

Neutron reflection from the liquid helium surface

T. R. Charlton¹, R. M. Dalgliesh¹, A. Ganshin², O. Kirichek¹, S. Langridge¹,
P. V. E. McClintock²

¹*ISIS, STFC, Rutherford Appleton Laboratory, Harwell, Didcot, UK*

²*Lancaster University, Lancaster, UK*

The importance of liquid helium for physics can hardly be overestimated. Its unique ability to remain liquid right down to absolute zero makes helium an ideal experimental sample for studying quantum phenomena in condensed matter. Helium-4 was first liquefied by Kamerlingh Onnes in 1908. Since then the properties of bulk helium liquid (including ⁴He, ³He and isotopic mixtures) have been studied in detail, resulting in the discovery of many fundamental phenomena including, especially, superfluidity. However the free surface of liquid helium has attracted significantly less attention than the bulk liquid, perhaps due to the relatively complex experimental methods that are needed for surface studies. Currently the best-known approaches are probably optical ellipsometry, surface tension measurements, and investigations of two-dimensional charge sheets above and below the surface.

Optical ellipsometry was exploited for measuring the thickness and density profile of the liquid helium surface [1]. It was shown in experiments that the vertical thickness within which the density changes from 90 to 10% of the bulk liquid density increases slightly with temperature from 1.4 to 2.1 K, with an average value of 9.4 Å at 1.8 K. There was also evidence of a rise of up to 3 Å between 2.1K and the superfluid transition temperature (λ point). A very close value for the vertical thickness, of 9.2 ± 1 Å at 1.13 K, was obtained in experiments where liquid-vapor density profiles of thick ⁴He films were measured using X-ray reflectivity [2].

The surface tension of liquid helium can be measured by various experimental methods [3, 4] and may be compared directly with the values derived from theoretical models [5, 6]. However, establishing the relationship between surface tension and the dynamical properties of superfluid helium remains a rather complex task [7].

Perhaps the most powerful experimental technique for studying the surface properties of liquid helium is through measurements on two-dimensional charge systems including surface electrons (SE) and layer of ions trapped under the surface of liquid helium [8]. The existence of the quantized capillary surface waves, or ripplons, earlier proposed by Atkins [7] in order to explain surface tension data, has been clearly proven in experiments with a two-dimensional plasma resonance in surface electrons [9]. Later, ripplons [10] as well as scattering of surface electrons with Fermi quasi-particles from bulk liquid [11] were detected on liquid ³He. Soon after that, it was found that the mobility of surface electrons on ³He deviates drastically from the single-electron-ripplon scattering theory below 70 mK [12]. In contrast, the mobility of surface electrons on ⁴He follows the single-electron-ripplon scattering theory down to 20 mK, the lowest temperature achieved in the experiment. This may suggest a crossover to the long-mean-free-path regime, implying a connection to the Rayleigh-like waves predicted theoretically in [13, 14]. A peculiar temperature dependence was also observed during

studies of the surface tension of liquid ^3He [4], behavior that could not be explained in terms of ripples and the quasi-particle-scattering contribution.

Capillary waves on the surface of thick superfluid ^4He films have been detected via the dielectric ponderomotive effect [15]. This method allowed measurement of the spectrum and damping of capillary waves within the wavelength range $3.3 - 20 \mu\text{m}$. The damping was orders of magnitude smaller than the value predicted theoretically.

In some respects, the SE, surface tension and capillary wave methods provide rather indirect ways of studying the liquid helium surface experimentally, inevitably leading to ambiguities in the interpretation of the data. A complementary experimental approach that might facilitate direct studies of the microscopic properties of liquid helium surface, thus illuminating this fascinating subject, is therefore much to be desired.

In this paper we propose a new experimental method based on the unique combination of neutron reflection and ultra-low temperature sample environment: small-angle neutron reflection from the liquid surface, opening new opportunities for studying the interface properties of quantum liquids. We also present our preliminary measurements of reflection from the superfluid and normal liquid ^4He surfaces.

