



Differences between CME-driven storms and CIR-driven storms

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[1] Twenty one differences between CME-driven geomagnetic storms and CIR-driven geomagnetic storms are tabulated. (CME-driven includes driving by CME sheaths, by magnetic clouds, and by ejecta; CIR-driven includes driving by the associated recurring high-speed streams.) These differences involve the bow shock, the magnetosheath, the radiation belts, the ring current, the aurora, the Earth's plasma sheet, magnetospheric convection, ULF pulsations, spacecraft charging in the magnetosphere, and the saturation of the polar cap potential. CME-driven storms are brief, have denser plasma sheets, have strong ring currents and Dst, have solar energetic particle events, and can produce great auroras and dangerous geomagnetically induced currents; CIR-driven storms are of longer duration, have hotter plasmas and stronger spacecraft charging, and produce high fluxes of relativistic electrons. Further, the magnetosphere is more likely to be preconditioned with dense plasmas prior to CIR-driven storms than it is prior to CME-driven storms. CME-driven storms pose more of a problem for Earth-based electrical systems; CIR-driven storms pose more of a problem for space-based assets.

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1. Introduction

[2] CME-driven and CIR-driven geomagnetic events differ, with the various forms of geomagnetic activity (ring current, aurora, convection, radiation belts, . . .) being manifested to different degrees in the two types of storms. In this paper a systematic comparison between the properties of CME-driven storms and CIR-driven storms is summarized. As will be seen, there are a large number of differences.

[3] It is not surprising that the two types of geomagnetic storms differ because the two types of drivers differ. Geomagnetic activity is chiefly driven by the value of vB_s of the solar wind [e.g., *Bargatze et al.*, 1986], where v is the solar wind speed and B_s is the southward component of the interplanetary magnetic field (in GSM coordinates). For CME drivers and CIR drivers, many things differ: the strength and morphology of the magnetic field, the temporal profile of the velocity v , the durations of the vB_s drivers, the presence of shocks, . . . [e.g., *Gonzalez et al.*, 1999]. Coronal mass ejections (CMEs) are preceded in time by a sheath of compressed solar wind, which itself is often preceded by an interplanetary shock. The strong magnetic field of the sheath can drive storms as well as can a strong magnetic field in the ejecta, and the interplanetary shock can add phenomenology to the geomagnetic storm on Earth. Additionally, the ejecta may contain a magnetic cloud. In this paper, CME-driven storms will refer to storms driven by all or some of the various components: shock, sheath,

ejecta, and cloud. Corotating interaction regions (CIRs) are followed in time by high-speed streams. Either the CIR or the high-speed stream or both can be drivers of storms. In this paper, CIR-driven storms refers to storms driven by either or both the CIR and the stream.

[4] The occurrence rate of CMEs peaks strongly during solar maximum [*Webb*, 1991; *Yashiro et al.*, 2004], so storms during solar maximum tend to be CME-driven storms [*Richardson et al.*, 2001]. The occurrence rate of well-formed CIR (with 27-day recurrence) peaks during the late declining phase of the solar cycle [*Mursula and Zeiger*, 1996], so storms in the declining phase tend to be CIR driven [*Richardson et al.*, 2001]. Note that interaction regions (IRs) produced by nonrecurrent high-speed streams can occur throughout the solar cycle [*Bobrov*, 1983; *L. Jian*, private communication, 2005] and can drive nonrecurring geomagnetic activity throughout the solar cycle [*Richardson et al.*, 2000], although the geomagnetic events that they drive are weaker than the recurring events of the declining phase [*Bobrov*, 1983].

[5] This study is a follow on to the multispacecraft study of *Denton et al.* [2006] which separately examined the magnetosphere's reaction to the passage of CMEs and CIRs. In that study, superposed epoch intervals of various quantities measured by spacecraft in the magnetosphere were constructed using the peak negative Dst as the zero epoch. Six of the differences listed in Table 1 are findings of the *Denton et al.* superposed epoch study. The *Denton et al.* [2006] study is itself a follow-up to a *Denton et al.* [2005] study in which the magnetospheric evolution during 283 geomagnetic storms was examined.

[6] The differences between CME-driven storms and CIR-driven storms are collected into Table 1. The differences are discussed individually in section 2. For some

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Table 1. A Summary of Some of the Important Differences Between CME-Driven Storms (Shock, Sheath, Ejecta, Cloud) and CIR-Driven Storms (CIR, High-Speed Stream)

Phenomenon	CME-Driven Storms	CIR-Driven Storms
Phase of the solar cycle when dominant	solar maximum	declining phase
Occurrence pattern	irregular	27-day repeating
Calm before the storm	sometimes	usually
Solar energetic particles (SEP)	sometimes	none
Storm sudden commencement (SSC)	common	infrequent
Mach number of the bow shock	moderate	high
β of magnetosheath flow	low	high
Plasma-sheet density	very superdense	superdense
Plasma-sheet temperature	hot	hotter
Plasma-sheet O^+/H^+ ratio	extremely high	elevated
Spacecraft surface charging	less severe	more severe
Ring current (Dst)	stronger	weaker
Global sawtooth oscillations	sometimes	no
ULF pulsations	shorter duration	longer duration
Dipole distortion	very strong	strong
Saturation of polar-cap potential	sometimes	no
Fluxes of relativistic electrons	less severe	more severe
Formation of new radiation belts	sometimes	no
Convection interval	shorter	longer
Great aurora	sometimes	rare
Geomagnetically induced current (GIC)	sometimes	no

other properties of storms the differences could not be discerned within the scope of the present study; those phenomena are briefly discussed in section 3. The findings are summarized in section 4.

2. Differences

[7] The differences between CME-driven storms and CIR-driven storms are collected into Table 1. The category “CME-driven” includes driving by ejecta, CME sheaths, and magnetic clouds. The category “CIR-driven” also includes driving by the high-speed-stream that follows the CIR. Each row of Table 1 is described in detail in the various subsections of this section. For some storm properties, the differences between CME-driven storms and CIR-driven storms are well known and the appropriate literature is pointed out. For other properties, original work is done here to discern the differences.

2.1. Phase of the Solar Cycle

[8] It is well known that CIR-driven storms generally occur in the late declining phase of the solar cycle and that CME-driven storms tend to occur at solar maximum [Gonzalez *et al.*, 1999; Yermolaev and Yermolaev, 2002; Richardson *et al.*, 2001, 2002]. There are exceptions to this general trend. CMEs occur throughout the solar cycle, but their occurrence frequencies and their velocities are both greatest during solar maximum [Webb and Howard, 1994; Ivanov and Obridko, 2001; Gopalswamy *et al.*, 2004]. Also, 27-day-recurring high-speed streams can also occur in phases of the solar cycle other than the declining phase. Additionally, interaction regions (IRs) that are not 27-day recurring occur throughout the solar cycle [Bobrov, 1983; Richardson *et al.*, 2000], although their geomagnetic effectiveness is weaker on average than is the geoeffectiveness of CIRs.

2.2. Occurrence Rate

[9] CIRs (corotating interaction regions), when they occur, pass the Earth every ~ 27 days owing to the rotation of

the Sun. During the declining phase of the solar cycle, coronal hole structure on the solar surface is long-lived and simple, and as a result 27-day-recurring high-speed-stream-driven storms are often found. This is shown in Figure 1a. Here the Kp index is used to create a storm occurrence database. The Kp index is broken into 27.27-day-long intervals, one interval for each rotation of the Sun. The data is then plotted as the day during the 27.27-day-long interval (horizontal) versus the fractional year (vertical). The timelines are the thin black lines. Time intervals when $Kp \geq 4^+$ are plotted as large black squares. 27-day repeating storms appear as the vertical clusters of black points. Several of these storm groups can be seen in 1993–1995 during the declining phase. Another group can be seen in late 1999 and early 2000. The storms are replotted in Figure 1b without the timeline. The McPherron online catalog (R. McPherron, private communication, 2005) (www.igpp.ucla.edu/public/rmcpherr/CDAWEventLists/) of CIR stream interfaces at Earth for January 1994 through December 1996 is used to produce the yellow points in Figure 1b. As can be seen by their association with CIRs, the groups of periodic storms in Figure 1a are high-speed-stream-driven storms (CIR-driven storms). As can be seen in Figure 1a, throughout the solar cycle other storms occur randomly. The occurrence frequency of these random storms is higher in the 2000–2002 solar maximum years. Plotted in purple in Figure 1b are times when CMEs are passing the Earth (the Cane and Richardson catalog [Cane and Richardson, 2003]) and plotted in yellow are times when magnetic clouds are passing the Earth (the Lepping online catalog http://lepmfi.gsfc.nasa.gov/mfi/mag_cloud_pub1.html, see Lepping *et al.* [2005]). The Cane and Richardson catalog goes from January 1996 through December 2002 and the Lepping catalog goes from February 1995 through August 2003. As can be seen, the randomly occurring storms in Figure 1a are associated with CMEs and with the magnetic clouds within CMEs.

