e and 1f are taken at the main auroral oval on the disk (blue crosses), main auroral oval at the limb (light blue pluses), and high-latitude polar region (red diamonds) along the spectral slit, with error bars (grey lines). The corresponding auroral structures are shown using the same colors in Figures 1b (left) and 1d (left) for 2 and 7 January, respectively. Figures 1a and 1c are shown in the same linear color scale. The aurora aperture of EXCEED is bounded by the two red lines in Figures 1a and 1c, and white vertical lines show the positions of the STIS slit for the spectral observations.

the coverage of the EXCEED auroral aperture in the northern hemisphere. The data set excluded times when Jupiter was eclipsed by the Earth or Hisaki was located in the southern Atlantic anomaly. Under the latter, the instruments were turned off to avoid effects of energetic particles precipitated from the terrestrial radiation belts. Here we analyze the data obtained from 21 December 2013 to 31 January 2014 when the EXCEED time coverage was maximum. EXCEED detects auroral emission in the wavelength range of 80–148 nm, covering



**Figure 2.** (a) Schematics of the observation geometry of EXCEED using the dumbbell-shaped slit, (b) imaging spectrum measured by EXCEED, (c) auroral spectra measured on 2 January by STIS integrated over the slit (black) and by EXCEED over the full polar region (red), and the relationships (d) between the two color ratios, CR<sub>STIS</sub> and CR<sub>EXCEED</sub>, determined from the STIS spectra (e) between the brightness at the wavelength bands less affected by absorption selected for the STIS (155–162 nm) and EXCEED (138.5–144.8 nm) and (f) between CR<sub>STIS</sub> and modified CR<sub>EXCEED</sub>, estimated using STIS spectra. The values for the main auroral oval (blue crosses), main auroral oval at the limb (light blue pluses), and high-latitude region (red diamonds) with error bars (grey lines) were derived. In Figure 2c, the grey hatched regions correspond to the H Lyman, He, and O emission lines from geocorona. The absorption cross section for methane [Parkinson *et al.*, 2006] is indicated by the blue line corresponding to the right-hand axis, and the two wavelength bands used for the CR estimations are shown using the horizontal lines. The dash-dotted line in Figures 2d and 2f shows the reference relation between CR<sub>STIS</sub> and CR<sub>EXCEED</sub> based on the CH<sub>4</sub> absorption effect. The light blue dotted line in Figure 2e is the best fit to the brightness values obtained for the main aurora at the limb.

part of the H<sub>2</sub> Lyman (B → X) and Werner (C → X) band emissions with a full width at half maximum resolution of 0.3 nm. We convert the photon counts arriving at the EXCEED detector into the photon flux at each wavelength referring to the effective area [Yoshikawa *et al.*, 2014]. Then, the number flux is converted into energy flux and integrated over a half hemisphere assuming isotropic emission to derive the total emitted power. Figure 2c shows the spectra measured using STIS (black line) and EXCEED (red line) close in time. Except for the geocoronal emission range (e.g., around 121.6 nm Lyman α) and spectral edges, the spectra observed using EXCEED and STIS are consistent. The data are integrated over certain wavelength bands (described in section 3.2) and over 10 min to improve the signal-to-noise ratio, such that the variations in the northern



auroral activity over timescales from a few tens of minutes to a few months are detected. The long temporal coverage of the average auroral activity provided by the EXCEED data set is complementary to the short temporal (approximately seconds) and spatial variations observed with STIS [e.g., Gérard *et al.*, 2003; Gustin *et al.*, 2004; Gérard *et al.*, 2014] (section 2.2).

### 3.2. CR for EXCEED

EXCEED covers a wavelength range up to 148 nm, such that the  $CR_{STIS}$  given by equation (1) is not directly applicable. An alternate  $CR_{EXCEED}$  is defined as

$$CR_{EXCEED} = I_{(138.5-144.8 \text{ nm})} / I_{(126.3-130 \text{ nm})}. \quad (2)$$

