

# The estimation of D-region electron densities from riometer data

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**Abstract.** At high latitude the hard electron precipitation associated with auroral activity is a major source of ionization for the D-region, one consequence being the absorption of radio waves. Direct measurements of the D-region electron density are not readily available, however. This paper investigates the relationship between the electron density at altitudes between 100 and 70 km and the total radio absorption observed with a riometer, with a view to using the latter to predict the former. Tables are given of the median electron density corresponding to 1 dB absorption at 27.6 MHz for each hour of the day, and it is shown that at certain heights the estimates will be accurate to within a factor of 1.6 on 50% of the occasions. A systematic variation with time of day is probably associated with a progressive hardening of the typical electron spectrum during the morning hours. There is also evidence for a seasonal effect possibly due to seasonal variations of the mesosphere.

**Key words.** Ionosphere (auroral ionosphere) – Radio science (ionospheric propagation; instruments and techniques)

## 1 Introduction

In the auroral regions the principal source of D-region ionization is the flux of energetic electrons during auroral activity. There are, of course, contributions from other sources, such as solar electromagnetic emissions (particularly Lyman- $\alpha$  and X-rays), and these are enhanced during solar flares. Solar protons cause ionization down to relatively low altitudes during solar proton events, but such events are much less frequent than auroral precipitation. Overall, the electron flux at energies from a few tens of keV to a few hundred keV may be assumed to be the dominant source most of the time, contributions from other ionization sources being generally less significant.

Direct observations in the D-region may be made by rocket-borne instruments, and although these techniques are

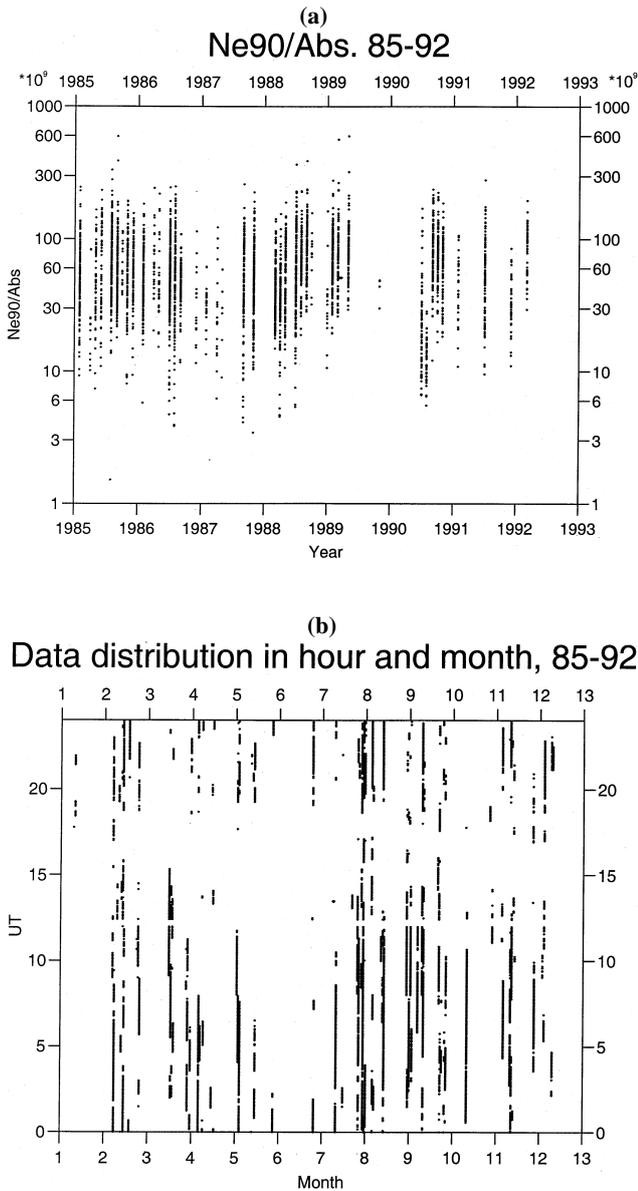
well established (Mechtly, 1974; Jacobsen and Friedrich, 1979), they are expensive and the data are limited because each flight produces only a single profile. Since the mid 1970s, and more particularly, since the EISCAT system began to operate in the early 1980s, it has been possible to measure electron densities in the D-region by incoherent scatter radar (Hunsucker, 1974; Ranta et al., 1985; Devlin et al., 1986). Under good conditions it is possible to obtain accurate measurements at 1 km height resolution and 10 s time resolution. Radar is considerably less sensitive than the rocket-borne probes but it may provide continuous runs of data over several hours. However, the radar data are still limited because incoherent scatter radars are few in number and immobile, and they are also expensive to operate.

Another way of detecting enhancements in the D-region is by measuring the absorption of a radio signal that has passed through the medium. This is the purpose of the riometer technique in which the intensity of the cosmic radio noise at some frequency, usually between 20 and 60 MHz, is monitored using a stable receiver at the ground (Hargreaves, 1969). The riometer is relatively cheap and simple to operate. Many of them have been, and still are, in use.