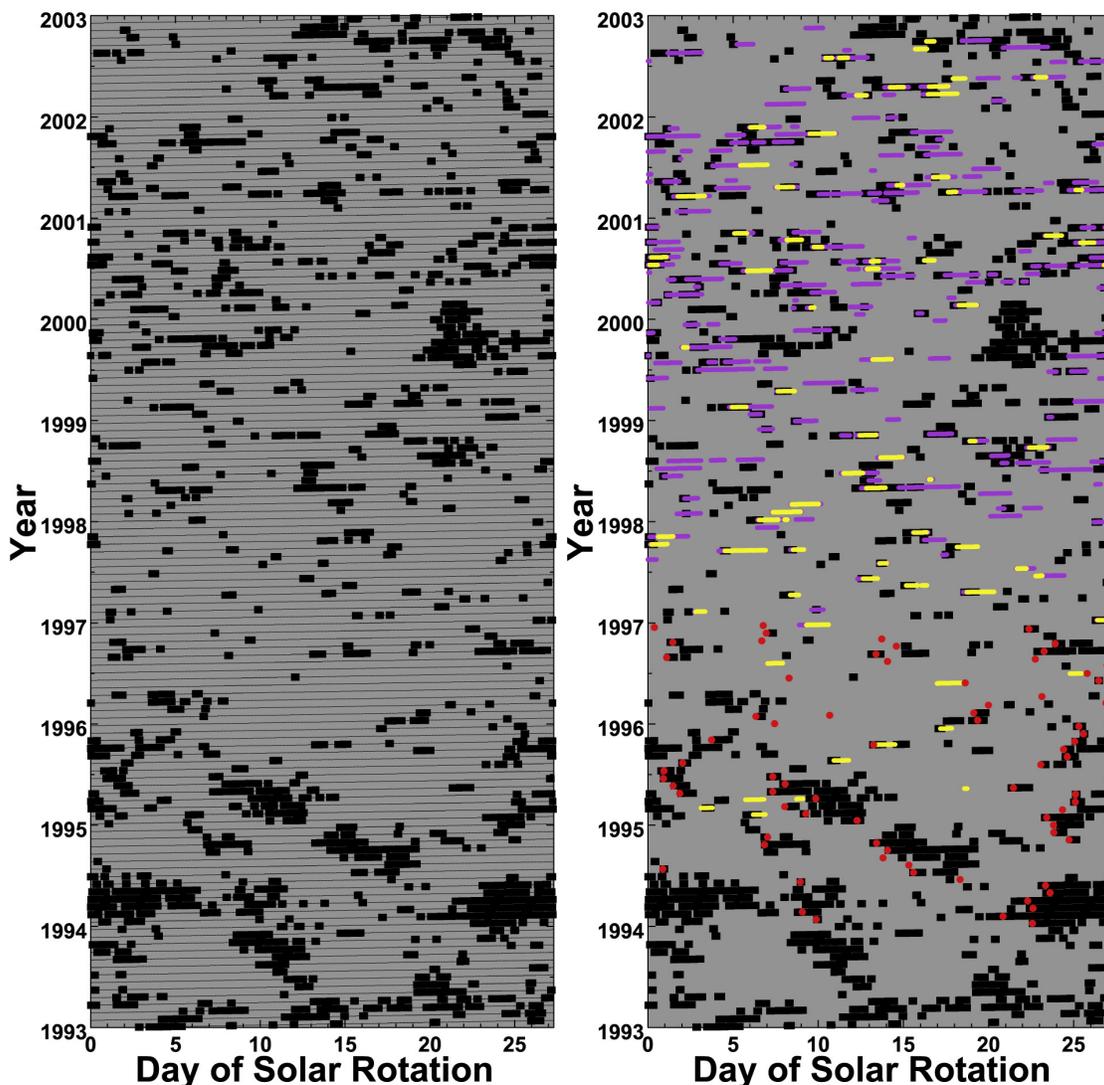


Figure 1. Times of high geomagnetic activity ($K_p \geq 4^+$) are plotted as the black points on a solar rotation versus year plot. In the left panel, the time line is plotted as the thin black lines and the storm times are plotted as the black boxes. In the right panel the storm times are plotted in black, the *Cane and Richardson* [2003] CME catalog is plotted in purple, the Lepping online magnetic cloud catalog is plotted in yellow, and the McPherron CIR-stream-interface catalog is plotted in red.

2.3. Calm Before the Storm

[10] *Borovsky and Steinberg* [2006] found that most recurring high-speed-stream-driven storms have an extended interval (1–2 days) of extreme geomagnetic calm ($K_p \geq 1^+$) within 24 hours prior to storm onset. These calms before the storms are caused by the same Russell-McPherron effect that produces an increased probability of high-speed-stream-driven-storm occurrence [cf. *Crooker and Cliver*, 1994; *Mursula and Zeiger*, 1996]: owing to a sector reversal just upstream of the CIR stream interface, when the IMF in the high-speed wind is favorable for a storm, the IMF in the slow wind is favorable for a calm. A survey of 83 high-speed-stream-driven storms by *Borovsky and Steinberg* [2006] found that 67% of the storms were immediately preceded by extended intervals of geomagnetic calm. Using the same analysis techniques, a survey by the present authors of 68 CME-driven storms (from the *Denton et al.* [2006]

catalog) finds that only 37% are immediately preceded by extended intervals of calm. Hence calms before the storms are predominantly a CIR-driven-storm phenomena. Calms before storms are of note because during the calm intervals the magnetosphere fills with dense plasmaspheric plasma and with dense LLBL (low-latitude boundary layer) plasma. The buildup of dense plasma in the magnetosphere can change the nature of a storm, yielding a stronger ring current (B. Lavraud et al., Magnetosphere preconditioning under northward IMF: Evidence from the study of CME and CIR geoeffectiveness, submitted to *Journal of Geophysical Research*, 2006), initially lower fluxes in the outer radiation belts [*Borovsky and Steinberg*, 2006], and perhaps changes to the reconnection rate at the dayside neutral line [*Borovsky and Steinberg*, 2006]. CIR-driven storms are more often preconditioned by calms than are CME-driven storms.

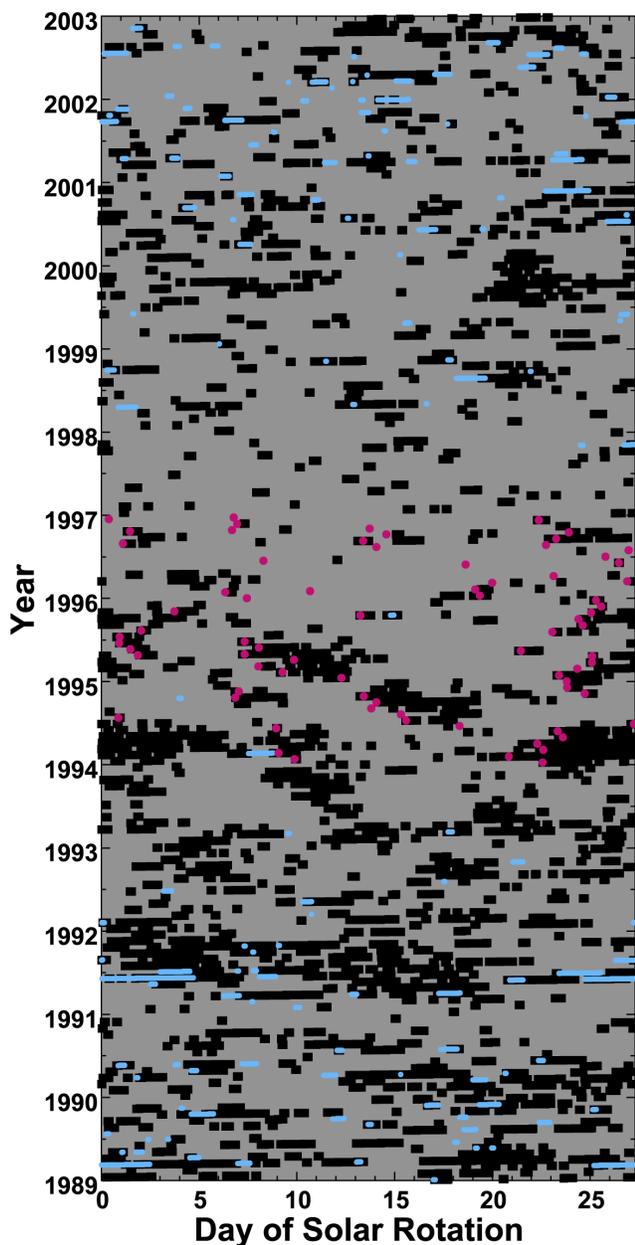


Figure 2. Times of high geomagnetic activity ($K_p \geq 4^+$) are plotted as the black points on a solar rotation versus year plot. The NOAA online SEP catalog is plotted in blue and the McPherron catalog of CIR stream interfaces for the years 1994–1996 is plotted in pink.

2.4. Solar Energetic-Particle (SEP) Events

[11] Solar energetic-particle events are enhanced fluxes of subrelativistic and relativistic ions that have durations of hours to days. SEP events are well known to be associated with solar flares and with strong interplanetary shocks driven by ejecta from the sun [Reames, 2003]. Flares and strong interplanetary shocks are both phenomena that are associated with CMEs, hence SEP events sometimes accompany CME-driven storms. The intensity of SEP events is correlated with the velocity of the coronal ejecta driving the shocks [Kahler, 2001]. On rare occasions, SEP events can be large enough to pose a hazard to aircraft passengers

and electronics [Dyer *et al.*, 2003] and to astronauts in high-latitude orbits [Reames, 1999a]. Energetic-particle events are also produced by CIRs (by CIR shocks in the heliosphere beyond ~ 2 AU, with the energetic particles traveling back to 1 AU) [Reames, 1999b], but their intensity at Earth is weak compared with CME-related SEP events [Mason and Sanderson, 1999], extending only to ~ 20 MeV/n. In Figure 2 it is demonstrated that major SEP events often accompany CME-driven storms and are not associated with CIR-driven storms. Here, for the years 1989–2003, storms (black boxes) are displayed in a solar-rotation versus time plot. The NOAA online SEP-event catalog (<http://www.sec.noaa.gov/ftpdir/indices/SPE.txt>) is plotted as the blue points. These events have integral fluxes of >10 MeV protons at geosynchronous orbit that exceed 10 protons/cm²/s/ster. As can be seen, there are large numbers of SEP events among the storms during the two solar maxima in the figure (the years 1989–1992 and 2000–2002) but there are no SEP events associated with the 27-day-recurring storm groups in the years 1993–1995 nor in the 27-day-recurring storm group in late 1999/early 2000. Also plotted in pink in Figure 2 is the McPherron catalog of CIR stream interfaces for the years 1994–1996. As can be seen, there is at most one SEP event that is associated with one of the 74 CIRs in the catalog, and this SEP is flare associated.

2.5. Storm Sudden Commencement (SSC)

[12] Most storm sudden commencements are associated with strong interplanetary shocks that compress the magnetosphere [Lucci *et al.*, 1988; Russell *et al.*, 1992]. The strong shocks are usually found ahead of fast CMEs. As pointed out by Kamide *et al.* [1998], CIR-driven storms generally lack sudden commencements [see also Venkatesan and Zhu, 1991]. An examination of storms by Taylor *et al.* [1994] supports the conclusion that SSCs tend to be driven by CMEs, with more gradual onsets being driven by high-speed streams. In Figure 3, SSCs (white circles) and storms (black boxes) are displayed on a solar rotation versus time plot. The SSCs are from the NOAA online SSC catalog (<http://www.ngdc.noaa.gov/stp/SOLAR/ftpSSC.html>). As can be seen in the figure, SSCs tend to not be associated with the groups of 27-day-recurring storms (cf. the years 1993, 1994, 1995, and 1999). Also, as can be seen in the figure, SSCs are associated with the randomly occurring storms. To quantify these observations, the Borovsky and Steinberg [2006] catalog of high-speed-stream-driven storms is examined and it is found that 21% of those storms are temporally associated with SSCs from the NOAA SSC catalog. Forming a CME-driven-storm catalog by requiring $K_p \geq 4^+$ in the Cane and Richardson [2003] CME catalog and the Lepping online [Lepping *et al.*, 2005] magnetic cloud catalog, it is found that 69% of those CME-driven storms are temporally associated with SSCs from the NOAA SSC catalog. Hence SSCs are common ($\sim 70\%$) for CME-driven storms and infrequent ($\sim 20\%$) for CIR-driven storms.