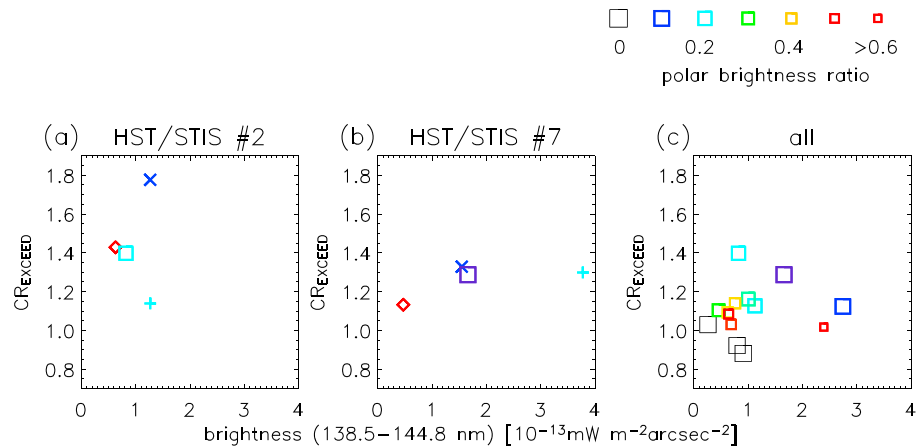
We selected these two wavelength ranges for  $CR_{EXCEED}$  using the following criteria: (i) the  $CH_4$  absorption cross section is significantly different in the two wavelength ranges; (ii)  $H_2$  self-absorption is not effective at these wavelengths, i.e.,  $>120$  nm [e.g., Gustin *et al.*, 2013]; and (iii) EXCEED has good sensitivity. A factor of  $(144.8 + 138.5) / (130 + 126.3) = 1.10$  is multiplied to the ratio of the intensities in power units for EXCEED, as for  $CR_{STIS}$  in section 2.2.

We used the 14 STIS spectral observations to determine the relationship between  $CR_{STIS}$  and  $CR_{EXCEED}$ . For each observation, a spectrum from the main oval (blue crosses), high-latitude (red diamonds), and limb auroral (light blue pluses) regions was obtained and used to derive the two CRs (Figure 2d). The expected ideal relation of  $CR_{STIS}$  and  $CR_{EXCEED}$  represented by the dash-dotted line is based on the  $CH_4$  absorption cross section (see Appendix A for details). The derived values follow the ideal relation in the small CR range ( $CR_{STIS} < 2.5$ ,  $CR_{EXCEED} < 1.2$ ) as expected; however, departures from the ideal relation increase at larger CR. Although the absorption by methane is dominant, acetylene ( $C_2H_2$ ) and ethane ( $C_2H_6$ ), which have a significant absorption cross section up to 145 nm, should also affect our estimations of the CR. The brightness derived from the less-absorbed wavelength bands selected for the STIS (155–162 nm) and EXCEED (138.5–144.8 nm) show two linear relations (Figure 2e). The brightness from the main aurora around the limb (light blue pluses) shows one linear relation with a slope of approximately 0.26 (light blue dotted line), and the brightness in the other regions (blue crosses and red diamonds) also show linearity with a shallower slope. The former would be largely contributed by high-altitude emission, which is less affected by the hydrocarbon absorption. For the disk emission subject to absorption,  $H_2$  emission at wavelengths of 138.5–144.8 nm is attenuated more than that at 155–162 nm by hydrocarbons. Because this attenuation decreases  $CR_{EXCEED}$  through its numerator,  $CR_{EXCEED}$  increases less rapidly than  $CR_{STIS}$ . When  $CR_{STIS}$  increases ( $>3$ ), the absorption at 138.5–144.8 nm becomes more important because the ethane optical depth becomes significant. We obtain the modified  $CR_{EXCEED}$  using the brightness in the wavelength range of 138.5–144.8 nm replaced by the expected linear relation (the dotted line in Figure 2e) referring to the unattenuated band of  $CR_{STIS}$ . In this modified estimation, the relation between  $CR_{STIS}$  and the modified  $CR_{EXCEED}$  closely follows the expected relation (shown in Figure 2f). Because EXCEED integrates over various regions with weak and strong hydrocarbon absorption, the application of this conversion factor is not simple, such that we do not apply this modification in this and a companion paper [Tao *et al.*, 2016]. This analysis also highlights the possible discrepancy in the brightness and  $CR_{EXCEED}$  caused by this wavelength band selection and the observing geometry. Owing to the monotonic relation of the CRs, the aforementioned modification does not qualitatively affect the main results of this as well as the companion papers.

Here we briefly estimate the  $CR_{EXCEED}$  variations comparable with the previous observations. The measurements of  $CR_{STIS}$  at the main oval vary between 1.5 and 5 and can reach  $\sim 10$  [Gustin *et al.*, 2002, 2004]. In the CR defined by Livengood and Moos [1990],  $CR = I_{(155.7-161.9 \text{ nm})} / I_{(123-130 \text{ nm})} \times (1300-1230) / (1619-1557)$ , wavelength ranges close to those used to define  $CR_{STIS}$  are used. They applied this result to earlier International Ultraviolet Explorer (IUE) observations to find CR variations mainly in the range of 1–5, sometimes increasing up to  $\sim 8$ . According to the ideal  $CR_{STIS}$ – $CR_{EXCEED}$  relation, a variation between the values of 0.7–2 with a maximum of  $\sim 5$  is expected.