This paper addresses the question of whether riometer data may realistically be used to indicate the electron density at a selected altitude. The depth to which an incoming electron penetrates into the atmosphere depends on its initial energy (Rees, 1963), the more energetic ones depositing their energy at lower altitudes. The resulting electron-density profile, therefore, depends on the original spectrum of the incident particles. Then, if the electron density is  $N_e(h)$ , the absorption of a radio wave of angular frequency  $\omega = 2\pi f$  is

$$A \text{ (dB)} = 4.6 \times 10^{-5} \int \frac{N_e(h)\nu(h)}{\omega^2 + \nu(h)^2} dh, \quad (1)$$

(Hargreaves, 1969) where  $\nu(h)$  is the electron-neutral collision frequency. While  $\nu(h)$  is a property of the atmosphere (i.e. proportional to pressure),  $N_e(h)$  depends on the incoming electron flux over a band of energies and thus, also on the particle energy. Since the spectrum of the ionizing electrons is not constant, one cannot expect to find a precise rela-

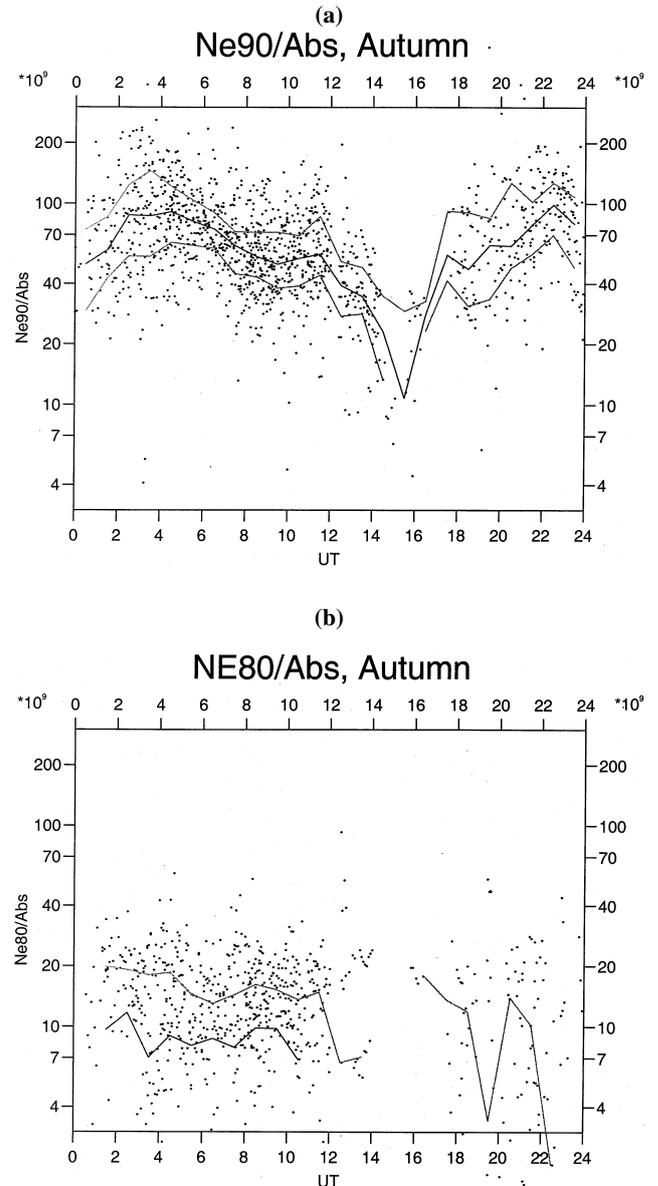


**Fig. 1.** (a) Value of 90 km electron density corresponding to 1 dB absorption for the whole data set between 1985 and 1992. (b) Distribution of data points with month and time of day.

relationship between radio absorption and electron density. But neither is the spectrum entirely random. Therefore, we may reasonably ask to what extent, if any, ground-based absorption measurements may be used to indicate D-region electron density, and, if so, what will be the accuracy of these estimates.

## 2 The present investigation and data used

Friedrich and Kirkwood (2000) have published average electron-density profiles corresponding to selected absorption levels. Most of the data came from EISCAT for the



**Fig. 2.** Values of  $N_e/A$  against time of day for (a) 90 km and (b) 80 km.

greater heights and from rockets for the lower ones below about 80 km. The profiles were divided according to whether or not the D-region was illuminated by the Sun. The present study seeks to extend those results in two directions: by specifying the likely accuracy of the estimates, and by looking at the change from night to day in greater detail.