2.6. Magnetosonic Mach Number of the Solar Wind

[13] Depending on the values of the Alfvén velocity $v_A = B/(4\pi n m_i)^{1/2}$ and ion-acoustic speed $C_s = (k_B(T_i + T_e)/m_i)^{1/2}$ in the solar wind, the magnetosonic Mach number $M_{ms} = v_{sw}/(v_A^2 + 5C_s^2/3)^{1/2}$ of the solar wind flow relative to the

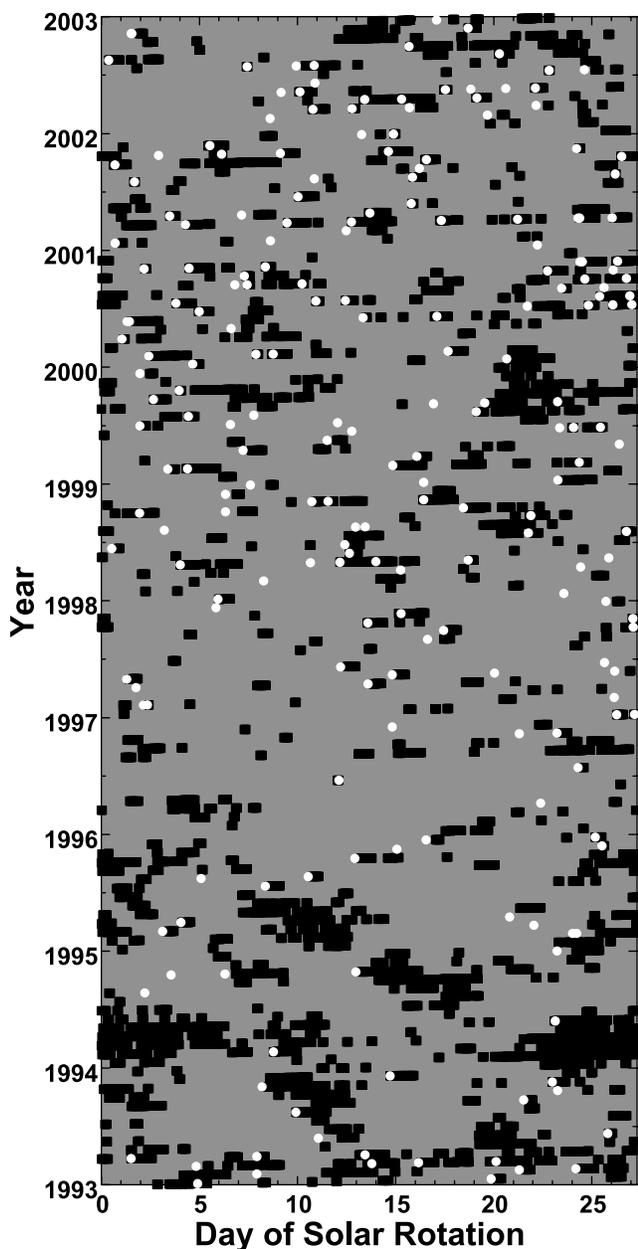


Figure 3. Times of high geomagnetic activity ($K_p \geq 4^+$) are plotted as the black points on a solar rotation versus year plot. Times of occurrence of SSCs (storm sudden commencements) at Earth are indicated which white circles. The SSCs occurrence times are from the NOAA online SSC catalog (<http://www.ngdc.noaa.gov/stp/SOLAR/ftpSSC.html>).

Earth can change. The magnetosonic Mach number M_{ms} sets the properties of the bow shock, which converts unshocked solar wind plasma into shocked magnetosheath plasma. The magnetosonic Mach number controls the compression ratio C of the shock, which is the multiplicative density change from upstream to downstream. The compression ratio C varies from 1 (at $M_{ms} = 1$) to 4 (at $M_{ms} \rightarrow \infty$) [Tidman and Krall, 1971]. In solar wind/magnetosphere coupling, it is actually the magnetosheath flow that drives activity in the Earth's magnetosphere, and

the bow shock converts solar wind flow into magnetosheath flow. Typically, the Mach number of the solar wind is high and the compression ratio is ~ 4 . As noted by Lopez *et al.* [2004], during some storms the Mach number of the solar wind can be low and the compression ratio of the bow shock can be substantially less than 4; in such cases the solar wind/magnetosphere coupling has a different parametric dependence on the solar wind than it does in the much more familiar high Mach number case. In the top panel of Figure 4 the magnetosonic Mach number M_{ms} of the solar wind is binned under various conditions. The electron temperature T_e of the solar wind is not measured regularly

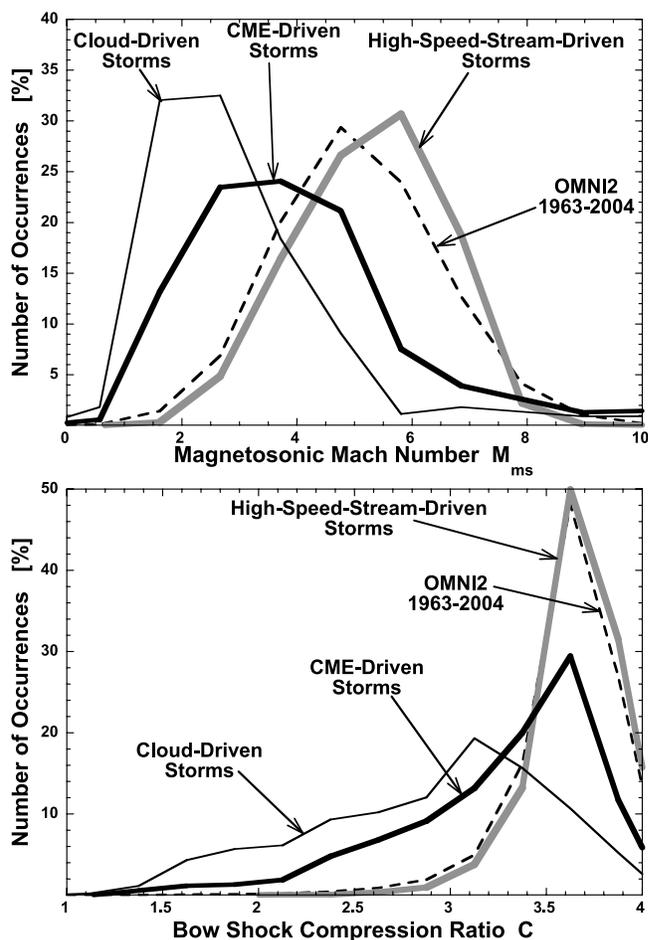


Figure 4. (top) For various data sets (as labeled) the magnetosonic Mach number M_{ms} of the solar wind is binned. The Mach number is given by $M_{ms} = v_{sw}/(v_A^2 + 5C_s^2/3)^{1/2}$, where $v_A = B^2/(4\pi n m_i)^{1/2}$ and $C_s = (k_B(T_e + T_i)/m_i)^{1/2}$ are the Alfvén and ion-acoustic speeds in the upstream solar wind. The median values of M_{ms} are 5.0 for the 1963–2004 OMNI2 data set, 5.3 for high-speed-stream-driven storms, 3.7 for CME-driven storms, and 2.7 for the magnetic cloud portions of CME-driven storms. (bottom) For various data sets (as labeled) the estimated compression ratio C of the Earth's bow shock is binned. The median values of C are 3.65 for the 1963–2004 OMNI2 data set, 3.69 for high-speed-stream-driven storms, 3.40 for CME-driven storms, and 3.01 for the magnetic-cloud portions of CME-driven storms.

and hence is not available in the OMNI data sets, so to calculate the ion acoustic speed $C_s = (k_B(T_i + T_e)/m_i)^{1/2}$ of the solar wind, $T_e = 1.5 \times 10^5 \text{ K} = 12.9 \text{ eV}$ is taken, which is approximately valid for all types of solar wind [Skoug *et al.*, 2000]. All other solar wind parameters are obtained from the OMNI2 solar wind data set [King and Papitashvili, 2005]. The dashed curve in Figure 4 is the occurrence distribution of magnetosonic Mach numbers in the entire 1963–2004 OMNI2 data set. The thick gray curve is the occurrence distribution for times when high-speed-stream-driven storms are ongoing. The collection of high-speed-stream-driven-storm times is the Borovsky and Steinberg [2006] catalog of storms driven by 27-day recurring high-speed streams. As can be seen, the Mach numbers during high-speed-stream-driven storms have a typical distribution of values (i.e., similar to OMNI2). The thick black curve in Figure 4 is for times when CME-driven storms are ongoing and the thin black curve is for times when the magnetic clouds are passing during CME-driven storms. The collection of CME-driven-storm times is obtained by requiring $K_p \geq 4^+$ in the Cane and Richardson [2003] CME catalog and the collection of cloud-driven storm times is obtained by requiring $K_p \geq 4^+$ in the Lepping online magnetic cloud catalog [cf. Lepping *et al.*, 2005]. As can be seen, the Mach numbers of the solar wind are anomalously low for CME-driven storms and even lower for the cloud-driven storms. In the bottom panel of Figure 4, the compression ratio C across the bow shock is binned for the same conditions as in the top panel. The compression ratio for the bow shock is estimated using the Rankine-Hugoniot shock jump conditions (equations (1.50) and (1.51) of Tidman and Krall [1971]) under the assumption of a quasi-perpendicular shock. The dashed curve in the bottom panel is the occurrence distribution of C for the entire 1963–2004 OMNI2 data set. As can be seen, the compression ratio C is typically near 4. The thick gray curve in the bottom panel is the distribution of C values for times when high-speed-stream-driven storms are ongoing; as can be seen C is typically near 4 for high-speed-stream-driven storms. The thick black curve in the bottom panel is for times when CME-driven storms are ongoing and the thin black curve is for the magnetic cloud portions of CME-driven storms. As can be seen, the bow shock compression ratio for CME-driven storms, and particularly for cloud-driven storms, is not near 4. For cloud-driven storms, 26.8% of the time C is less than 2.5, which is the midway point between the limits of 1 and 4. Because of the systematic difference in the bow shock compression ratio, it is expected that the dependence of geomagnetic activity on solar wind parameters will differ for CME-driven and high-speed-stream-driven storms.

2.7. Plasma- β of the Magnetosheath

[14] For the solar wind parameters that lead to global sawtooth oscillations of the magnetosphere, it has been noted that the plasma- β value of the magnetosheath is anomalously low [Borovsky, 2004; J. E. Borovsky *et al.*, The solar-wind driving of global sawtooth oscillations and periodic substorms: What determines the periodicity?, submitted to *Annales Geophysicae*, 2006, hereinafter referred to as Borovsky *et al.*, submitted manuscript, 2006], often less than unity. The plasma- β is a measure of how strongly distorted the magnetic field can be owing to diamagnetic

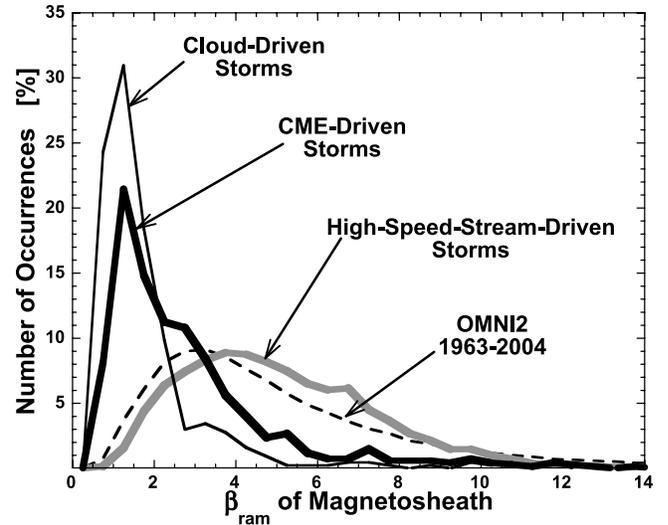


Figure 5. For various data sets (as labeled) the plasma- β_{ram} value of the magnetosheath estimated from the Rankine-Hugoniot shock-jump conditions and upstream solar wind parameters is binned. The median values of β_{ram} for the various distributions are 4.42 for the 1963–2004 OMNI2 data set, 4.75 for high-speed-stream-driven storms, 2.29 for storms driven by CMEs, and 1.43 for the magnetic-cloud portions of CME-driven storms.