### 3.3. Spatial Integration Effect of the EXCEED Observations

We check (1) the effect of spatial integration along the slit on the derived emitted powers and CR using the STIS spectral data set and (2) a hypothesis that a large CR and a small brightness derived from the EXCEED data can also represent a relative enhancement of the high-latitude emission. We derive the  $CR_{EXCEED}$



**Figure 3.** Relationships between the auroral emitted power at wavelengths of 155–162 nm and the CR using the similar format as Figure 1f for observations (a) on 2 January 2014 and (b) on 7 January 2014. The squares indicate the values estimated from the observation integrated over the entire auroral region along the slit in each observation, and those for (c) all events are shown. The color and size of the squares represent the ratio of the brightness of the high-latitude region to the total auroral brightness integrated along the slit according to the legend at the top.

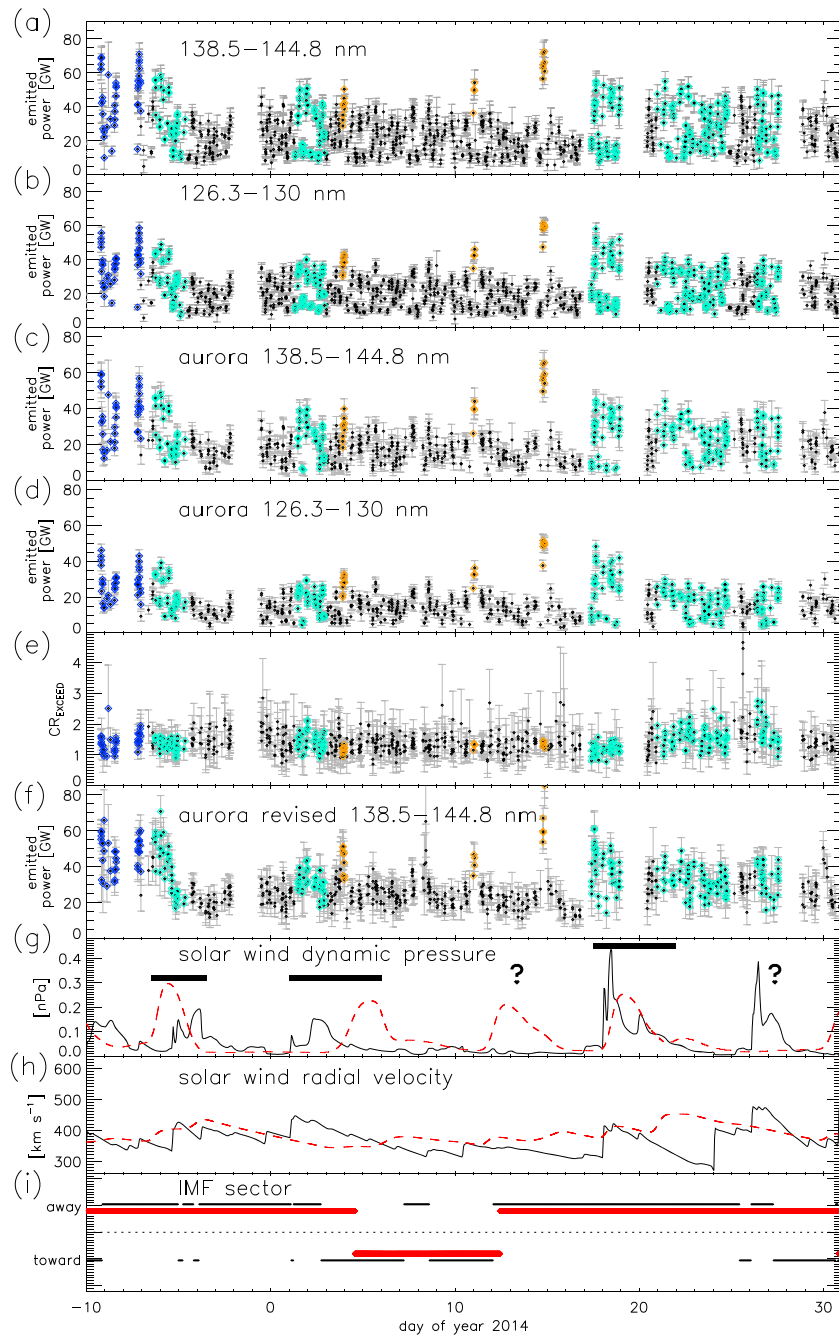
brightness relation from the auroral brightness spatially averaged over the entire auroral region under the STIS slit, shown by the squares in Figure 3. The relation derived from the observations performed on 2 January (event #2) are located almost centrally between the CR<sub>EXCEED</sub>-brightness relations of different auroral components, i.e., the main aurora (blue cross), high-latitude emission (red diamond), and main aurora at the limb (light blue plus), as shown in Figure 3a. The ratio of the high-latitude brightness to the brightness integrated along the slit is 0.24 for this event. A similar distribution is seen from the observation on 7 January (event #7), with a smaller high-latitude-to-total brightness ratio of 0.14 owing to the faint high-latitude feature captured under the slit (Figure 3b). Therefore, the CR<sub>EXCEED</sub>-brightness relations derived from the auroral brightness integrated over the full auroral region are intermediate between the values determined for different spatial regions. This can be representative of the EXCEED field of view over the entire northern auroral region as described in the following sections. Figure 3c shows the spatially integrated CR<sub>EXCEED</sub>-brightness relations, color coded according to the polar brightness ratio, for all cases. Relatively large polar brightness events (>0.4) are seen with small average brightness, (0.4–1.2) × 10<sup>-13</sup> mW m<sup>-2</sup> arc sec<sup>-2</sup>, and CR<sub>EXCEED</sub> ~1–1.2. The only exceptional event with a high brightness of ~2.4 × 10<sup>-12</sup> mW m<sup>-2</sup> arc sec<sup>-2</sup> and CR<sub>EXCEED</sub> of ~1 is observed on 13 January. In addition, events with a smaller polar brightness ratio of ~0.2 are also seen at similar brightness-CR<sub>EXCEED</sub> values. Therefore, this data set cannot establish the spatially resolved analysis.