The electron density values are taken from EISCAT UHF common programme runs over the years 1985 to 1992, inclusive. The heights 70 to 100 km were selected at 5 km intervals. Each of the electron-density profiles was inspected to eliminate poor data. In cases where the return was too weak for a reliable electron density to be determined, the value zero was substituted. The absorption measurements are taken as averages from four riometers within 10 km of the radar

**Table 1.** Median electron density corresponding to 1 dB absorption for each height and each hour of the day, using data from all years and months. The \* indicates an upper quartile

UT	No. of Values	Height						
		100 km	95 km	90 km	85 km	80 km	75 km	70 km
0.5	188	.231E12	.129E12	.414E11	.745E10	.219E10*	.000E00*	.000E00*
1.5	195	.152E12	.112E12	.481E11	.100E11	.490E10*	.000E00*	.000E00*
2.5	237	.179E12	.129E12	.617E11	.148E11	.960E10*	.000E00*	.000E00*
3.5	226	.163E12	.120E12	.551E11	.161E11	.326E10	.000E00*	.000E00*
4.5	254	.142E12	.128E12	.805E11	.273E11	.589E10	.243E10*	.000E00*
5.5	271	.980E11	.895E11	.714E11	.255E11	.640E10	.437E10*	.000E00*
6.5	258	.965E11	.901E11	.688E11	.266E11	.734E10	.463E10*	.000E00*
7.5	252	.852E11	.771E11	.625E11	.236E11	.702E10	.373E10*	.000E00*
8.5	260	.643E11	.636E11	.539E11	.251E11	.855E10	.447E10*	.000E00*
9.5	243	.747E11	.657E11	.518E11	.235E11	.774E10	.527E10*	.000E00*
10.5	236	.720E11	.625E11	.500E11	.202E11	.572E10	.495E10*	.000E00*
11.5	198	.768E11	.660E11	.512E11	.191E11	.124E11*	.000E00*	.000E00*
12.5	114	.624E11	.572E11	.356E11	.132E11	.676E10*	.000E00*	.000E00*
13.5	156	.581E11	.449E11	.328E11	.764E10	.000E00*	.000E00*	.000E00*
14.5	90	.571E11	.434E11	.305E11	.903E10	.485E10*	.000E00*	.000E00*
15.5	37	.407E11	.293E11	.222E11	.603E10	.000E00*	.000E00*	.000E00*
16.5	27	.551E11	.405E11	.280E11	.188E11*	.000E00*	.000E00*	.000E00*
17.5	45	.181E12	.789E11	.386E11	.137E11	.936E10*	.000E00*	.000E00*
18.5	65	.190E12	.889E11	.455E11	.172E11	.989E10*	.000E00*	.000E00*
19.5	133	.144E12	.841E11	.355E11	.165E11*	.000E00*	.000E00*	.000E00*
20.5	172	.117E12	.648E11	.340E11	.557E10	.290E10*	.000E00*	.000E00*
21.5	236	.176E12	.116E12	.516E11	.123E11	.536E10*	.000E00*	.000E00*
22.5	218	.191E12	.142E12	.690E11	.134E11	.395E10*	.000E00*	.000E00*
23.5	186	.248E12	.157E12	.518E11	.863E10	.000E00*	.000E00*	.000E00*

site at Ramfjordmoen, Norway (69.6° N, 19.2° E,  $L = 6.2$ ), specifically two each at Ramfjordmoen and at Lavangsdalen. The absorption values are converted to 27.6 MHz, X-mode and vertical incidence (Friedrich et al., 2002). Both data sets are at a 5-min interval. Since the riometer data are relatively less accurate at small values, and to reduce the effect of solar contributions to the absorption (which could amount to a few tenths of a dB at noon), only values of at least 0.5 dB are used in the following analysis.

### 3 Data analysis

Since the absorption at a given height is proportional to the electron density (Eq. 1), we work with the ratio  $N_e(h)/A$ , where  $N_e(h)$  is the electron density in  $\text{m}^{-3}$  at height  $h$  and  $A$  is the total absorption at 27.6 MHz in dB. Figure 1a shows the values of  $N_e(90)/A$  over the whole data set against the year. Its value spreads by about a factor of 10, but there is no indication of a systematic variation from year to year. We shall, therefore, take all the years together. Figure 1b shows the distribution of data points with month and time of day. This distribution is less even than we might wish. The gap in the afternoon corresponds with the low occurrence of auroral absorption at that time of day, which is well known (Harg-

reaves and Cowley, 1967), but the gaps during the summer period, and during the winter, from Christmas to the end of January, are plainly observational and due to the schedules of radar operations.

As an example of the results, Fig. 2 shows the values of  $N_e/A$  for 90 and 80 km for the autumn period (months 8.25–11.25) against time of day, with medians and quartiles superimposed. At 90 km there is a marked variation with time of day, particularly during the morning hours. That variation is not apparent at 80 km. In that case, however, medians were obtained only for the hours 01:00 to 11:00 UT. At other times only the upper quartiles are shown. The lower quartiles are lost in the zeros of the data set, and are not plotted; they fall below the limit of measurement.

Table 1 summarizes the median values of electron density corresponding to 1 dB absorption for each hour of the UT day. When no median was obtained, the upper quartile is marked with an asterisk. This table uses the values from all months. The relative lack of data in summer and winter makes it more difficult to look for seasonal variations, but we will return to this point in a later section.