currents in the plasma. β can be defined two ways: a thermal β defined as $\beta_{\text{th}} = 8\pi n k_B(T_i + T_e)/B^2$ or a ram β defined as $\beta_{\text{ram}} = 4\pi n m_i v^2/B^2$ where v is the plasma flow velocity [Parsons and Jellison, 1983; Borovsky, 1992]. β_{th} is a measure of the currents when a warm plasma is confined by a magnetic field; β_{ram} is a measure of the currents when a plasma flow is deflected or slowed. Although $\beta_{\text{th}} \approx \beta_{\text{ram}}$ in the magnetosheath plasma, we are specifically interested in the value of β_{ram} , which is the ratio of the ram pressure $n m_i v^2$ of a flow to the magnetic-field pressure $B^2/8\pi$ of the plasma. (Note also that $\beta_{\text{ram}} = M_A^2$, where M_A is here the Alfvén Mach number of the magnetosheath flow.) If $\beta_{\text{ram}} < 1$, a flow has difficulty distorting magnetic-field lines. Ordinarily, in the magnetosheath $\beta_{\text{ram}} \gg 1$ and the flow of the magnetosheath is little impeded by the magnetic field in the magnetosheath. In that case a gasdynamic flow picture [e.g., Alksne and Webster, 1970; Spreiter and Stahara, 1994; Stahara, 2002] is valid with the magnetic field lines passively convected by the magnetosheath flow; strong draping of the field lines over the magnetosphere results. Global MHD simulations of the solar-wind-driven magnetosphere when $\beta_{\text{ram}} < 1$ in the magnetosheath [Borovsky *et al.*, 2004] find that the magnetosheath magnetic field lines are stiff and do not drape; the stiff field lines squeeze the magnetosphere into an unusual flattened shape and produce an asymmetric magnetosheath flow pattern that is very different from the gasdynamic picture. The β_{ram} of the magnetosheath flow can be estimated from upstream solar wind parameters by using the Rankine-Hugoniot quasi-perpendicular shock jump conditions. Taking $T_e = 12.9 \text{ eV}$ for the solar wind (see section 2.6) along with measured values of v , T_i , B , and n in the solar wind, equations (1.50) and (1.51) of Tidman and Krall [1971] yield β_{ram} values for the magnetosheath plasma. In Figure 5

these β_{ram} values are binned for various conditions. The dashed curve is the occurrence distribution of β_{ram} for the entire 1963–2004 OMNI2 data set. As can be seen from the OMNI2 curve, β_{ram} in the magnetosheath is typically high. The thick gray curve is for times when high-speed-stream-driven storms are ongoing, the thick black curve is for times when CME-driven storms are ongoing, and the thin black curve is for times when magnetic clouds are passing the Earth during storms. As can be seen, for high-speed-stream-driven storms β_{ram} is slightly higher than typical. The β_{ram} values for CME-driven storms are anomalously low and the β_{ram} values for the magnetic cloud portions of CME-driven storms are even lower. For the OMNI2 data set, 0.75% of the time $\beta_{\text{ram}} < 1$ in the magnetosheath. For high-speed-stream-driven storms, only 0.17% of the time is $\beta_{\text{ram}} < 1$, whereas for CME-driven storms 8.14% of the time $\beta_{\text{ram}} < 1$ and for cloud-driven storms 24.1% of the time $\beta_{\text{ram}} < 1$.

2.8. Plasma Sheet Density

[15] The density of the Earth's plasma sheet is governed by the solar wind density with a few hour time lag [Borovsky *et al.*, 1998a]. In a superposed epoch study, Denton *et al.* [2006] compared the Earth's plasma sheet density for CME-driven storms with its density for CIR-driven storms as functions of time during the storm and local time around geosynchronous orbit. Typical nonstorm nightside densities are $\sim 0.7 \text{ cm}^{-3}$. For both CME-driven and CIR-driven storms, Denton *et al.* found that superdense plasma sheets [cf. Borovsky *et al.*, 1997] occur, with the density being twice or more the typical value. In comparing the two drivers, CMEs drive a plasma sheet that is more dense and the high-density values persist longer. The relatively brief superdense intervals in CIR-driven storms appear to coincide with the CIR, which has a high solar wind density; this is in contrast to the enhanced geomagnetic activity, which persists long into the high-speed stream that follows the CIR. This is because high-speed streams have only modest densities although they have strong vB_z . In a similar fashion, the superdense intervals of CME-driven storms probably coincide with the passing of the higher-density CME sheath. Note that CME-driven storms without a superdense plasma sheet can occur [e.g., Thomsen *et al.*, 1998a; Seki *et al.*, 2005].

2.9. Plasma Sheet Temperature

[16] In a superposed epoch study, Denton *et al.* [2006] compared the Earth's plasma sheet ion and electron temperatures for CME-driven storms with CIR-driven storms, making the examination as functions of time during the storm and local time of the spacecraft measurement. For both types of storms the ion and electron temperatures are substantially elevated over typical values, probably because the solar wind speed is high and the plasma sheet temperature is related to the solar wind velocity [Borovsky *et al.*, 1998a]. Comparing the plasma sheet temperature for the two types of storms (Figure 4 of Denton *et al.*), the plasma sheet is hotter for the CIR-driven storms than it is for the CME-driven storms and the elevated temperatures persist longer (days) for CIR-driven storms. This is particularly evident when the electron temperatures are examined.

2.10. Plasma Sheet Composition

[17] In the superposed epoch study, Denton *et al.* [2006] compared the estimated O^+/H^+ ratio in the Earth's plasma sheet for CME-driven storms with CIR-driven storms as functions of time during the storm and local time around geosynchronous orbit. Both types of storms produced elevated ratios, but CMEs that drive large Dst storms producing very elevated O^+/H^+ ratios (exceeding unity). Hence CME-driven storms can have very strong ionospheric outflows of O^+ and CME-driven storms can produce plasma sheets that are predominantly O^+ by number density.

2.11. Spacecraft Surface Charging

[18] In the superposed epoch study, Denton *et al.* [2006] compared the measured voltages of geosynchronous-orbit spacecraft with respect to the ambient plasma for CME-driven storms with CIR-driven storms. The comparison is made as functions of time during the storm and local time of the spacecraft. For both types of storms the hot plasmas of the magnetosphere drive substantial spacecraft charging, even outside of the Earth's shadow. The spacecraft charging that occurs during CIR-driven storms is more severe than it is for CME-driven storms: the voltages attained are higher (greater than -1000 volts), the region of local time over which the severe charging occurs is wider across the nightside of the magnetosphere, and the high voltages persist longer (days). This result is consistent with the well-known fact that spacecraft charging is driven by hot electrons [Whipple, 1981; Mullen *et al.*, 1986], and, as discussed in section 2.5 above, the electron temperatures of the magnetosphere are hotter and more persistent for CIR-driven storms than they are for CME-driven storms. Note that voltages of differential surface charging of spacecraft are proportional to the voltages of surface charging [cf. Borovsky *et al.*, 1998b, Figure 10], so differential surface charging is anticipated to be more severe during CIR-driven storms than it is during CME-driven storms.

2.12. Ring Current and Dst

[19] In a survey of Dst perturbations, Yermolaev and Yermolaev [2002] found that CMEs are more effective at producing strong Dst values than are CIRs and high-speed streams. This conclusion is borne out by the Denton *et al.* [2006] survey that looked at the Earth's reaction to CMEs and CIRs in the solar wind, which also found that CMEs produced larger Dst perturbations. This conclusion is also borne out by Figure 6, where storms and Dst perturbations are displayed on a solar rotation versus time plot. Times when storms are ongoing are indicated by black boxes and times when $\text{Dst} < -100$ are indicated by white circles. As can be seen, most of the strong Dst perturbations are associated with randomly occurring storms at solar maximum (CME-driven storms). Strong Dst perturbations are rare in the 27-day-recurring storms groups (high-speed-stream-driven storms). Dst is a measure of the strength of the ring current and partial ring current in the magnetosphere [Liemohn *et al.*, 2001; Jordanova *et al.*, 2003]. Since a stronger ring current is produced by a plasma sheet that is denser and cooler and magnetospheric convection that is stronger [Thomsen *et al.*, 1998b; Kozyra and Liemohn, 2003], all of which are more true for CMEs than CIRs, it

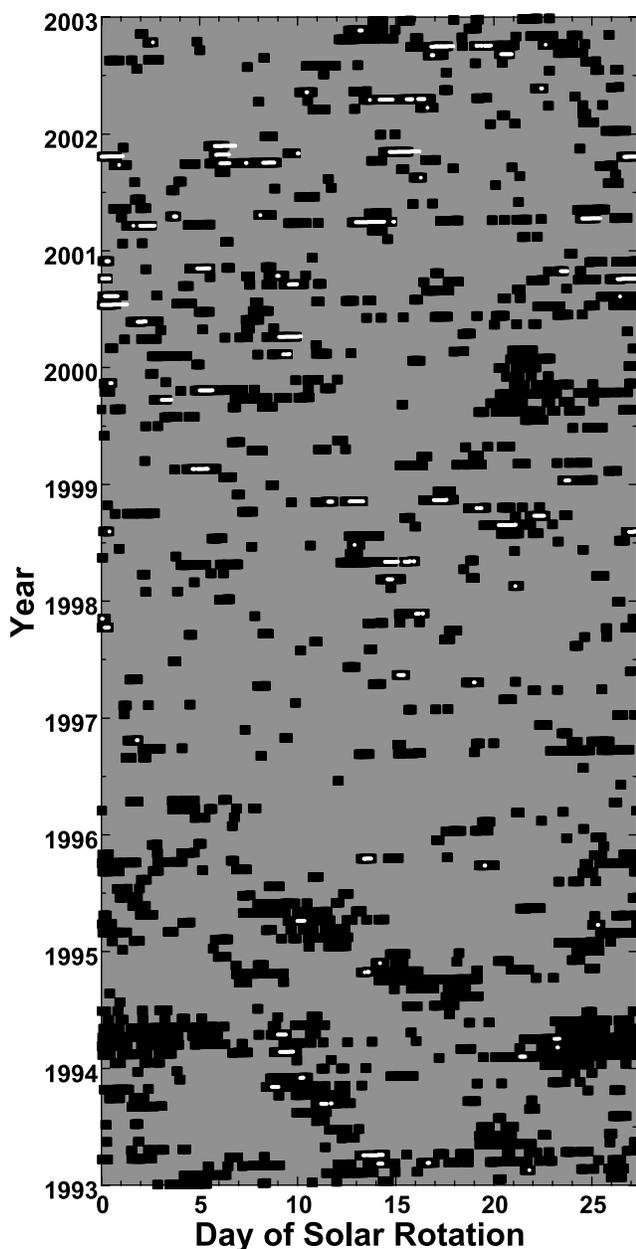


Figure 6. Times of high geomagnetic activity ($K_p \geq 4^+$) are plotted as the black boxes on a solar rotation versus year plot. Times when $Dst < -100$ are plotted as white points.

makes sense that Dst is stronger for CME-driven storms than it is for CIR-driven storms.