Note that a range of classical optical phenomena such as reflection, refraction and interference can be also observed for slow neutrons [16, 17]. For a neutron approaching a plane surface at a glancing angle θ , the refractive index at the boundary between two media is defined as:

$$n = \frac{k'}{k} = \frac{\cos \theta}{\cos \theta'} \quad (1)$$

where k' and k are the neutron wave vectors inside and outside the medium. Since $n < 1$, a limiting case occurs when $\cos \theta' = 1$, characterized by the critical angle θ_c such that $\cos \theta_c = n$. For $\theta < \theta_c$ we have total external reflection. The phenomenon is analogous to the total internal reflection of light; the difference is in the sign of $(n - 1)$. For the thermal neutrons $\theta_c \ll 1$ we can write $\cos \theta \approx 1 - \theta^2/2$, which leads to:

$$\theta_c = \sqrt{2 \cdot (1 - n)} \quad (2)$$

Total reflection of slow neutrons was first reported by Fermi and co-workers [18, 19] and since been extensively applied to range of problems. The specular reflection of neutrons gives information on the neutron refractive index profile normal to a surface. The refractive index is simply related to the scattering length density and hence specular neutron reflection can provide important information about the composition of surfaces and interfaces. Similar information can be obtained from x-ray reflection experiments. There are many instances, however, where the neutron method provides distinct advantages. For example the neutron's ability to penetrate through bulk samples and their containers, including cryostats, is much higher than that of x-rays. It allows for an easy combination of neutron scattering with an ultra low temperature sample environment.

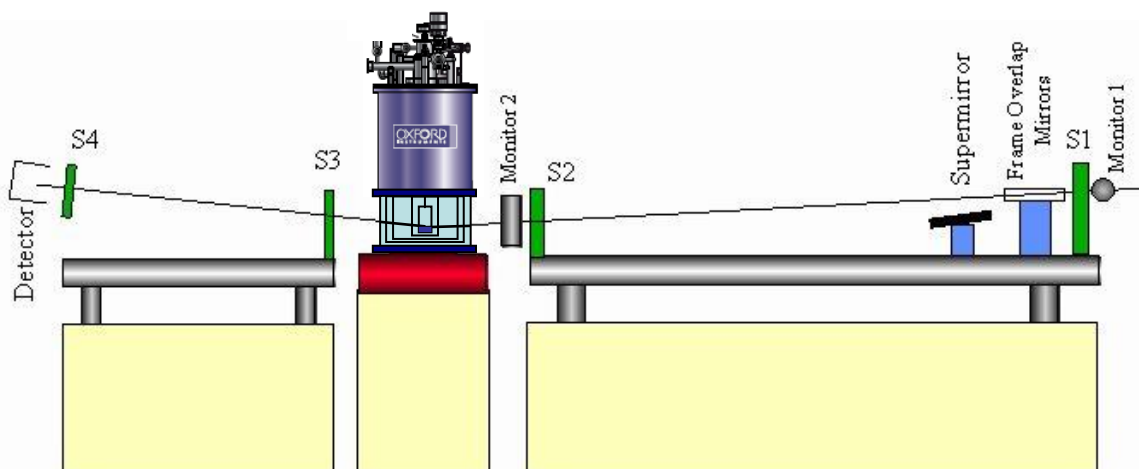


Fig. 1 Schematic diagram of the CRISP reflectometer showing the major components. Items used for polarization handling have been removed from the drawing since they were not used in the experiments described.