When the value zero appears in Table 1, it means that even the upper quartile was below the threshold of measurement. We can estimate that threshold from the distribution of measured electron density (Fig. 3). Since each value is

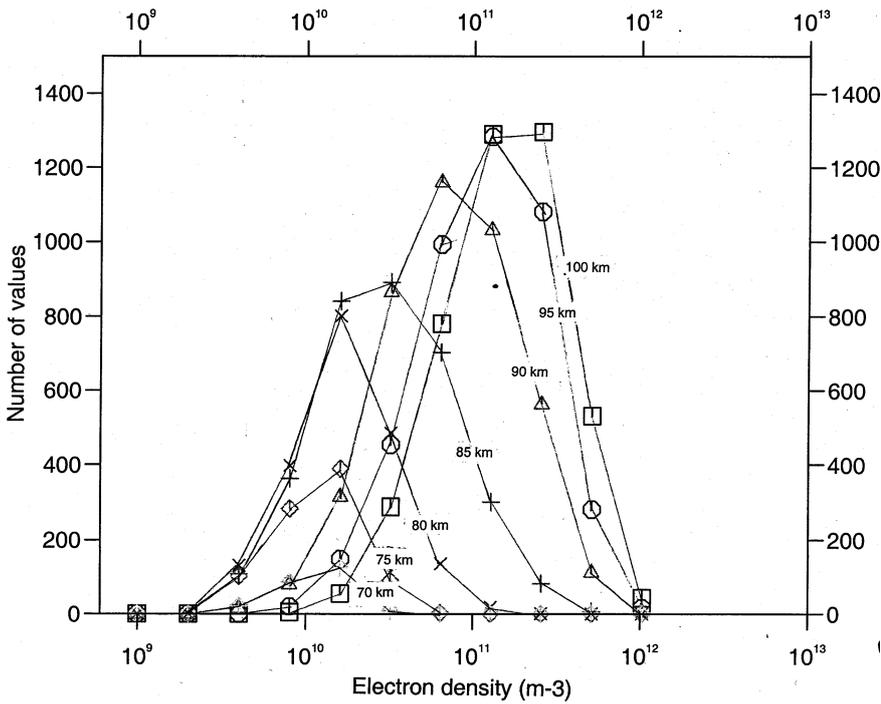


Fig. 3. Histograms of electron density from 100 km to 70 km.

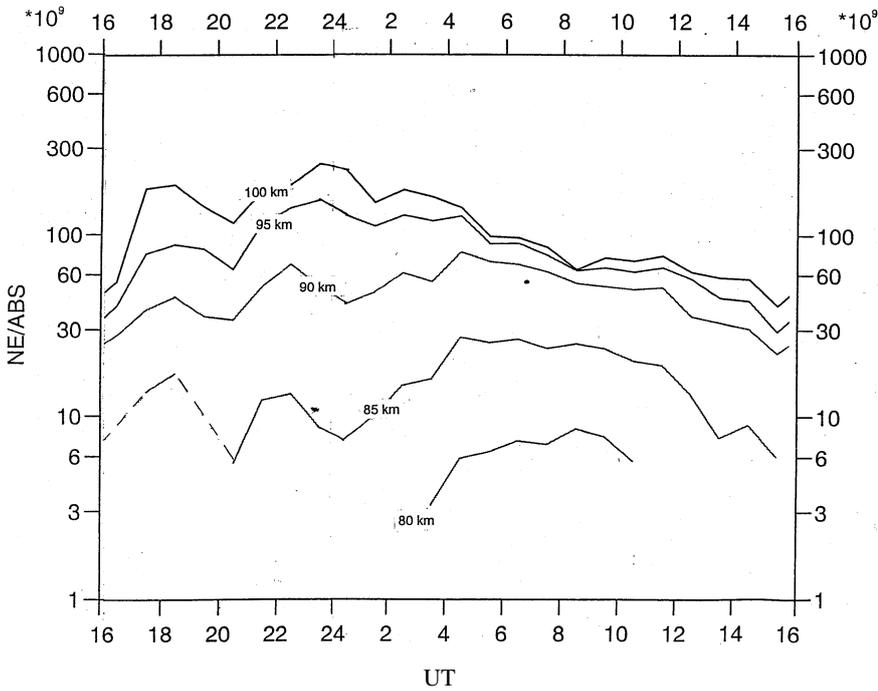


Fig. 4. Medians for whole data set, 1985-1992. Variation of median electron density (for 1 dB absorption) with time of day. At EISCAT local magnetic midnight is about 21:30 UT. The data show a daily cycle starting about six hours before magnetic midnight.

taken from a complete profile covering 70 to 100 km, there should be the same number of readings for each altitude. But whereas we see similar distributions, though with different medians, for 100, 95 and 90 km, fewer values appear at and below 80 km. Clearly, the numbers fall off significantly below about  $10^{10} \text{ m}^{-3}$ , the missing values having been designated as zero during the initial editing of the data. This may be checked by counting the number of zeros in each height