### 3.4. Temporal Variations of the EXCEED Observations

Figure 4 shows the time variations of the auroral emitted power observed with EXCEED from 21 December 2013 (day of year 2014, DOY -10) to 31 January 2014 (DOY 31). The power emitted at wavelengths of 138.5–144.8 nm (Figure 4a) and 126.3–130 nm (Figure 4b) varies over several timescales. Over the timescale of one planetary rotation, the emitted power at 138.5–144.8 nm changes from ~10 to ~40 GW. The upper limit of the observed emitted power varies more than the lower limit, i.e., ~20 GW at DOY 15, ~40 GW at DOY 0–10, and ~60 GW at DOY ~20 and DOY -7. In addition, two different types of enhancement can be seen. The first type is short-term, occurring within one planetary rotation, on DOYs 4, 11, and 14, as reported by Kimura *et al.* [2015] (the periods shown in orange color in Figures 4a–4f). The other type of enhancement is long term (lasting several planetary rotations, the periods shown in light blue color) at DOY -10 to -6, DOY 1–2, and DOY 17–27. The power emitted at 126.3–130 nm also varies similarly.

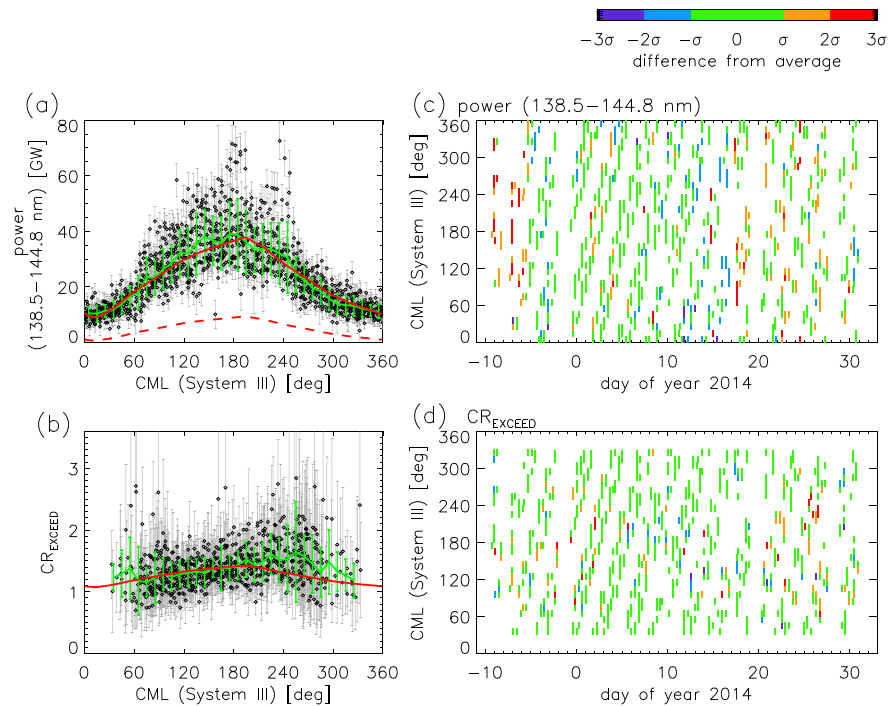
The power emitted at wavelengths of 138.5–144.8 nm has a clear CML dependence, as shown in Figure 5a. The dashed red line shows the length of the dayside auroral oval on the visible disk as a function of the CML. This length is obtained from the region in the northern ionosphere mapping to an *L* value of 30 using the VIP4 magnetic field model plus the ring current contribution [Connerney *et al.*, 1998, Table 4]. The red solid line shows this scaled profile with a constant added to match the magnitude of the observed emitted