Our experiments were performed on the CRISP Instrument, which was designed as a general purpose reflectometer for the investigation of a wide spectrum of interfaces and surfaces. It uses a broad-band neutron time-of-flight method for determination of the wavelength at fixed angles. Fig. 1 shows a schematic diagram of the instrument. 1. The sample geometry is horizontal, allowing investigation of liquid surfaces. The instrument uses the pulsed neutron beam from the ISIS neutron facility and passage via a 20 K hydrogen moderator provides neutron wavelengths in the range 0.5 - 6.5 Å at the source frequency of 50 Hz, extending up to a maximum of 13 Å if operated at 25 Hz. The incident beam is well collimated by both coarse and fine adjustable collimating slits (S1 - S4) to give a variable beam size and angular divergence, with typical dimensions of 40 mm wide (horizontal direction) and anything up to 10 mm in height (vertical direction). A variable aperture disc chopper defines the wavelength band, and prompt pulse suppression is achieved by a nimonic chopper. Additional frame overlap suppression is provided by the nickel-coated silicon wafer frame overlap mirrors, which remove wavelengths greater than 13 Å. The instrument has two types of detectors, a He single detector and a 1-D position-sensitive multidetector. The instrument can also be run in a polarized mode allowing polarized neutron reflectivity measurements as well being able to perform full polarization analyses.

The experimental arrangement is extremely flexible and, with the neutron beam inclined at 1.5° to the horizontal, provides well for the study of liquid surfaces. Provision for liquid surfaces angles of less than 1.5° is achieved by the insertion of a supermirror. The sample position is designed to be vibrationally isolated from the rest of the instrument, further aided by active anti-vibrational damping. Much of the instrument is automated allowing for precision control and a high degree of reproducibility.

A large variety of sample environment equipment including cryostats and ultra low temperature refrigerators is available. For the experiment to be described, we used an Oxford Instruments Variox^{BL} cryostat for neutron scattering experiments with a variable temperature insert (VTI). The temperature range of the VTI is 1.25 - 300K. A cylindrical cell made of aluminium alloy 6082 with inner diameter 31.4 mm, internal length 74.0 mm

and wall thickness 0.5 mm was attached to the sample stick of the VTI. The temperature of the cell was measured with a calibrated Cernox sensor.

In order to estimate the influence of the cell and sample environment on the signal from the scattered neutrons, we performed neutron reflection calibration measurements from the surface of liquid D₂O at room temperature in the same experimental setup. The test demonstrated a negligible influence of the sample environment on the experimental data.

In our experiments we attempted to study neutron reflection from the surfaces of both liquid ⁴He and ³He. However, due to the large cross-section for neutron scattering from the ³He nucleus in the vapour phase, we could not detect a reflection signal from the ³He liquid surface within the experimental temperature range down to 1.25K. The ³He vapour pressure decreases rapidly with temperature, and we therefore we believe that the temperature reduction below ~1K should solve this problem.

In the experiment we condense approximately 10 litres of helium gas (at STP) into the cell, which is kept at a temperature of 2 K. After the reflection signal is detected, the temperature of the cell is stabilized and controlled for a sufficient period of time to collect neutron reflection data over a range of values of wave vector transfer Q_z perpendicular to the reflecting surface (see the diagram in the inset of Fig. 2), which typically takes a few hours.