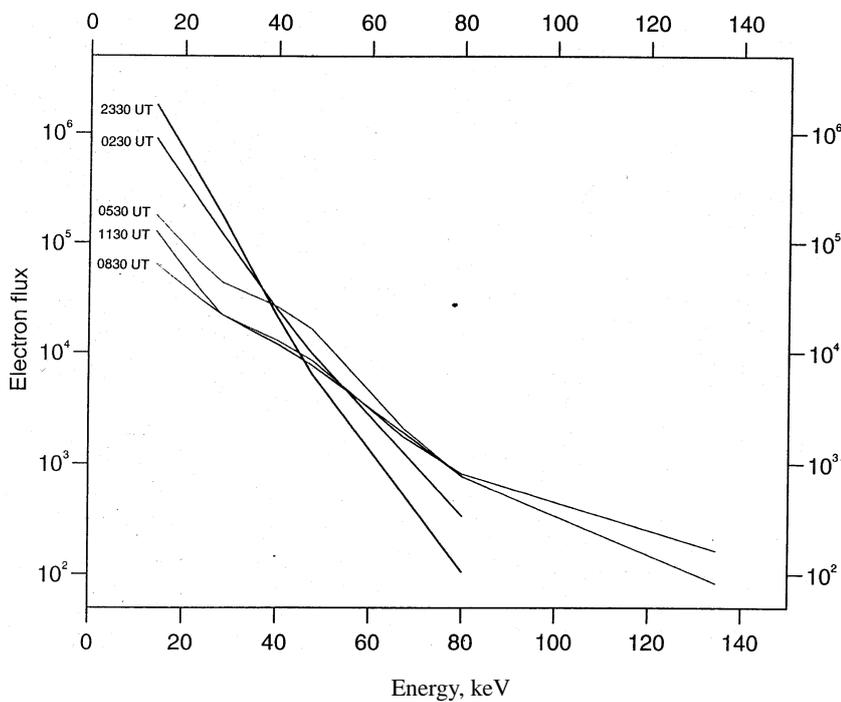
group (Table 2). Down to 90 km fewer than 5% of the readings fell below threshold, whereas at 70 km only 6% were above it. It must be remembered that this result is for conditions where the 27.6 MHz absorption was at least 0.5 dB. A different absorption limit would produce a different distribution in Table 2, though the absolute threshold of approximately  $10^{10} \text{ m}^{-3}$  would presumably be the same.

**Table 2.** Numbers of electron-density values and of zeros in the data set

Height	100 km	95 km	90 km	85 km	80 km	75 km	70 km
Number of electron density values	4295	4282	4149	3300	1968	870	244
Number of zeros	1	14	147	996	2328	3426	4152
% of zeros	0.02	0.33	3.4	23	54	79	94

**Table 3.** Spread of values in terms of median and quartiles

Height	100 km	95 km	90 km	85 km	80 km	75 km	70 km
<u>Median</u>							
Lower quartile	1.59	1.63	1.60	2.44	–	–	–
Upper quartile	1.52	1.49	1.50	1.71	2.12	–	–
<u>Upper quartile</u>							
Lower quartile	2.42	2.43	2.40	4.17	–	–	–



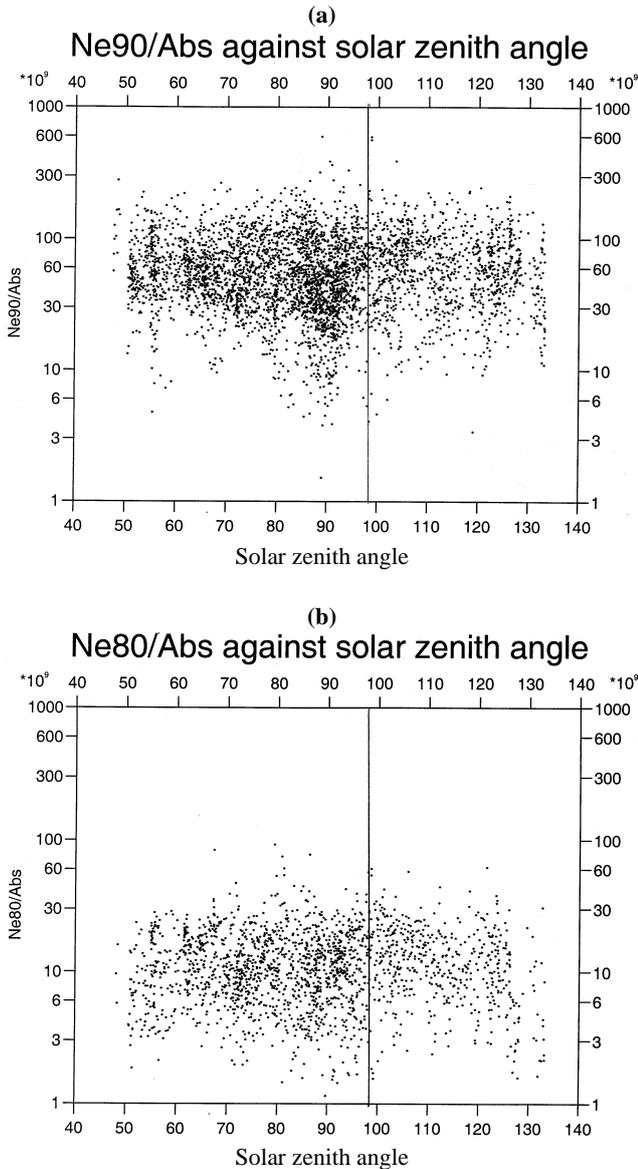
**Fig. 5.** Spectra deduced from the median electron-density profiles from night to morning.