**Figure 4.** Time variations of the total power emitted at wavelengths (a) of 138.5–144.8 nm and (b) 126.3–130 nm; (c, d) those from aurora (with background subtraction, see the text); (e) CR, observed with EXCEED (black points) with error bars representing photon statistics errors (grey lines); (f) the emitted power scaled for visibility of the auroral region; and time variations of the solar wind (g) dynamic pressure, (h) radial velocity, and (i) the IMF sector, estimated using a 1-D model (black lines) and SUSANOO (thick or dashed red lines). The periods of short-term and long-term auroral emitted power enhancements are shown in orange and light-blue colors, respectively, except for the intermittent observation before DOY 7 as shown in blue. The solar wind dynamic pressure enhancements observed in both the 1-D and SUSANOO models are shown using the horizontal thick bars in Figure 4g.

power profile. The average of the emitted power in 36 CML bins is shown using the green lines. Their profile in CML is comparable with that of the scaled visible auroral area (red solid line). The EXCEED observation includes emission from the entire region including those inside and outside the auroral oval, which adds variations in the CML dependence. According to the observation of Jupiter’s auroral emission performed using



**Figure 5.** (a) Power emitted at wavelengths of 138.5–144.8 nm and (b)  $CR_{EXCEED}$  shown as a function of the system III CML over DOYs from –11 to 31 2014, and the color maps of (c) the power emitted at wavelengths of 138.5–144.8 nm and (d)  $CR_{EXCEED}$  normalized by standard deviations  $\sigma$  in 36 CML bins (as in color bars), as functions of the time and the CML. In Figures 5a and 5b, the black dots are observed points and the grey lines show their errors, and the average values in 36 CML bins are shown using the green lines with error bars representing the standard deviation. The red solid lines in Figures 5a and 5b are the fits to the average of each parameter using the relative northern auroral area shown by the red dashed line in Figure 5a.

the Cassini/Ultraviolet Imaging Spectrograph at 111.5–191.2 nm, the background disk emission can be small compared to the auroral emission in the wavelength range covered by EXCEED below 148 nm [Pryor *et al.*, 2005]. Airglow emission and aurora at the limb can affect the derived emitted power in addition to the background disk emissions, as seen by the nonzero emitted power detected at a CML of  $\sim 0^\circ$  when the aurorae are on the antiobserver/nightside of the planet. In order to minimize these effects, the auroral emitted powers are derived by subtracting a 5 day running average of the emitted powers measured when  $0^\circ < CML < 30^\circ$ . The subtracted power (7–15 GW) is smaller than the auroral dynamic variation (15–80 GW).

The revised auroral emitted power at wavelengths of 138.5–144.8 nm is shown in Figure 4c, and that at 126.3–130 nm is shown in Figure 4d.  $CR_{EXCEED}$  is derived as the ratio of the background-subtracted emitted powers, and large signal-to-noise ( $>1.5$ ) data are used. Figure 4e shows the observed  $CR_{EXCEED}$  during this period is in the range of 0.8–2 and sometimes enhances to  $\sim 4.5$ , which is comparable with the previous observations described in section 3.2. The values of  $CR_{EXCEED}$  during the short- or long-term auroral enhancements are similar to those before and after the events, except for the  $CR_{EXCEED}$  enhancement after DOY 20. Figure 4f shows the auroral emitted power at wavelengths of 138.5–144.8 nm scaled for visibility of the auroral region by multiplying the factor (maximum visible auroral length at all CMLs)/(visible auroral length at instantaneous CML). A short-term enhancement on DOY 8 becomes visible in addition to the three other short-term events. Long-term variations are similar to those described above with modified amplitudes.