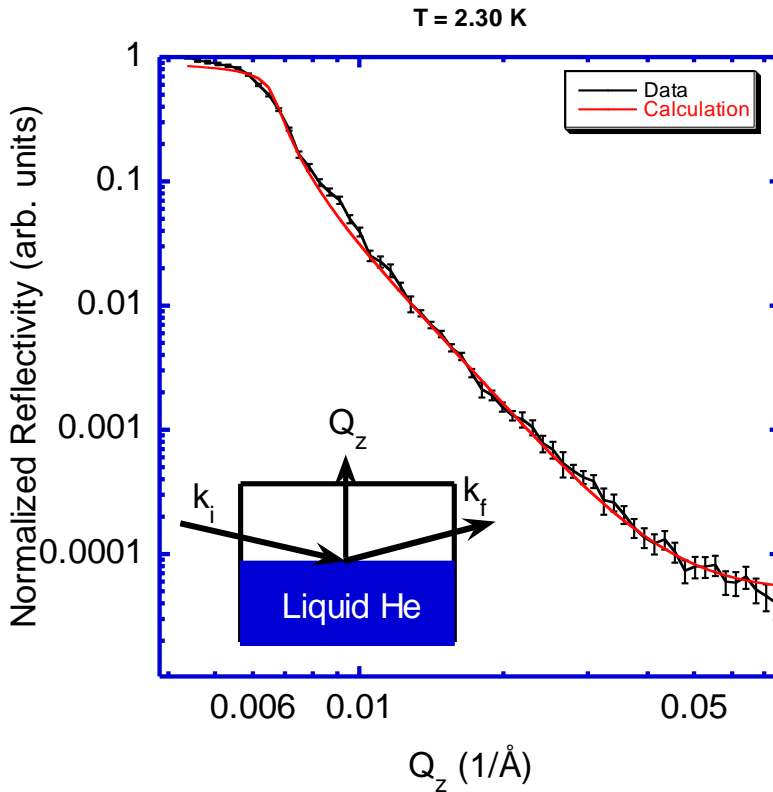


Fig. 2 The reflectivity of the liquid ⁴He surface as a function of Q_z measured at 2.3 K. The inset shows the orientation of the incoming k_i and outgoing k_f wave vectors as well as the momentum transfer vector Q_z in relation to the liquid helium surface.

We have measured the reflectivity of the liquid ^4He surface as a function of Q_z at temperatures of 2.3 K and 1.54 K, obtaining the results presented in Figs. 2 and 3 respectively. As a general feature both data sets show a slight downward slope of the reflectivity approaching the critical value, $Q_c \sim 0.006 \text{ 1/\AA}$, attributable to absorption of neutrons by traces of ^3He at the natural isotopic ratio of about 2×10^{-7} in the supplied ^4He gas. Without absorption the signal below Q_c would be completely flat. Past Q_c , from $Q_z > 0.08 \text{ 1/\AA}$ where the Born approximation applies we observe the typical $\sim Q_z^{-4}$ decrease in the reflectivity curve indicating, qualitatively, that the surface is smooth.

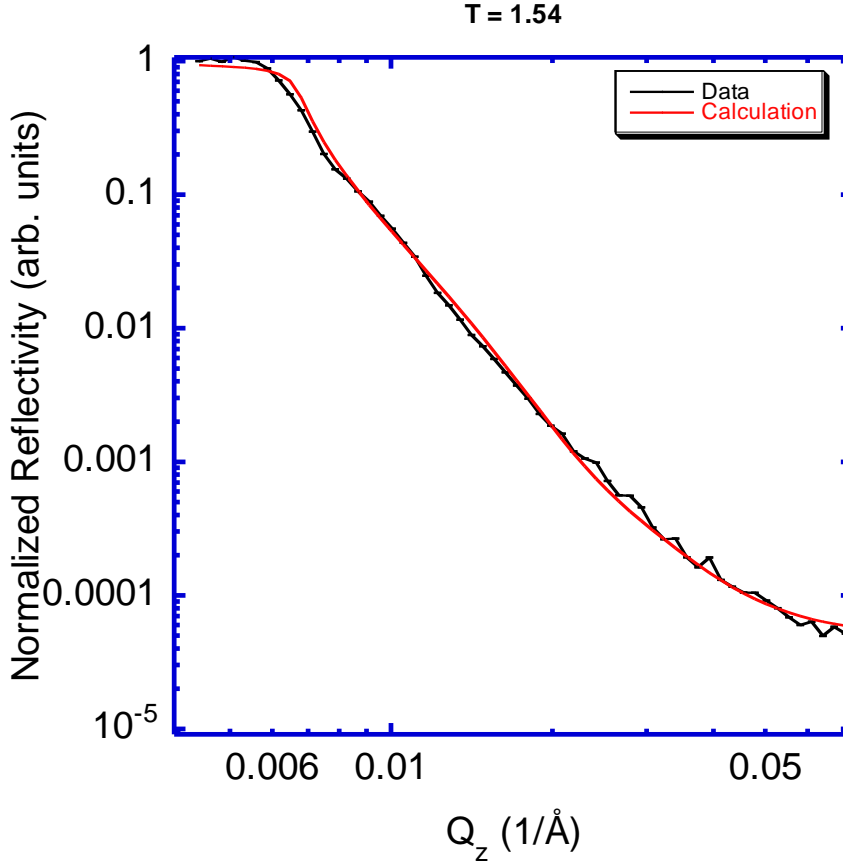


Fig. 3 The reflectivity of liquid ^4He surface as a function of Q_z measured at 1.54 K.