#### 4 Prediction model for the whole year

The overall results are summarized in Fig. 4, which gives the hourly medians for each height. Note that at 75 km and below the medians were all below the radar threshold. At 80 km the values were above the threshold only between 03:00 and 11:00 UT. Systematic trends are apparent. At 100 and 95 km the ratio  $N_e/A$  decreases steadily from the midnight sector

into the daytime, whereas the lower heights show an increase followed by a decrease during the morning hours. The dips at 20:00 UT are probably statistical.

In order to serve as a prediction model, the information in Fig. 4 needs to be supplemented by a knowledge of the spread of the values contributing to each median. Table 3 shows the relationship between quartiles and median at the various heights. At 90 km and above, the ratios (upper quar-



**Fig. 6.**  $N_e/A$  against the solar zenith angle.

tile)/median and median/(lower quartile) are both about 1.5 or 1.6. This means that there is a 50% chance that a single estimate will be within 1.6 of the true value. The spread is somewhat worse at the lower heights.

## 5 Significance of the diurnal variation

The electron-density profile may be inverted to give an indication of the spectrum of incoming particles that produced it (Hargreaves and Devlin, 1990). Some assumptions have to be made about the atmospheric composition and about the effective recombination coefficient as a function of altitude. Although the absolute values may be questioned, this approach is useful for interpreting changes of profile in terms of a changing spectrum. Treating in this way the median

profiles for 23:00–24:00 UT and each 3 h to 11:00–12:00 UT (Fig. 5) shows a hardening of the typical spectrum with time. The earliest spectrum is almost exponential with characteristic energy 6.5 keV, a value consistent with other determinations for that time of day (Devlin et al., 1986). Later in the morning, the characteristic energy is considerably greater, about 16 keV. This hardening is also consistent with previous observations of individual events (Hargreaves and Devlin, 1990).

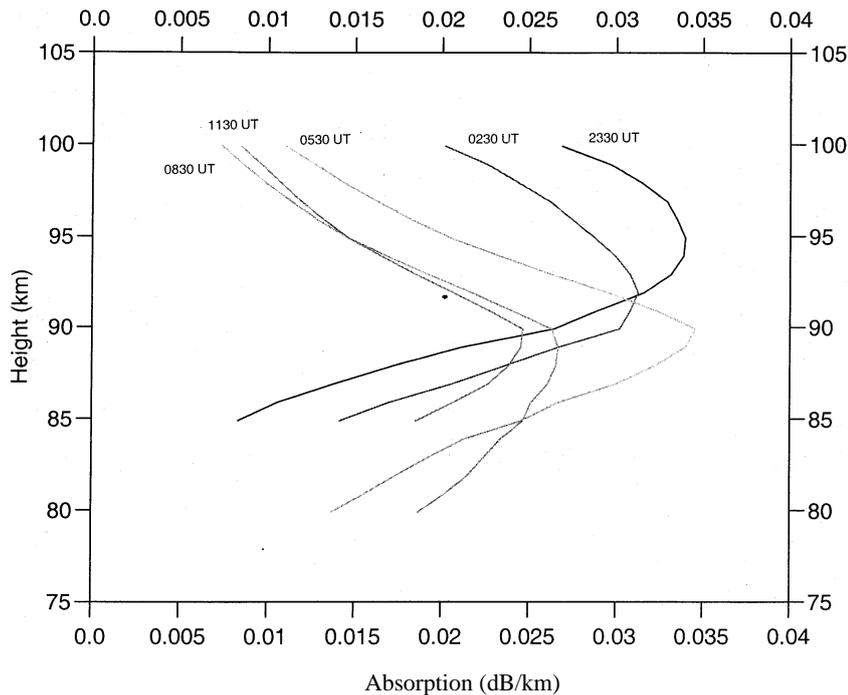
## 6 The question of photodetachment

In polar-cap absorption events (PCA), which result from the incidence of solar protons, the level of absorption changes drastically between day and night, according to whether or not the ionosphere is sunlit (Bailey, 1957). Such a transition is expected because the effective electron recombination in the lower ionosphere increases drastically in the absence of daylight, due to attachment of electrons to neutrals and the enhanced formation of cluster ions. PCA occurs lower in the atmosphere than auroral absorption (AA), and attempts to identify the same effect in AA (most notably the analysis by Armstrong et al., 1977) have failed to find one.