2.13. Global Sawtooth Oscillations

[20] Global sawtooth oscillations are periodic (3–4 hours) oscillations of the magnetosphere that are characterized by a slow global distortion of the magnetosphere followed by a rapid crash back to a less-distorted morphology [Borovsky *et al.*, 2001]. They are observed to be driven by magnetic clouds in CMEs that have southward magnetic field in the range of -5 to -15 nT and that tend to have parameters that lead to polar cap saturation (Borovsky *et al.*, submitted manuscript, 2006). This is low Mach number solar wind and for these parameters the Earth's magnetosheath is a $\beta < 1$ plasma,

which has very unusual flow properties (see section 2.7). Why these solar wind conditions lead to sawtooth oscillations of the magnetosphere during storms is presently an area of active research. Global sawtooth oscillations are not seen during CIR-driven storms (Borovsky *et al.*, submitted manuscript, 2006).

2.14. ULF Pulsations

[21] ULF oscillations in the magnetosphere have been reported for CME-driven storms [Baker *et al.*, 1998] and high-speed-stream-driven storms [Kessel *et al.*, 2004], with pulsation activity peaking in the declining phase of the solar cycle [Zieger, 1991]. In Figure 7 the power of ULF pulsations P_{ULF} in the dawnside magnetosphere is binned for various conditions. The data set used is the 1987–2000 daily average of the dawnside power obtained from the IMAGE and SAMNET ground-based magnetometer networks, courtesy of P. O'Brien (private communication, 2005); a description of the data set is found in the work of O'Brien *et al.* [2003]. The dashed curve is the occurrence distribution of $\text{Log}_{10}(P_{ULF})$ for the entire 1987–2000 P_{ULF} data set. The thick gray curve is the occurrence distribution of $\text{Log}_{10}(P_{ULF})$ during high-speed-stream-driven storms, the thick black curve is the occurrence distribution of $\text{Log}_{10}(P_{ULF})$ during CME-driven storms, and the thin black curve is the occurrence distribution of $\text{Log}_{10}(P_{ULF})$ during the magnetic cloud portions of CME-driven storms. As can be seen in Figure 7, P_{ULF} is elevated above normal for all types of storms. The amplitude of Pc5 oscillations in the magnetosphere is proportional to the solar wind velocity [Engebretson *et al.*, 1998; Mathie and Mann, 2001]. Electron acceleration by Pc5 oscillations are believed to be responsible for the production of high fluxes of relativistic electrons in the outer radiation belts [Rostoker *et al.*, 1998], particularly during long intervals of active ULF pulsations [Mathie and Mann, 2000; O'Brien *et al.*, 2001], which

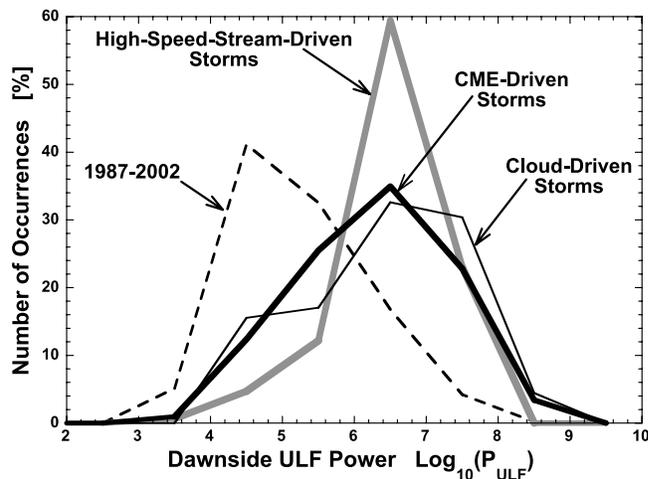


Figure 7. For various data sets (as labeled) the logarithm of the dawnside ULF power $\text{Log}_{10}(P_{ULF})$, where P_{ULF} is in units of nT^2/Hz , is binned. The median values of $\text{Log}_{10}(P_{ULF})$ for the various distributions are 5.09 for the 1963–2004 OMNI2 data set, 6.57 for high-speed-stream-driven storms, 6.10 for storms driven by CMEs, and 6.56 for the magnetic cloud portions of CME-driven storms.

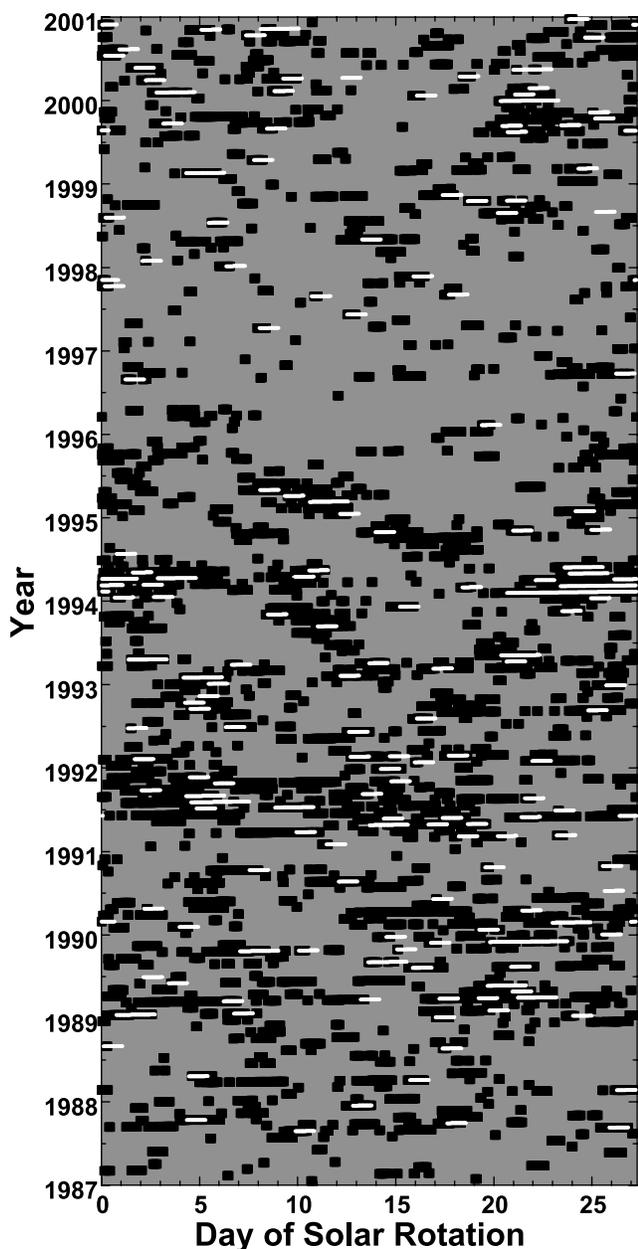


Figure 8. Times of high geomagnetic activity ($K_p \geq 4^+$) are plotted as the black boxes on a solar rotation versus year plot. Times when the dawnside ULF power P_{ULF} exceeds 1×10^7 nT²/Hz are plotted as white circles. The ULF power was provided by Paul O'Brien [O'Brien *et al.*, 2003].

occur during longer intervals of high-velocity solar wind associated with recurrent high-speed streams. This is shown in Figure 8, where storms and intervals of high ULF power are shown on a solar rotation versus time plot for the years 1987–2000. The storms and intervals when $P_{ULF} > 1 \times 10^7$ nT²/Hz are indicated as white circles. As can be seen, the longest (multiday) intervals of elevated ULF power are found with the 27-day recurring storm groups, particularly in the year 1994. The intervals of active ULF pulsations are longer for CIR-driven storms than they are for CME-driven storms because the intervals of elevated solar wind velocity are longer.

Evidence for this can be found in Figure 9. For full years when there is thorough data coverage of the solar wind (the years 1995–present), a measure of the duration of high-speed flows is plotted in Figure 9 as the green bars. This measure is made as follows. For each year of data, the durations of all flows wherein the solar wind velocity is sustained at 550 km/s or higher are calculated. Then, for each year, the 12 longest-duration fast flows are taken. The height of the green bar in Figure 9 is the mean duration for those 12 longest flows that year. The sunspot number is also plotted as the black bars in Figure 9, and intervals when 27-day recurring storms (which are high-speed-stream driven) were ongoing are marked with the horizontal blue bands (see section 2.2 for the methodology of finding the recurring storm groups). As can be seen, when recurring CIR-driven storms were ongoing (which tends to be in the declining phase of the solar cycle), the fast flow durations were longer. Since the power in ULF pulsations is proportional to the solar wind speed, it follows that the duration of intervals of enhanced ULF activity are longer in years when CIR-driven storms are present than for the rest of the solar cycle. This is also shown in Figure 9, where a 0.5-year running average of the measured flux of 1.1–1.5 MeV electrons at geosynchronous orbit is plotted as the red curve. (See section 2.17 for a description of this measurement.) This flux is elevated in the years when the durations of the high-speed flows were long (green bars).

2.15. Distortion of the Dipole

[22] CIR-driven storms produce modest stretching of the dipole magnetic field toward a tail-like geometry (with the near-equatorial field at geosynchronous orbit rotated $\sim 30^\circ$

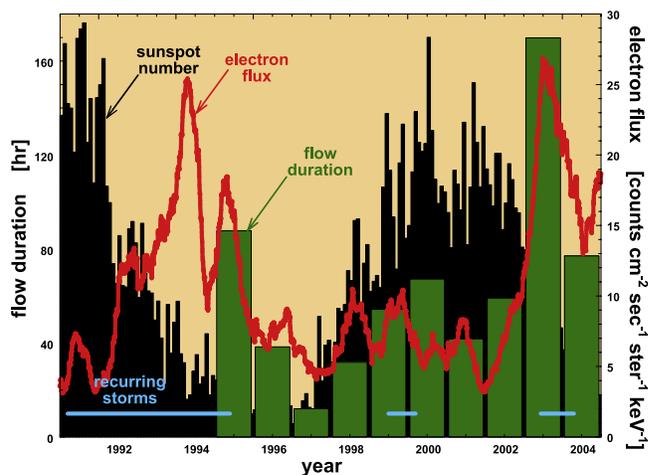


Figure 9. As functions of time, the sunspot number (black bars, left axis), duration of high-speed flows of each year (green bars, left axis), and a 0.5-year running average of the flux of 1.1–1.5 MeV electrons at geosynchronous orbit (red curve, right axis) are plotted. The flow duration plotted is the mean value of the 12 longest-duration fast flow that year, where fast flows are $v \geq 550$ km/s, obtained with the OMNI2 data set [King and Papitashvili, 2005]. The energetic electron fluxes are multispacecraft averages obtained from the SOPA instruments [Reeves *et al.*, 1997]. Time intervals when 27-day recurring storms were ongoing are marked with the blue horizontal bands.