To extract quantitative information from the data we compare the reflectivity profile with an optical model. In its simplest form, the model reduces to the quantum mechanical problem of a particle incident on a potential step. In our case, the optical potential has three distinct regions (vacuum, near surface, and bulk) with the transition between adjacent steps calculated as an interface roughness. The reflectivity is calculated from the optical potential using a recursive definition of the reflectivity, taking into account multiple reflections [20]. The numerical description of the interface roughness follows that of reference [21]. The results of the calculation are shown as lines through the data in Fig. 2 and 3 with the optical potential shown in Fig. 4. From the calculations it is clear that, near the surface, a region with higher scattering power exists in the 1.54K case as compared to the 2.30K calculations. Also the transition from bulk to vacuum at

2.30K occurs gradually over a large 800 Å distance while, at 1.54K, the transition is compressed into ~400 Å. We can suggest a possible explanation for this observation.

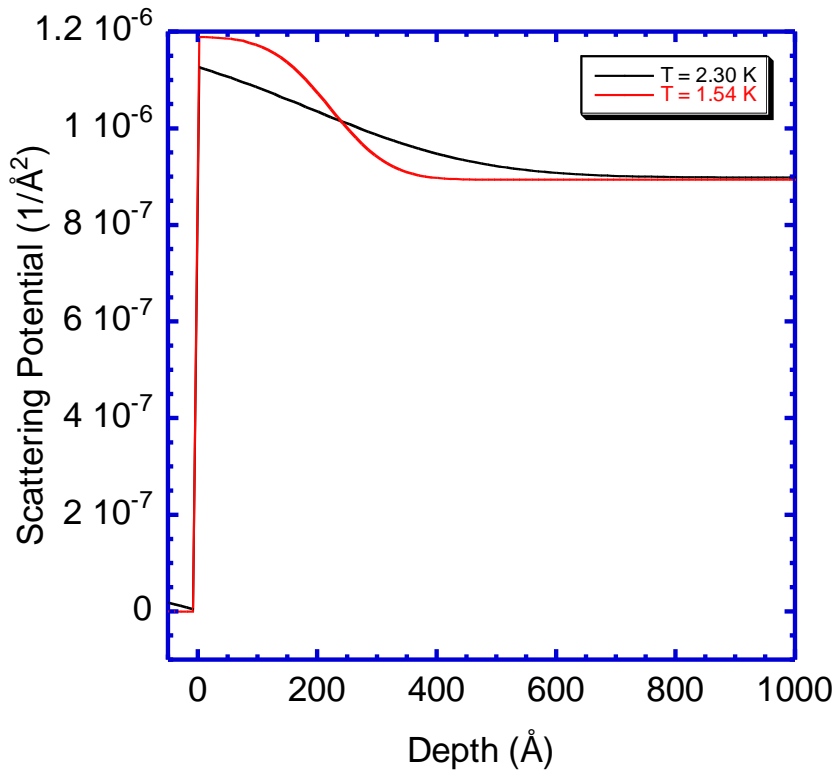


Fig. 4 The scattering potential profiles obtained by numerical simulation, using the reflectivity data for 2.3 K and 1.54 K. The abscissa corresponds to vacuum below 0 Å, to the near-surface region between 0 Å and 800 Å, and to bulk liquid ^4He above 800 Å. Note that the optical potential is proportional to the density and scattering power of the material.

Experiments on liquid ^4He with both neutron scattering [22, 23] and X-ray diffraction [24] indicate a significant change in the structure of liquid as it cools through the superfluid transition. This observation was also supported by Monte Carlo simulations [25] where it was shown that, when Bose condensation takes place, the liquid ^4He becomes more disordered in Cartesian space with some rearrangement of atoms occurring to create the necessary space for delocalization to occur. If there is an enhanced Bose-Einstein condensate fraction close to the surface of the superfluid [26, 27], an increase in disorder could be expected, possibly accounting for the alteration in the neutron scattering structure factor. However thorough theoretical modelling is now required to establish whether the results observed in the experiment can be accounted for quantitatively on this basis.