### 3.5. Solar Wind Model and Comparison With Aurora

We compare the auroral variation with solar wind parameters predicted using different models. A one-dimensional (1-D) magnetohydrodynamic (MHD) model propagates the observed solar wind conditions around Earth to Jupiter [Tao *et al.*, 2005]. For the model input, we use OMNI 1 h data, which are calibrated solar wind archive based on solar wind observations around the Earth (e.g., [http://omniweb.gsfc.nasa.gov/html/ow\\_data.html](http://omniweb.gsfc.nasa.gov/html/ow_data.html)). During

the observations from 21 December 2013 to 31 January 2014 of interest here, Jupiter was located at opposition on 6 January and the Earth-Sun-Jupiter angle was small enough ( $<30^\circ$ ) to estimate the arrival time of solar wind pressure enhancement with a good accuracy of  $\sim 1$  day [Tao *et al.*, 2005]. The weakness of the 1-D model is the treatment of the longitudinal variation, e.g., localized disturbance related to coronal mass ejections, which pass either Jupiter or Earth. We confirm that the longitudinal-limited structures are not probable during the investigated term before mid-January referring to 3-D models of ENLIL [e.g., *Odstrcil and Pizzo*, 1999] and SUSANOO [Shiota *et al.*, 2014]. We also refer to a 3-D model. SUSANOO solves 3-D MHD propagation of the solar wind parameters from the vicinity of the Sun based on synoptic maps of the photospheric magnetic field provided by the Global Oscillation Network Group (GONG) and empirical models. Good accuracy in predicting the interplanetary magnetic field (IMF) sectors using the 3-D model is evaluated at the positions of planets (see details in Shiota *et al.* [2014]). A difference in the predicted arrival time of solar wind at Jupiter of  $\pm 1$  day can be brought by an ambiguity of  $\pm 20 \text{ km s}^{-1}$  if the propagation velocity is  $400 \text{ km s}^{-1}$ , while 3-D MHD models also provide at least this ambiguity so far.

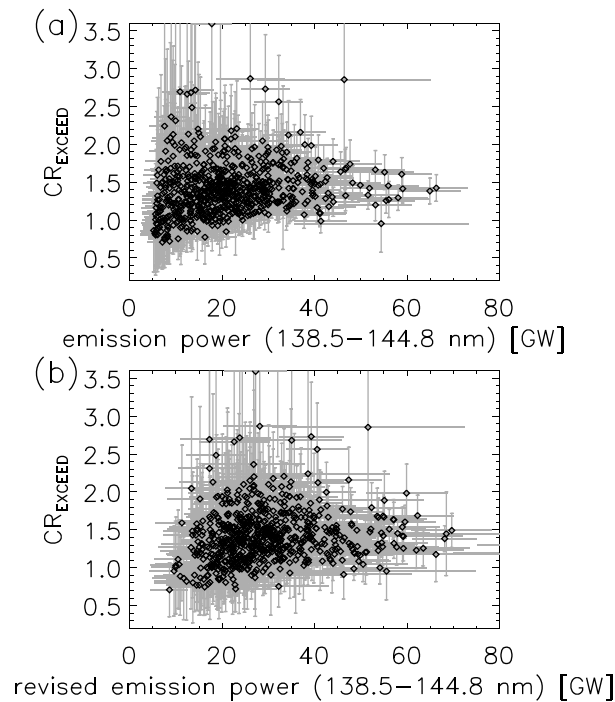
Large enhancements in dynamic pressure ( $>0.1 \text{ nPa}$ ) lasting a few days on DOYs  $\sim -6, 1-4$ , and  $17-21$  are predicted by both the 1-D and 3-D SUSANOO models with less than 2 day difference in the arrival time between the two models (Figure 4g). A pressure enhancement on DOYs  $12-14$  is only predicted with the 3-D model, while that on DOY  $26-28$  is only found using the 1-D model. The enhanced pressure events obtained in both models are close to the auroral brightness enhancement lasting longer than one planetary rotation (light blue color in Figure 4). The short-term (less than one planetary rotation) enhancements of the auroral brightness have a shorter duration than the solar wind pressure variation. The auroral response to the solar wind pressure enhancement on DOY  $26-28$  is not clear, which would be partly due to the continued auroral enhancement after DOY 17. The cross-correlation coefficients between the auroral revised power averaged over 0.2 day and the solar wind model show a weak correlation with maximum correlation coefficients of 0.28 for the 1-D model and 0.35 for the 3-D model with a 1 day lag, despite the ambiguity of the solar wind models. The solar wind radial velocity is small ( $<500 \text{ km s}^{-1}$ , Figure 4h) in both models. The CR shows a less clear dependence on the solar wind parameters. The IMF sector is mainly “away,” i.e., directed outward from the Sun, before DOY  $\sim 3$  and after DOY 12, and becomes “toward” to the Sun at other times (Figure 4i). From the 1-D model, the sector is judged from the azimuthal component of the IMF. The IMF sectors estimated using the 1-D model are consistent with those obtained using the 3-D model (Figure 4i).