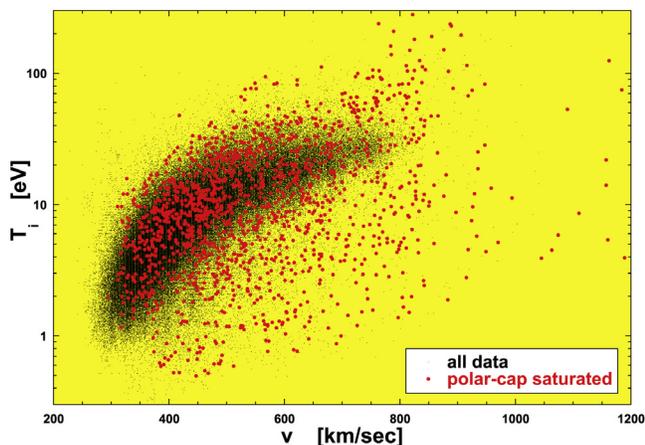


Figure 10. Using the hourly averaged solar wind values in the 1963–2004 OMNI2 data set, the ion temperature of the solar wind is plotted as a function of the velocity of the solar wind (black points). During time intervals when the conditions of the solar wind should result in a saturation of the Earth’s cross-polar-cap potential (expression (2)), the points are replotted in red. As can be seen, the red points form a population of outliers around the main solar wind population.

from dipolar) with the stretching angle proportional to the local pressure of the ion plasma sheet [see *Borovsky et al.*, 1998b, Figure 18]. During times when global sawtooth oscillations occur, which are driven by CMEs, the dipole magnetic field can show fully tail-like stretching at geosynchronous orbit (with the near-equatorial field rotated $\sim 70^\circ$ from dipolar) (*Borovsky et al.*, submitted manuscript, 2006). Two further studies indicate that large distortions of the dipole magnetic field are associated with CME-driven storms. First, the *Tsyganenko et al.* [2003] catalog of big distortion events for the inner magnetosphere all were large Dst events that occurred at solar maximum, which is indicative of driving by CMEs. Second, the flankside lobe encounters by geosynchronous spacecraft at the equator that were tabulated during an interval of solar maximum found that most lobe encounters were associated with large-Dst storms [*Moldwin et al.*, 1995], which is also consistent with driving by CMEs.

2.16. Saturation of Polar Cap Potential

[23] Typically, under vB_s driving by the solar wind, the cross-polar-cap potential φ is linearly proportional to vB_s [*Reiff and Luhmann*, 1986], where v is the solar wind velocity and B_s is the southward (GSM) component of the IMF, but under rare circumstances when the solar wind conditions are favorable, the cross-polar-cap potential φ will reach a saturation value and will no longer be proportional to vB_s . During several CME-driven storms this saturation has been analyzed [*Ober et al.*, 2003; *Boudouridis et al.*, 2004; *Hairston et al.*, 2005]. Equation (1) of *Siscoe et al.* [2004] describes this saturation: saturation occurs when the second term in the denominator of that equation exceeds approximately 2, which occurs when

$$v_A \Sigma_P / 806 > 2, \quad (1)$$

where v_A is the Alfvén velocity in the unshocked solar wind (in units of km/s) and Σ_P is the height-integrated Pederson conductivity of the dayside ionosphere (in units of mho). Taking $\Sigma_P = 0.77 F_{10.7}^{1/2}$ (from equation (4) of *Ober et al.* [2003]), where $F_{10.7}$ is the radio emission of the Sun at 10.7 cm which is a proxy for solar EUV emission which produces ionospheric ionization [cf. *Balan et al.*, 1993], and $F_{10.7} \approx S_N$ (from Figure 1 and Table 2 of *Floyd et al.* [2005]), where S_N is the international sunspot number, expression (1) is rewritten as

$$Q \equiv v_A S_N^{1/2} / 1050 > 2, \quad (2)$$

where Q is a saturation parameter. The OMNI2 data set is used to find times when expression (2) is satisfied: the results are displayed in Figures 10 and 11. In Figure 10 the temperature-speed plot of the solar wind using the OMNI2 data set [*King and Papitashvili*, 2005] is shown (black points) and the data points wherein expression (2) is satisfied are plotted in red. Expression (2) is satisfied 0.96% of the time in the 1963–2004 OMNI2 data set. These red points where the solar wind will drive the polar cap voltage into saturation are outliers to the normal solar wind population; this is consistent with the red points being CMEs [see, e.g., *Elliott et al.*, 2005, Figure 1a]. In Figure 11 the sunspot number S_N is plotted as a function of time as the black points, with the values at times when expression (2) is satisfied plotted in red. As can be seen, the time intervals when expression (2) is satisfied occur during solar maxima; this is again consistent with those times being CMEs. A detailed examination of ACE solar-wind data for time intervals when expression (2) is satisfied (*J. Steinberg*, private communication, 2005) finds that most of those intervals are in CMEs, typically CMEs with low number density. A minor fraction of the points in the OMNI2 data set satisfying expression (2) are low-density pressure balance structures embedded in the solar wind. Polar cap

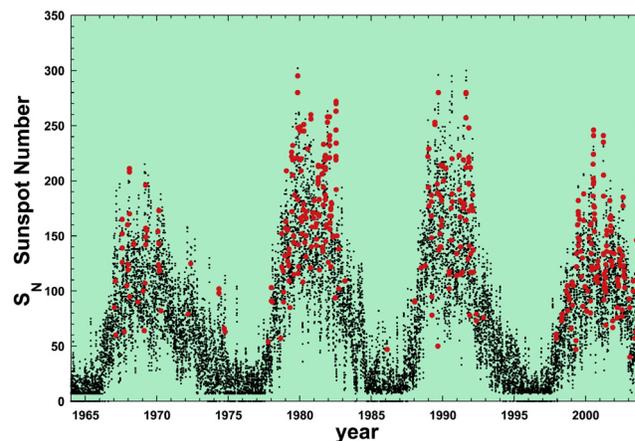
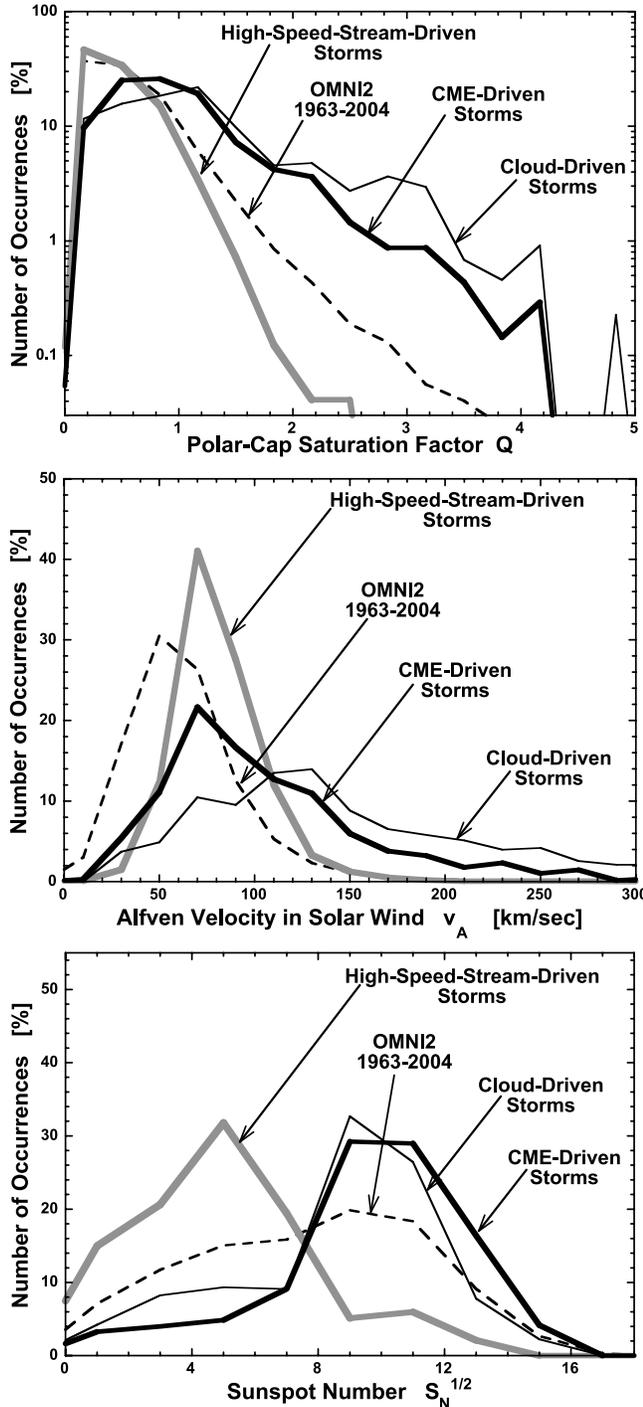


Figure 11. The sunspot number S_N is plotted as a function of time for 1963–2004 (black points). Using the OMNI2 data set, time intervals in which the Earth’s polar-cap voltage should saturate are calculated (expression (2)) and the sunspot number for those time intervals are replotted in red. As can be seen, those times occur mainly during solar maximum.

saturation is usually thought of as a phenomenon associated with strong storms, but note that according to expression (2) and Figure 10 the polar cap potential will saturate for drivers that are not extremely strong (i.e., it will saturate for slow CMEs also). To further explore the polar cap saturation for CME- and CIR-driven storms, the value of Q is binned under various circumstances in the top panel of Figure 12. The dashed curve is the occurrence distribution of Q for the entire 1963–2004 OMNI2 data set, the thick gray curve is the occurrence distribution of Q for high-speed-stream-driven storms, the thick black curve is the occurrence distribution of Q for CME-driven storms, and the thin black



curve is the occurrence distribution of Q for the magnetic cloud portions of CME-driven storms. As can be seen, saturation of the polar cap ($Q > 2$) occurs rarely for high-speed-stream-driven storms but commonly for CME-driven storms. Even considering partial saturation, for high-speed-stream-driven storms only 4.4% of the time is $Q > 1$, whereas for CME-driven storms 39.3% of the time $Q > 1$ and for cloud-driven storms 54.2% of the time $Q > 1$. To explore whether the polar cap saturation is simply owed to the high values of Σ_P during solar maximum when CMEs tend to occur, the Alfven velocity v_A in the solar wind and the square root of the sunspot number $S_N^{1/2}$ are binned in the middle and bottom panels of Figure 12. Examining the various curves in the middle and bottom panels it can be concluded that the higher values of $Q = v_A \Sigma_P / 806 = v_A S_N^{1/2} / 1050$ for CME-driven storms are owed both to higher values of v_A in the solar wind and to higher values of $S_N^{1/2}$.

2.17. Fluxes of Relativistic Electrons

[24] The Earth's outer electron radiation belt is very dynamic. It is well known that the fluxes of relativistic electrons in the outer radiation belt are maximized by recurring high-speed-stream-driven storms during the declining phase of the solar cycle [Paulikas and Blake, 1976; Love et al., 2000; Lam, 2004]. Sufficiently high fluxes of electrons in this belt lead to internal charging in spacecraft, which leads to operational problems for spacecraft in geosynchronous orbit [Romanova et al., 2005]. The frequency of these problems peaks during the declining phase of the solar cycle where CIR-driven storms peak [Wrenn et al., 2002]. In Figure 13 the production of intense fluxes of relativistic electrons in the outer radiation belts by the two types of storms is explored. For the years 1993–2001 geomagnetic storms are indicated (black boxes) on a solar rotation versus time plot. The 27-day-recurring storm groups (CIR-driven storms) are clearly seen in 1994 and 1995 and in late 1999/early 2000. The McPherron catalog of CIR stream interfaces for the years 1994–1996 is plotted in yellow. In Figure 13, times when the multispacecraft daily average of 1.1–1.5 MeV electrons at geosynchronous orbit as measured by the SOPA energetic particle instruments [Belian et al., 1992] exceeds 30 electrons/cm²/s/ster/keV are indicated in red. The relativistic-electron fluxes

Figure 12. In the top panel the polar-cap-saturation factor $Q = v_A S_N^{1/2} / 1050$ is binned or various data sets (as labeled). Saturation of the polar cap occurs for $Q > 2$. The median values of Q for the various distributions are 0.45 for the 1963–2004 OMNI2 data set, 0.36 for high-speed-stream-driven storms, 0.87 for CME-driven storms, and 1.06 for cloud-driven storms. In the middle panel the Alfven velocity v_A in the upstream solar wind is binned. The median values of v_A for the various distributions are 60 km/s for the 1963–2004 OMNI2 data set, 78 km/s for high-speed-stream-driven storms, 95 km/s for CME-driven storms, and 131 km/s for cloud-driven storms. In the bottom panel the square root of the sunspot number $S_N^{1/2}$ is plotted. The median values of $S_N^{1/2}$ for the various distributions are 8.0 for the 1963–2004 OMNI2 data set, 4.8 for high-speed-stream-driven storms, 9.9 for CME-driven storms, and 8.7 for cloud-driven storms.