In Fig. 6 we show the quantity  $N_e/A$  against the solar zenith angle ( $\chi$ ), for two altitudes. It will be noted that there is no sharp change between  $\chi = 90$  and  $100^\circ$ , as would be the case in PCA. We note, further, that the absorption profiles (Fig. 7) corresponding to the spectra of Fig. 5 peak between 90 and 95 km; the contribution at and below 75 km is likely to be small. Incoherent-scatter radar studies of PCA, however, show that most of the day/night variation of electron density in those events is strongly height dependent, and is rather small above 75 km (Reagan and Watt, 1976; Collis and Rietveld, 1990; Hargreaves et al., 1993). For large ionisation rates, such as those considered in the present context (absorption larger than 0.5 dB), the cluster ledge is most likely to be always below 80 km. We conclude that the diurnal variation shown in Fig. 4 and Table 1 is a consequence of a changing electron spectrum, not of solar illumination.

## 7 Seasonal variation

The scheme described above takes into account daily but not seasonal variations. To include both variables requires the data to be sorted into 288 boxes to cover each hour of each month, and since the distribution is uneven some of these boxes will contain too few data points for a significant result. Therefore, median values of the ratio  $N_e/A$  were derived only for those boxes containing at least 30 points. These values are given in Tables A1–5, covering heights 100 to 80 km, in the Appendix. As before, zero means that the median value was below the threshold of observation. Based on the general statistics (Sect. 3), it is estimated that about half the medians in the tables will not be within 10% of the true value. In comparing values, therefore, differences less



**Fig. 7.** Partial absorption profiles corresponding to the median electron-density profiles in the night-to-morning sector.

than 10% should not be considered significant. Differences of more than 50% are likely to be significant.

The best coverage across all months is between 04:00 and 07:00 UT. We can also compare the variations during the day by taking together November to March (as winter) and July to August (as summer), though both these groups lack data from 14:00 to 20:00 UT. The reader may like to draw his or her own conclusions, but the authors offer the following comments:

1. At 100 and 95 km, the value of electron density corresponding to 1 dB absorption is relatively high in summer around 00:00 UT, but relatively low for several hours before that.
2. At 90 km, the summer maximum at 00:00 UT has vanished, though values are still low before that time.
3. At 85 km, the low values before 00:00 UT are now below the threshold. There is now also a dip during the early morning (05:00 to 07:00 UT).
4. At 80 km, all summer values are low relative to those at other times of the year.

For a given spectrum of incoming energetic electrons, the ionization rate and the effective recombination coefficient, and, therefore, the ionization profile, are all influenced by the vertical profiles of density and pressure in the mesosphere. Further, the resulting radio absorption depends on the electron-neutral collision frequency, which is also pressure dependent. Direct measurements of atmospheric density at high latitude show a reduction in summer at 100 and 95 km,

and a summer increase at 85 and 80 km. The seasonal variations indicated by Tables A1–5 may well be a consequence of mesospheric changes, but due to the inherent limitations of the data they should be regarded as indications rather than definitive determinations.

## 8 Conclusions and discussion

1. An analysis of incoherent-scatter and riometer data over a period of 8 years, including more than 4000 individual measurements, shows that an observation of auroral radio absorption may be interpreted to estimate the approximate electron density at heights between 100 and 80 km. The results presented apply only if the absorption is at least 0.5 dB. The results are given in terms of the medians, and upper and lower quartiles of the electron density, corresponding to 1 dB absorption measured with a 27.6 MHz riometer pointed vertically.
2. The uncertainty of the procedure is such that at the greater heights (100, 95 and 90 km) the estimate will be accurate to within a factor of 1.6 on 50% of the occasions. At lower heights the error appears to be rather worse.
3. The sensitivity of the EISCAT UHF radar when using 5-min averaging is found to be  $10^{10} \text{ m}^{-3}$ , or a little below, in the D-region. Due to this limit no median values of the electron density corresponding to 1 dB absorption could be determined at or below 75 km, and the 80-km values are restricted to one period of the day. The upper

quartile could be specified at 80 and 75 km, but not at 70 km.

4. Since the typical spectrum of incoming energetic electrons varies with the time of day, the estimates of electron density show a diurnal variation. There is no evidence that solar illumination has any effect at 80 km or above.
5. Dividing the data according to both time-of-day and season suggests that a seasonal effect is also present, but due to data limitations the actual values should be treated with caution.

The absorption values used in this investigation are those actually measured by a simple riometer system using a vertical 5-element Yagi antenna. It is estimated that such an antenna records absorption about 20% greater than would be measured by an ideal system having a zenithal pencil beam. This factor should be taken into account where appropriate.