### 3.6. CML Dependence

The CML dependence of the auroral emission was examined in previous studies [e.g., *Livengood and Moos*, 1990]. We analyze it using the EXCEED data to evaluate its stability and to find whether the EXCEED parameters are consistent with it. The power emitted at wavelengths of  $138.5-144.8 \text{ nm}$  has a clear CML dependence reflecting the fraction of the aurora visible to an observer as seen in section 3.3 and Figure 5a. The  $\text{CR}_{\text{EXCEED}}$  with large signal-to-noise ratio ( $>1.5$ ) is shown as a function of CML in Figure 5b. Both the power and CR vary with the CML. The emitted power maximizes at CML of  $\sim 170^\circ$ . This is a bit lower than that the CML at which the visible auroral area maximizes. The  $\text{CR}_{\text{EXCEED}}$  maximizes at a slightly higher CML of  $\sim 260^\circ$ . The average and variance  $\sigma$  of these parameters in each CML bin are shown using the green line and its error bar, respectively. Because the variance of the emitted power is large enough compared to its errors (grey lines in Figure 4a), the deviation from the average behavior is significant. On the other hand, the error is large and comparable to the variance  $\sigma$  for the  $\text{CR}_{\text{EXCEED}}$  case for several points.

Next, we show the temporal variation of the CML dependence of the emitted power (Figure 5c) and  $\text{CR}_{\text{EXCEED}}$  (Figure 5d). The differences from the average values in each CML bin are shown using the color maps, where the white parts indicate no available data. It is difficult to find specific CML dependences of the auroral brightness enhanced events in this dataset. For example, a short-term enhancement occurs at CMLs of  $120^\circ-180^\circ$  on DOY 11 (small enhancement by  $\sim \sigma$ ), and at CMLs of  $150^\circ-250^\circ$  on DOY 14 (large enhancement by  $> 2\sigma$ ). Long-term events, e.g., until DOY-5, DOY 1-2, and DOY 17-27, show enhancements over large CML ranges. The enhancement of the flux ratio by  $> \sigma$  above the averaged profile also does not show a clear CML dependence.

The CML offsets of the peak emitted power and CR values from the auroral visibility profile exhibit the same trends as reported by *Livengood and Moos* [1990], i.e., the peak emitted power at a lower CML and the peak



**Figure 6.** (a) Relationship between the power emitted at wavelengths of 138.5–144.8 nm and  $CR_{EXCEED}$  and (b) relationship between the scaled emitted power (scaled for visibility of the auroral region) and  $CR_{EXCEED}$ . The observed relations are shown using the black points.

reaching a maximum  $CR_{EXCEED}$  of  $\sim 3$ . At a power greater than 25 GW, the upper edge of the envelope decreases, seeming to approach  $CR_{EXCEED}$  of  $\sim 1.4$  for the greatest emitted power. This behavior suggests that the large emitted power events are mainly caused by primary particles with low average energies that are not significantly attenuated by the hydrocarbons. A few points with lower  $CR_{EXCEED}$  values appear at the lowest power values and others outside the triangular concentration have large uncertainties. This is comparable with the STIS results (Figure 1f) with less scatter because of the spatial integration (Figure 3c and section 3.3).

#### 4. Conclusions

Auroral spectra with a good time resolution ( $\sim 10$  min) and a long coverage of over  $\sim 40$  days obtained using Hisaki/EXCEED provide a unique opportunity to investigate the temporal variation of Jupiter's auroral parameters. The brightness-CR relation obtained from the EXCEED observations is compared with that obtained from the spatially resolved STIS observations. The main results from the observations from the end of 2013 to January 2014 are summarized as follows.

1. The EXCEED results are consistent with the STIS and previous observations in their auroral spectral profile, CML dependence, and CR-brightness distribution despite the limitation of the different wavelength coverage and large field of view of EXCEED.
2. The enhancement of the auroral brightness by a factor of 2–5 over short-duration ( $<$  one planetary rotation) and long-duration ( $>$  one planetary rotation) intervals are observed at both wavelength bands that are absorbed and unabsorbed by hydrocarbons. Therefore, compared to the brightness variation, the CR is relatively constant during this enhancement. Because the temporal variations of the solar wind dynamic pressure are different between the 1-D and 3-D models, decisive assessments of correlations between the aurora and solar wind parameters cannot be made. This study still suggests that the long-term large enhancement of the auroral power could be correlated with the large solar wind dynamic pressure enhancement. Further statistical survey is planned as future work.

CR at a higher CML. An enhancement of the CR in the dawnside region ( $\sim 8$  h magnetic local time) was reported using the spatially resolved HST spectral analysis by Gustin *et al.* [2004]. The distorted northern main auroral oval provides a better view of the dawnside at CMLs larger than  $\sim 200^\circ$ , which might cause the shift in the CML of the peak CR.