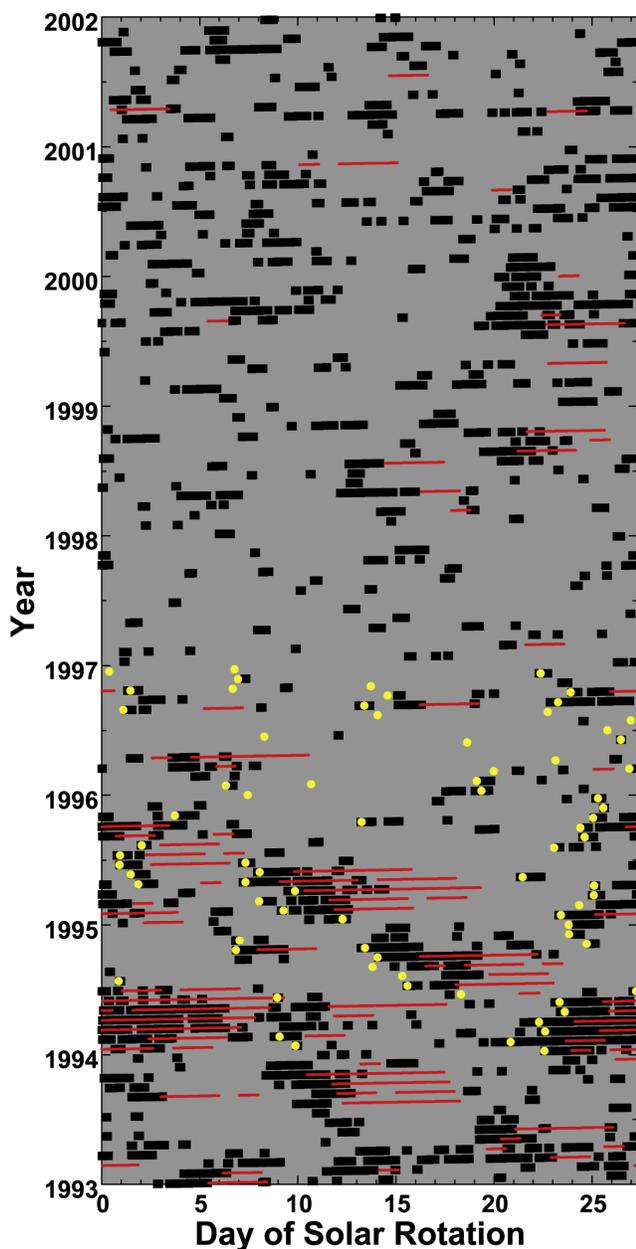


Figure 13. Times of high geomagnetic activity ($K_p \geq 4^+$) are plotted as the black points on a solar-rotation versus year plot. Times when the daily flux of 1.1–1.5 MeV electrons at geosynchronous orbit exceeds 30 counts/cm²/s/ster/keV are indicated in red and the McPherron catalog of CIR stream interfaces for the years 1994–1996 is plotted in yellow. The relativistic electron flux was provided by Reiner Friedel.

were provided by R. Friedel (private communication, 2005). As can be seen in the figure, the high fluxes of relativistic electrons are predominantly associated with CIR-driven storms. There are some relativistic electron events that are associated with randomly occurring storms, e.g., in the years 1996–2002, but a comparison of these events with the Cane and Richardson CME catalog and the Lepping magnetic cloud catalog finds that only one event is clearly CME related. It may be that the nonrecurrent relativistic electron flux events are associated with IRs (i.e., nonrecur-

rent high-speed-stream-driven storms). In Figure 13, note the well-known delay in the appearance of intense relativistic electron fluxes from the onsets of the storms [cf. Friedel *et al.*, 2002].

2.18. Formation of New Ion and Electron Radiation Belts

[25] The production of a new ion radiation belt is known to be caused by the capture of SEP ions during a strong shock compression of the magnetosphere [Lorentzen *et al.*, 2002; Hudson *et al.*, 2004]. SEPs and strong shocks are both CME-driven phenomena, and the formation of a new ion radiation belt is therefore a manifestation of CME-driven storms. The formation of a new inner electron radiation belt has been observed coincident with the compression of the magnetosphere by SSCs [Blake *et al.*, 1992; Tverskaya *et al.*, 2003; Obara and Li, 2003; Looper *et al.*, 2005]. The mechanism is believed to be a displacement of trapped electrons during a rapid compression and relaxation of the magnetosphere caused by a CME-driven interplanetary shock [Li *et al.*, 1993]. No formation of a new radiation belt has been observed in association with a CIR.

2.19. Magnetospheric Convection Interval

[26] In the superposed-epoch study of the reaction of the magnetosphere to CMEs and CIRs, Denton *et al.* [2006] examined the temporal profile of the Kp index. Kp is a measure of the strength of convection in the magnetosphere [Thomsen, 2004]. Denton *et al.* found that CMEs produce a Kp temporal profile that is peaked near the time of Dst maximum and that falls quickly (in a fraction of a day) back to prestorm levels, whereas CIRs produce a Kp temporal profile that peaks near the time of Dst maximum but falls slowly (over a few days) back to prestorm levels. Hence the interval of enhanced convection is longer for CIR-driven storms than it is for CME-driven storms. This can also be seen in both panels of Figure 1, where the durations of the storms in the 27-day recurring groups (high-speed-stream-driven storms) is notably longer than the durations of the randomly occurring storms (CME-driven storms). Note that the maximum values of Kp, i.e., the strongest peak convection, can be higher for CME-driven storms [e.g., Richardson *et al.*, 2001, Figure 1].

2.20. Great Aurora

[27] Great auroras have “high visual brightness with a progression to exceptionally low latitudes” [Jones, 1992]. An examination of the catalog of great aurora events of Jones [1992], supplemented by Yevlashin [2000], Pallamraju and Chakrabarti [2005], and Mikhalev *et al.* [2004], finds that the great aurora chiefly occur during solar maximum. A display of the solar cycle dependence of great aurora is shown in Figure 14, where the sunspot number is plotted as a function of time and the reported occurrences of great aurora are indicated by the circles. As can be seen, there is a concentration of great auroras at the solar maxima, but they occur in other phases in the solar cycle, too. In Figure 15 intervals of $K_p \geq 4^+$ (black points) and incidents of great auroras (tan points) are displayed on a 27.27-day solar-rotation versus year plot. Here, data for the 55 years 1950–2005 are shown. Recurrent storm activity (high-speed-stream-driven storms) is seen as a black patch. As

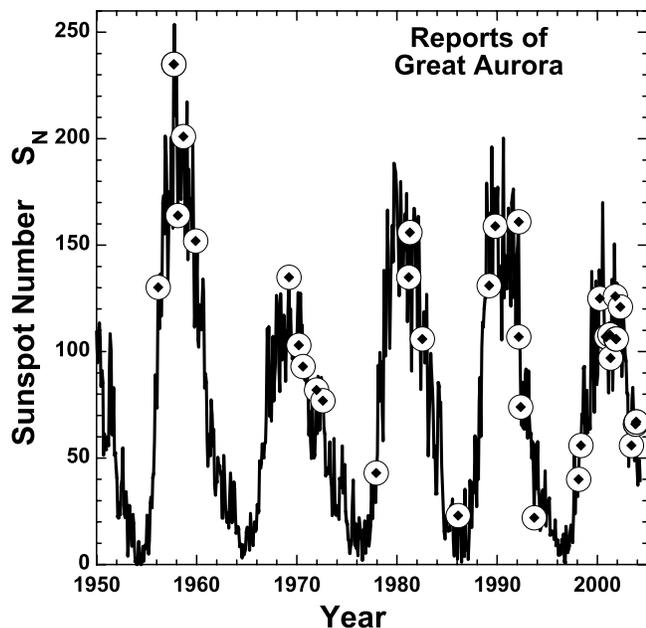


Figure 14. The sunspot number is plotted as a function of time. Occurrences of great aurora [from Jones, 1992; Yevlashin, 2000; Pallamraju and Chakrabarti, 2005; Mikhalev *et al.*, 2004] are plotted as the circles with diamond centers.

can be seen, the great aurora generally are not associated with the recurrent activity, rather they are associated with isolated storms. Perusing through Figure 15 at high resolutions, two exceptions are found: 13–14 July 1982, where a great aurora incident occurred during a recurrent storm, and 13 September 1993, where a great aurora incident occurred during a recurrent high-speed stream. The July 1982 incident was flare-associated, so this major storm was likely CME associated. The September 1993 incident was not flare-associated and no solar wind data is available to discern whether or not a magnetic cloud was embedded in the high-speed wind; no SSC occurred, so if a cloud was present, it was not extraordinarily fast. From Figure 15 and the subsequent analysis it is concluded that 1 out of 44 incidents of great aurora ($\sim 2\%$) are caused by CIR-driven storms. Of the four incidents of great aurora that occurred during the years of the McPherron CIR catalog (the years 1994–1996 and 2002–2003), none was temporally associated with CIRs. Of the 11 incidents of great aurora that occurred during the years when both the Cane and Richardson CME catalog and the Lepping magnetic cloud catalog were available (the years 1998–2002), 10 were temporally associated with CMEs ($\sim 91\%$). Hence from all of this evidence it is concluded that great aurora are associated chiefly with CME-driven storms and rarely with CIR-driven storms. Supporting this notion, Jones [1974] points out a study finding that low-latitude aurora “can usually be associated with solar flares occurring 1 to 4 days earlier.” Also, Legrand and Simon [1988] examined the latitudes of CME-associated and CIR-associated (“shock” and “stream”) aurora and found that the stream-associated aurora did not progress equatorward beyond 56° , whereas the CME-associated aurora did. Consistent with this, of the 26 low-latitude auroral events measured by STELAB in

1989–2004 [Shiokawa *et al.*, 2005], only one event occurred during the 1993–1995 era of recurring high-speed streams. Since the aurora do not progress as far equatorward for CIR-driven storms, it may also follow that the polar-cap area does not get as large for CIR-driven storms as it can during CME-driven storms.