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## Appendix

The following tables give the data for heights of 100 to 80 km divided by both time of day and season. The values are median electron densities corresponding to 1 dB absorption at 27.6 MHz, and are only given for those boxes having at least 30 data points. Each box is designated according to the central time and month concerned; e.g. “time 5.5, month 6.5” means 05:00 to 06:00 UT in the month of July.

**Table A1.** Median Electron density at 100 km for 1 dB absorption

Month	0.5	1.5	2.5	3.5	4.5	5.5	6.5	7.5	8.5	9.5	10.5	11.5
<u>Hour (UT)</u>												
0.5		180					341					
1.5							189	144				
2.5		182						140				
3.5								78.6	178			
4.5					155		143		142		166	
5.5			128	66.5	88.7		103		109		128	
6.5		58.9	104				118		103		115	
7.5							111	62.5	76.3		81.2	
8.5							101	46.6	75.3		52.0	
9.5			54.6				97.8		91.4			
10.5							114	76.6	70.5			
11.5								111	76.6			
12.5												
13.5							95.7		56.9			
14.5												
15.5												
16.5												
17.5												
18.5												
19.5												
20.5		287					64.0					
21.5		200					84.6					102
22.5		208					103				244	
23.5							280					

**Table A2.** Median Electron density at 95 km for 1 dB absorption

Month	0.5	1.5	2.5	3.5	4.5	5.5	6.5	7.5	8.5	9.5	10.5	11.5
<u>Hour (UT)</u>												
0.5		95.1					208					
1.5							123	87.7				
2.5		131						116				
3.5								71.0	142			
4.5					132		147		136		136	
5.5			96.4	61.1	95.4		75.3		99.1		105	
6.5		58.7	70.1				97.8		82.4		104	
7.5							113	55.3	64.3		78.0	
8.5							105	52.5	64.4		48.4	
9.5			50.9				80.7		72.4			
10.5							93.1	61.5	58.7			
11.5								78.3	61.9			
12.5												
13.5							53.8		41.5			
14.5												
15.5												
16.5												
17.5												
18.5												
19.5												
20.5		111					47.5					
21.5		155					58.1					79.2
22.5		160					85.4				171	
23.5							190					

**Table A3.** Median Electron density at 90 km for 1 dB absorption

Month	0.5	1.5	2.5	3.5	4.5	5.5	6.5	7.5	8.5	9.5	10.5	11.5
<u>Hour (UT)</u>												
0.5		40.6					41.1					
1.5							44.4	39.2				
2.5		61.7						72.5				
3.5								51.4	85.5			
4.5					78.4		108		82.1		57.6	
5.5			48.3	41.2	84.8		42.0		71.5		79.9	
6.5		53.7	58.8				61.4		74.7		71.0	
7.5							78.6	61.2	60.3		56.9	
8.5							68.2	63.3	52.3		37.9	
9.5			48.4				49.2		48.6			
10.5							53.4	57.7	50.0			
11.5								57.9	51.4			
12.5												
13.5							30.9		32.8			
14.5												
15.5												
16.5												
17.5												
18.5												
19.5												
20.5		49.5					21.1					
21.5		87.6					30.7					34.0
22.5		86.4					39.2				70.3	
23.5							46.9					

**Table A4.** Median Electron density at 85 km for 1 dB absorption

Month	0.5	1.5	2.5	3.5	4.5	5.5	6.5	7.5	8.5	9.5	10.5	11.5
<u>Hour (UT)</u>												
0.5		0					4.77					
1.5							6.40	7.04				
2.5		16.8						8.87				
3.5								13.3	26.0			
4.5					27.7		22.1		28.9		18.0	
5.5			26.1	18.1	34.6		9.52		21.7		41.2	
6.5		35.3	33.6				14.2		24.1		34.4	
7.5							21.3	7.31	20.2		33.3	
8.5							21.9	27.7	24.8		23.6	
9.5			28.7				12.0		22.7			
10.5							10.7	23.3	18.5			
11.5								18.1	20.0			
12.5												
13.5							0		11.8			
14.5												
15.5												
16.5												
17.5												
18.5												
19.5												
20.5		25.6					0					
21.5		27.3					0					5.77
22.5		29.5					0				13.9	
23.5							7.94					

**Table A5.** Median Electron density at 80 km for 1 dB absorption

Month	0.5	1.5	2.5	3.5	4.5	5.5	6.5	7.5	8.5	9.5	10.5	11.5
<u>Hour (UT)</u>												
0.5		0					0					
1.5							0	0				
2.5		0						0				
3.5								0	5.22			
4.5					7.81		2.55		6.84		4.01	
5.5			10.5	3.99	8.36		0		0		10.5	
6.5		12.4	9.77				0		9.62		8.77	
7.5							0	0	7.54		11.9	
8.5							4.20	8.03	9.33		9.09	
9.5			10.9				0		0			
10.5							0	7.91	0			
11.5								0	0			
12.5												
13.5							0		0			
14.5												
15.5												
16.5												
17.5												
18.5												
19.5												
20.5		9.71					0					
21.5		7.34					0					0
22.5		9.07					0				0	
23.5							0					