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## A dynamic model for the planetary bearings in a double planetary gear set

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### Abstract

Double planetary gear sets (DPGSs) are widely applied to high power and high torque mechanical transmission systems due to the high loading capacity, steady transmission, and high transmission efficiency. Most previous works studied the planetary bearing vibration characteristics without the gear excitations. This work establishes a dynamic model for a DPGS containing all components (sun gear, ring gear, carrier, inner planet, outer planet, planetary bearing roller, and planetary bearing cage). The components' excitations in this work are more comprehensive than the reported references. The proposed model considers the planetary bearing roller and cage dynamics. Moreover, the gear interactions, planetary gear-bearing-carrier interactions, and roller-cage interactions are contained in the proposed model. The vibrations of planetary bearing roller, planet, carrier, and planetary bearing cage are analyzed. The effect of sun gear rotation speed on the planetary bearing contact force, planetary bearing roller-cage impact force, and vibrations of planetary bearing parts are studied. The simulated and experimental results are compared to prove the correctness of proposed model. Moreover, the results from proposed model and model without the cage are compared to show the advantage of proposed model.

**Keywords:** Double planetary gear set (DPGS); planetary bearing; dynamic model; vibrations

### 1. Introduction

Double planetary gear sets (DPGSs) are widely applied to high power and high torque mechanical transmission systems due to the high loading capacity, steady-transmission, and high transmission efficiency. As one of the key parts of the DPGS, the planetary bearing rotates around the DPGS axis besides its own axis. Due to the special work conditions, the planetary bearings are prone to failure during the operations. The vibrations can reflect the planetary bearing working conditions, which can be helpful for the fault diagnosis and vibration suppression. Thus, it is critical to establish the DPGS dynamic model to analyze the planetary bearing vibrations.

Many studies [1,2] were conducted to research the planetary gear set dynamics. Kahraman [3] studied the planetary gear set (PGS) vibrations by establishing a torsional dynamic model. Parker [4] proposed a two-dimensional dynamic model for PGS. The effect of the tooth profile modification on the PGS vibrations are analyzed in their work [5]. Öztürk et al. [6] proposed a new time-varying meshing stiffness (TVMS) calculation method considering the tooth profile modification to study its' effect on the PGS vibrations. Fatourehchi et al. [7] established a torsional dynamic model to find the best mesh phasing to reduce the PGS vibrations. Chen and Shao [8] proposed a TVMS model considering the ring gear tooth root crack to study its effect on the PGS vibrations. Shen et al. [9] analyzed the effect of the gear wear on the PGS vibrations. In their work, the wear depth is obtained by Archard model, and the effect of the gear wear on the TVMS and gear backlash was considered. Xiang et al. [10] proposed a torsional dynamic model of PGS including the comprehensive gear error, TVMS, and piece-wise backlash nonlinearities. Luo et al. [11] proposed a PGS dynamic model considering the sliding friction between the meshing contact. They analyzed the effect of the spalling defects on the PGS vibrations. Khoozani et al. [12] researched the effect of the gyroscopic effects on the PGS

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vibrations. Kim et al. [13] pointed out that the pressures angles and contact ratios will be changed with the time, which will affect the PGS vibrations. Thus, they proposed a PGS dynamic model including the time-varying pressure angles and contact ratios. Xu et al. [14] proposed a TVMS model containing the positioning errors and studied the effect of positioning errors on the PGS vibrations. Hu et al. [15] established a multistage PGS dynamic model and studied the effect of sun gear crack on the PGS vibrations. Wang et al. [16] proposed a DPS-bearing-rotor system dynamic model and analyzed the effect of support bearing misalignments on the system vibrations. The above analysis indicates that most researchers ignored the planetary bearings in dynamic modeling. However, as one of the most important parts in the PGS, the planetary bearing is one of important excitation sources during the actual operations.

The planetary bearing was considered in some studies. Guo and Parker [17] proposed a dynamic model considering the planetary bearing clearance. In their work, the roller was approximated as a spring system; and the roller dynamic characteristics were not considered. Liu et al. [18,19] established a flexible-rigid coupling dynamic model for PGS. They studied the effect of the planetary bearing fault and clearance on the PGS vibrations. Liu et al. [20] developed a dynamic finite element (FE) model for a PGS. They studied the effect of a planetary bearing localized fault on the PGS vibrations. Denni et al. [21] proposed a dynamic model for planetary bearing. In their work, they considered the hydrodynamic lubrication effect in the roller/cage contact. However, the proposed model ignored the effect of the gear meshing excitation. They studied the planetary bearing separately. Moshrefzadeh and Fasana [22] developed a dynamic model for PGS considering the bearing local defect. However, the planetary bearing dynamics were ignored in their work. Based on the above discussions, note that the previous works studied the PGS vibrations ignoring the bearing dynamics, or they studied the planetary bearing vibrations ignoring the gear excitations. However, the planetary bearing is one of the parts of PGS, the planetary bearing vibrations obtained by a dynamic model without the gear excitations are not reasonable. Liu et al. [23,24] studied the effect of the cage flexibility on the bearing skidding dynamic performance under accelerating conditions. Wu et al. [25] studied the distribution of lubricating oil in roller bearing. Besides, the vibration studies for DPGSs are fewer. Thus, it is important to propose a new dynamic model for DPGS including the bearing and gear excitations comprehensively.