2.21. Geomagnetically Induced Currents (GIC)

[28] Geomagnetically induced currents that are hazardous to Earth-based electrical power systems are believed to be caused by rapid intensification of ionospheric currents that

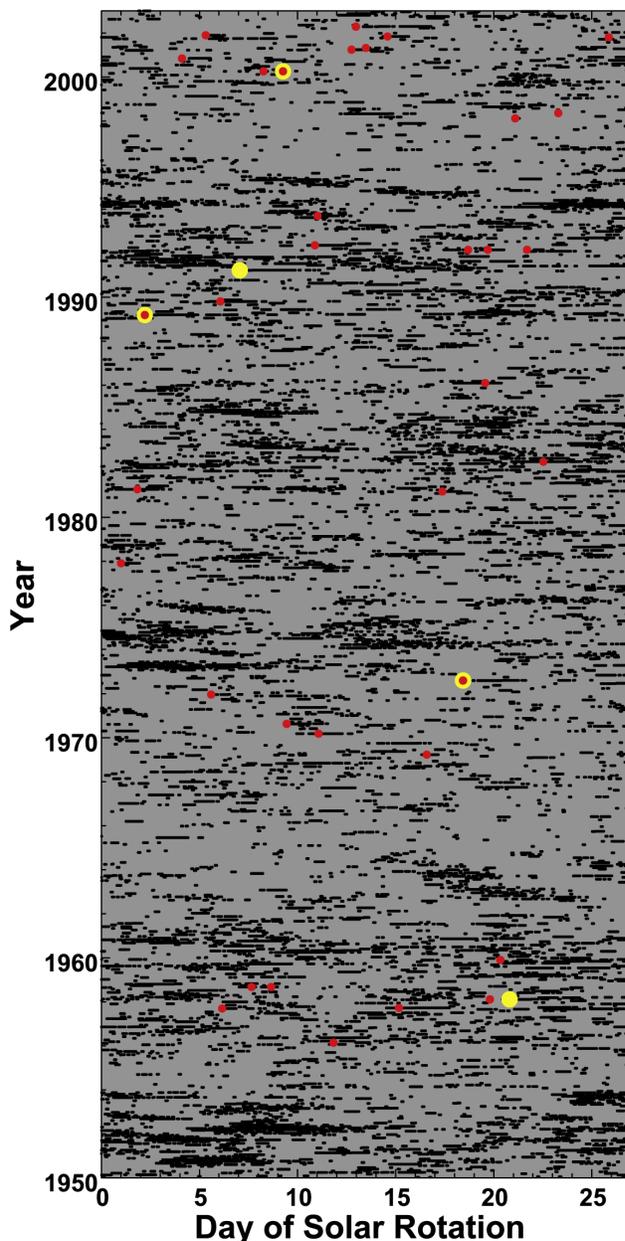


Figure 15. Times of high geomagnetic activity ($K_p \geq 4^+$) are plotted as the black boxes on a solar rotation versus year plot. Historical occurrences of great aurora are plotted as the red circles and historical occurrences of damaging geomagnetically induced currents (GIC) are plotted as the yellow circles.

can be driven by shock compression of the magnetosphere or by a substorm [Pulkkinen *et al.*, 2005]. Of the four events in the catalog of Boteler [2001], plus the March 1991 event [Bolduc, 2002], the April 2000 event [Pulkkinen *et al.*, 2003], and the October 2003 event [Kappenman, 2005], all have occurred at solar maximum or in the early declining phase. This is demonstrated in Figure 16 where the sunspot number is plotted versus time with the seven GIC events indicated as the circles with diamonds. From direct solar wind measurements or from their temporal association with large solar flares, all seven of these GIC events are associated with CME-driven storms. Examining their locations relative to storms on solar rotation versus time plots (see Figure 15), none of the GIC events are associated with 27-day recurring geomagnetic activity. Hence it is concluded that hazardous GIC are predominantly associated with CME-driven storms rather than CIR-driven storms.

3. Other Stormtime Differences

[29] Several stormtime phenomena have not been or could not be analyzed for inclusion in Table 1. Some of these pertain to observations of the ionosphere. It is hoped that the discussion of this section will motivate future research to clarify how these phenomena fit into CME-driven and CIR-driven geomagnetic storms.

[30] Large equatorial bubbles, which appear in the topside ionosphere when Dst is increasing in magnitude [Huang *et al.*, 2002; Basu *et al.*, 2005], have been reported for many CME-driven magnetic storms [e.g., Huang *et al.*, 2001; Abdu *et al.*, 2003; Basu *et al.*, 2005]. Large equatorial bubbles do occur for Dst perturbations at solar minimum but at a much lesser rate than they do at solar maximum [Huang *et al.*, 2002]. The bubbles are attributed to penetration electric fields associated with build up of the ring current: since ring current buildup is greater for CME-driven storms, one would expect the bubbles to be primarily a CME-driven storm phenomenon.

[31] SAPS (subauroral polarization stream) and SAID (subauroral ion drift) ionospheric flow streams are associated with a deep penetration of the ion plasma sheet into the dipolar magnetosphere [Anderson *et al.*, 1993; Garner *et al.*, 2004]. SAPS and SAID events have been observed and analyzed for several CME-driven storms [e.g., Foster *et al.*, 1994; Garner *et al.*, 2004; Foster and Vo, 2002], but no comparisons have been made for their occurrence probabilities and properties for CIR versus CME-driven storms.

[32] A recently discovered phenomenon is the disappearance of AKR (auroral kilometric radiation) during CME-driven storms coincident with the appearance of a superdense plasma sheet [Seki *et al.*, 2005; Morioka *et al.*, 2003]. The AKR-disappearance events in the literature are all associated with CME-driven storms, except the 11 January 1994 event, which corresponds to a superdense plasma sheet in a recurring CIR-driven storm. A comparison of the nature of the AKR disappearance for the hotter superdense plasma sheets of CIR-driven storms and the cooler superdense plasma sheets of CME-driven storms has not been made.

[33] The injection of cool dense plasma sheet (CDPS) material (which is probably low-latitude boundary layer (LLBL) plasma) into the nightside of the dipolar magneto-

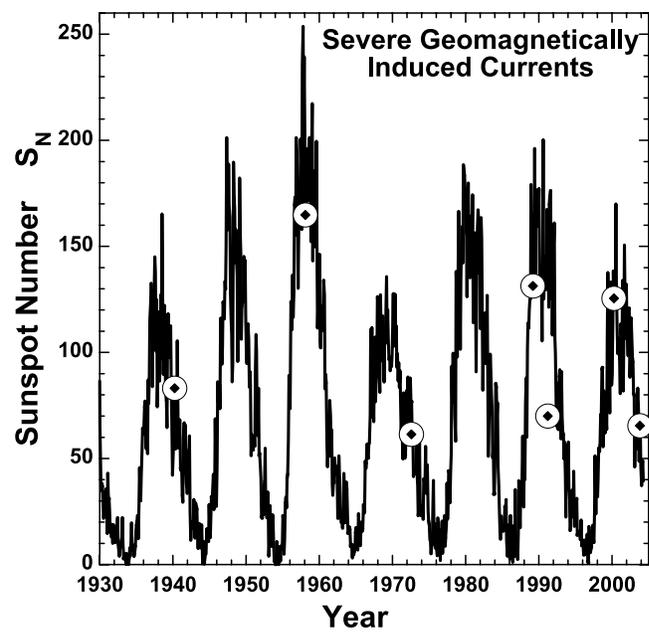


Figure 16. The sunspot number is plotted as a function of time. Occurrences of hazardous GIC (geomagnetically induced current) [from Boteler, 2001; Bolduc, 2002; Pulkkinen *et al.*, 2003; Kappenman, 2005] are plotted as the circles with diamond centers.

sphere has been documented to occur for both types of storms. The collection of all known nightside dense-plasma injections find that they are statistically found at times of elevated Kp (3.1 ± 1.2) and mildly elevated values of Dst (-21 ± 21 nT) [cf. Lavraud *et al.*, 2005]. This injection has been seen for CME-driven storms (15 of the 30 events in Table 1 of Thomsen *et al.* [2003] are in the Cane and Richardson [2003] catalog of shock-sheath-CME events) and this injection has been estimated to occur for $\sim 50\%$ of high-speed-stream-driven storms [Borovsky and Steinberg, 2006].

[34] One phenomenon that is similar for both types of storms is the appearance of a plasmaspheric drainage plume in the dayside magnetosphere in the early phases of a storm. As magnetospheric convection increases, the outer portions of the plasmasphere are transferred from closed cold plasma drift trajectories to open trajectories, and the material of the outer plasmasphere convects from the middle magnetosphere to the dayside neutral line. Plasmaspheric drainage plumes have been studied for CIR-driven storms [Borovsky *et al.*, 1998b] and for CME-driven storms [Thomsen *et al.*, 1998a]. A comparison of the details of drainage plume phenomenon for the two types of storms has not been made.

[35] The tongue of ionization and polar cap patches are higher-density ionospheric plasma convecting into and across the polar cap from the dayside. They are related to the plasmaspheric drainage plume seen in the magnetosphere when geomagnetic activity increases [Su *et al.*, 2001]. High-density polar cap events have been investigated for several CME-driven storms [e.g., Sojka *et al.*, 1992; Crowley *et al.*, 2000; Foster *et al.*, 2005], but a comparison of their occurrence rate and properties between CME-driven storms and CIR-driven storms has not been performed.

[36] In this report, no attempt has been made to sort out the occurrences and intensities of various plasma waves in

the magnetosphere that are thought to be important for the physics of storms, such as whistler-mode chorus [Villalon and Burke, 1995; Smith et al., 2004], electromagnetic ion cyclotron waves [Jordanova et al., 1998; Summers and Thorne, 2003], and plasmaspheric hiss [Horne et al., 2003].

[37] Finally, the authors have made no attempt here to compare the morphology and phenomenology of the aurora during CME-driven storms and CIR-driven storms.

4. Summary

[38] The differences between CME-driven storms and CIR-driven storms were explored and the findings were collected into Table 1. In section 2 the 21 phenomena listed in Table 1 were discussed, item by item.

[39] In a nutshell, CME-driven storms are brief, have denser plasma sheets, have stronger ring currents and Dst perturbation, have solar energetic particle events, and they can produce new radiation belts, great auroras, and dangerous geomagnetically induced currents. CIR-driven storms are of longer duration, have hotter plasma sheets and hence stronger spacecraft charging, and produce higher fluxes of relativistic electrons. The magnetosphere is more likely to be preconditioned by a geomagnetically calm interval prior to a CIR-driven storm than it is prior to a CME-driven storm. CME-driven storms are more hazardous to Earth-based systems and CIR-driven storms are more hazardous to space-based assets, particularly at geosynchronous orbit. CME-driven storms occur randomly in time with an occurrence frequency that is highest during solar maximum whereas CIR-driven storms recur with a 27-day period primarily in the declining phase of the solar cycle.

[40] Analogous to geomagnetic storms, terrestrial thunderstorms also vary, with differing outputs of rain, wind, hail, lightning, and tornadoes. Also, thunderstorms have different types of drivers: radiation-driven convection (the island effect) or air-mass interaction (frontal uplift). Historically, and still typically, geomagnetic storms are gauged only by the size of their Dst signature [e.g., Gonzalez et al., 1994; Loewe and Pross, 1997]: this is akin to gauging the severity of thunderstorms only by their output of hail.

[41] One lesson that can be taken from this study is that when geomagnetic storms are studied, CME-driven storms and CIR-driven storms should be studied separately. A second lesson that can be taken from this study is that Dst alone is a poor indicator of the properties of storms.

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