#### 3.7. CR-Brightness Relations From the EXCEED Observations

The CR-power relations from the EXCEED observations are shown in Figure 6a. Because the high-latitude-integrated power varies with the CML owing to the auroral aperture, we change the power by multiplying by the factor (the maximum visible auroral length in the all CML)/(the visible auroral length at instantaneous CMLs). The CR-power relation is modified as shown in Figure 6b. The distribution in the auroral emitted power and  $CR_{EXCEED}$  map exhibits a triangular envelope. The lower edge is at  $CR_{EXCEED} = \sim 0.7$ –1.5 with the emitted power increase from  $\sim 10$  to 70 GW. The upper edge of the envelope increases linearly with the emitted power from  $\sim 8$  to 25 GW,

3. The variability of the integrated auroral emitted power over the polar region is mainly attributed to the rotation of Jupiter's main auroral oval with the planet. The auroral CR also varies with a low dependence on the planetary rotation phase. A clear CML dependence of either short- or long-term auroral emitted power enhancements is not identified from the averaged CML data.

The quantitative estimation of the auroral parameters and further exploration of electron origin are described in the companion paper [Tao *et al.*, 2016]. The EXCEED observations and these findings cover spatially integrated auroral features and therefore represent the activity of the auroral region as a whole. Localized auroral features should exist, as shown by previous auroral observations [e.g., Gérard *et al.*, 2014], which will be one of the targets of upcoming Juno observations. In addition, examination of these characteristics is also planned using the next EXCEED Jupiter observing season, from the end of 2014 to early 2015 and beyond. Comparison with direct solar wind monitoring by Juno during its cruising phase is also expected.

## Appendix A

The relation between color ratios  $CR_{EXCEED}$  and  $CR_{STIS}$ , which are referring to different wavelength ranges, are represented by the hydrocarbon absorption cross sections [Gustin *et al.*, 2002]. The observed spectral intensity including absorption by dominant absorber  $CH_4$  can be expressed in terms of absorption cross section  $\sigma$  of  $CH_4$ ,  $CH_4$  column density  $N_{CH_4}$ , and unabsorbed height-integrated spectrum intensity  $I'$  as

$$I'_{(138.5-144.8 \text{ nm})} = I'_{(138.5-144.8 \text{ nm})} \exp(-N_{CH_4} \sigma_{(138.5-144.8 \text{ nm})}). \quad (A1)$$

Using this and similar relations for other wavelength ranges, we obtain the following expressions for the CRs as

$$\begin{aligned} CR_{STIS} &= I'_{(155-162 \text{ nm})} / I'_{(123-130 \text{ nm})} \\ &= I'_{(155-162 \text{ nm})} / I'_{(123-130 \text{ nm})} \exp\{-N_{CH_4}(\sigma_{(155-162 \text{ nm})} - \sigma_{(123-130 \text{ nm})})\}, \end{aligned} \quad (A2)$$

$$\begin{aligned} CR_{EXCEED} &= I'_{(138.5-144.8 \text{ nm})} / I'_{(126.3-130 \text{ nm})} \\ &= I'_{(138.5-144.8 \text{ nm})} / I'_{(126.3-130 \text{ nm})} \exp\{-N_{CH_4}(\sigma_{(138.5-144.8 \text{ nm})} - \sigma_{(126.3-130 \text{ nm})})\}. \end{aligned} \quad (A3)$$

Substituting  $N_{CH_4}$  obtained from equation (A2) into equation (A3),

$$CR_{EXCEED} = I'_{(138.5-144.8 \text{ nm})} / I'_{(126.3-130 \text{ nm})} \times \{CR_{STIS} I'_{(123-130 \text{ nm})} / I'_{(155-162 \text{ nm})}\}^\beta \quad (A4)$$

$$\beta \equiv (\sigma_{(138.5-144.8 \text{ nm})} - \sigma_{(126.3-130 \text{ nm})}) / (\sigma_{(155-162 \text{ nm})} - \sigma_{(123-130 \text{ nm})}), \quad (A5)$$

where  $I'_{(155-162 \text{ nm})} / I'_{(123-130 \text{ nm})} = 1.1$  [e.g., Grodent *et al.*, 2001; Gérard *et al.*, 2014],  $I'_{(138.5-144.8 \text{ nm})} / I'_{(126.3-130 \text{ nm})} = 0.59$  (estimated using the STIS spectra), and  $\sigma_{(126.3-130 \text{ nm})} = 1.73 \times 10^{-17} \text{ cm}^2$ ,  $\sigma_{(138.5-144.8 \text{ nm})} = 5.70 \times 10^{-19} \text{ cm}^2$ ,  $\sigma_{(123-130 \text{ nm})} = 1.74 \times 10^{-17} \text{ cm}^2$ , and  $\sigma_{(155-162 \text{ nm})} = 5.33 \times 10^{-24} \text{ cm}^2$  are derived from Parkinson *et al.* [2006].

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