In conclusion, we have observed neutron reflection from the ^4He surface for the first time. The surface is smoother in the superfluid state at 1.54 K than in the case of the

normal liquid at 2.3 K. We have also observed a surface layer ~ 200 Å thick with slightly different neutron scattering structure factor in superfluid ^4He . The new experimental method described in this paper opens up fresh opportunities in the study of quantum fluids and solid quantum surfaces and interfaces. For example it could be used for direct observation of Fomin surface excitations on liquid ^3He [13], or in the search for a superfluid transition in 2D ^3He on the surface of micro-separated ^3He - ^4He liquid mixture [28]

References

1. D.V. Osborne, *J. Phys.: Condens. Matter* **1**, 289 (1989)
2. L.B. Lurio, T.A. Rabedeau, P.S. Pershan, Isaac F. Silvera, M. Deutsch, S.D. Kosowsky, and B.M. Ocko *Phys. Rev. Lett.* **68**, 2628, (1992)
3. K.N. Zinoveva, *ZhETF* **29**, 899 (1955)
4. K. Matsumoto, Y. Okuda, M. Suzuki, and S. Misawa *J. of Low Temp. Phys.* **125**, 59 (2001)
5. K.R. Atkins *Can. J. Phys.* **31**, 1165 (1953)
6. A.D. Singh, *Phys. Rev.* **125**, 802 (1962)
7. K.R. Atkins and Y. Narahara *Phys. Rev.* **138**, A437 (1965)
8. Yu. Monarkha and K. Kono *Two-Dimensional Coulomb Liquids and Solids*, Springer (2004)
9. C.C. Grimes and G. Adams *Phys. Rev. Lett.* **36**, 145, (1976)
10. K. Shirahama, S. Ito, H. Suto and K. Kono *J. Low Temp. Phys.* **101**, 439 (1995)
11. K. Shirahama, O.I. Kirichek and K. Kono *Phys. Rev. Lett.* **79**, 4218, (1997)
12. O. Kirichek, M. Saitoh, K. Kono, F.I.B. Williams *Phys. Rev. Letts*, **86**, 4064 (2001)
13. I.A. Fomin, *Sov. Phys. JETP* **34**, 1371 (1972)
14. Yu.B. Ivanov, *Sov. Phys. JETP* **52**, 549 (1980)
15. P. Roche, G. Deville, K.O. Keshishev, N.J. Appleyard, and F.I.B. Williams *Phys. Rev. Lett.* **75**, 3316, (1995)
16. A.G. Klein and S.A. Werner *Rep. Prog. Phys.* **46**, 259 (1982)
17. J. Penfold and R.K. Thomas *J. Phys.: Condens. Matter* **2**, 1369 (1990)
18. E. Fermi and W. Zinn *Phys. Rev.* **70**, 103 (1946)
19. E. Fermi and W. Marshall *Phys. Rev.* **71**, 666 (1947)
20. L. G. Parratt, *Phys. Rev.*, **95(2)**, 359 (1954)
21. L. Nevot and P. Croce, *Rev. Phys. Appl.* **15**, 761 (1980)
22. V.F. Sears and E.C. Svensson, *Phys. Rev. Lett.* **43**, 2009 (1979)
23. E.C. Svensson, V.F. Sears, A.D.B. Woods, and P. Martel *Phys. Rev. B* **21**, 3638 (1980)
24. H.N. Robkoff, D.A. Ewen, and R.B. Hallock *Phys. Rev. Lett.* **43**, 2006 (1979)
25. J. Mayers, *Phys. Rev. Lett.* **84**, 314 (2000)
26. A. Griffin and S. Stringari *Phys. Rev. Lett.* **76**, 259 (1996)
27. A.F.G. Wyatt *Nature* **391**, 56 (1998)
28. P.A. Sheldon and R.B. Hallock *Phys. Rev. Lett.* **77**, 2973 (1996)