This work establishes a dynamic model for DPGS considering the bearing and gear excitations comprehensively. The proposed model considers the planetary bearing roller and cage dynamics. Moreover, the gear interactions, planetary gear-bearing-carrier interactions, and roller-cage interaction are considered in the proposed model. The TVMS is obtained by the method in Ref. [26]. In the proposed method, the mesh deformations are calculated by the gear relative meshing displacements. Moreover, the bearing contact force of two contact parts is calculated by the relative displacement in a Cartesian coordinate system. Besides, the roller-cage impact force is obtained by the relative angular displacement between the roller and cage. The friction moment calculation method in Ref. [27] is used. An experiment is conducted to prove the correctness of proposed model. The effect of the sun gear rotation speed on the planetary bearing vibrations is studied. In comparison with the previous studies, this work establishes an improved dynamic model for DPGS including all components (sun gear, ring gear, carrier, inner planet, outer planet, planetary bearing roller, and planetary bearing cage). The components' excitations in this work are more comprehensive. Moreover, the results from proposed model and model without the cage are compared to show the advantage of proposed model.

## **2. Problem description**

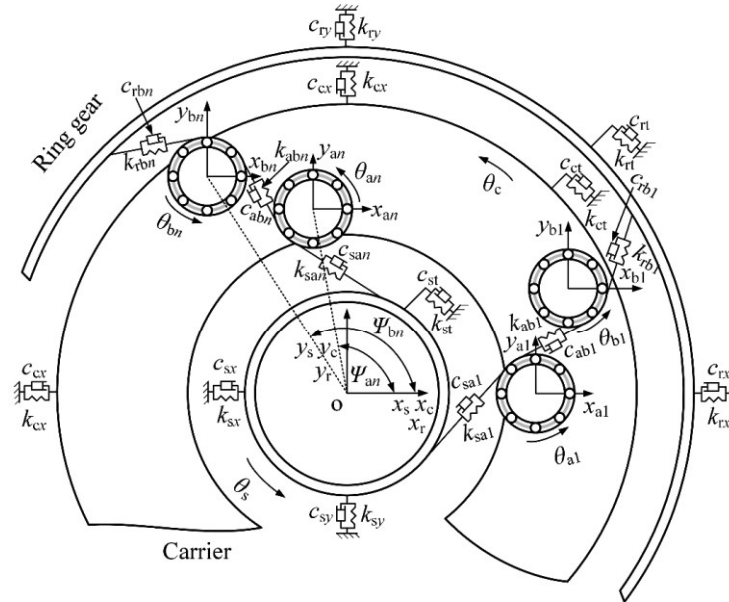
As one of the key parts of DPGS, the planetary bearing rotates around the DPGS axis besides its own axis. Due to the special work conditions, the planetary bearings are prone to

failure during the operations. However, few studies studied the planetary bearing dynamics and ignored the gear excitations. In the actual operations of a DPGS, the excitations are complicated, which include the gear mesh excitations, TVMS excitations, bearing roller-ring contact excitations, bearing cage-ring impact excitations, and bearing roller-cage impact excitations, as shown in Fig. 1. Thus, the planetary bearing vibrations obtained by the model ignored the gear tooth deformation excitations and TVMS excitations are not reasonable. Thus, an improved dynamic model for DPGS including all components (sun gear, ring gear, carrier, inner planet, outer planet, planetary bearing roller, and planetary bearing cage) is proposed in the following sections.

### 3. Dynamic modeling of a DPGS

The dynamic models presented in this paper are based on the following assumptions: (1) Each planetary gear is uniformly distributed along the circumference; and their mass and rotational inertia are same. (2) The counterclockwise is defined as the positive direction of rotation. (3) The effect of the gear backlash on the TVMS and mesh deformation are ignored. (4) The flexibilities of components are ignored too.

(a)



(b)

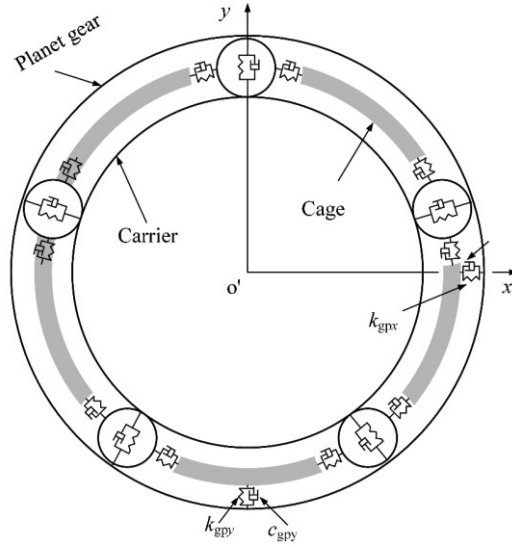


Fig. 1. (a) A dynamic model of DPGS and (b) planetary bearing dynamic model.

Fig. 1 shows the proposed DPGS dynamic model. The DPGS consists of one ring gear (r), one carrier (c), one sun gear (s),  $N$  inner planet, and  $N$  outer planetary gear. Three degrees of freedom of each part are considered, which include the transverse movement along the  $x$ - and  $y$ -axis, and rotation about the  $z$ -axis. A fixed coordinate system with the carrier center as the origin,  $N$  rotating coordinate systems with the inner planetary gear center as the origin, and  $N$  rotating coordinate systems with the outer planetary gear center as the origin. The rotating coordinate systems rotate around the fixed coordinate system  $o$ . In the proposed model, the motion of sun gear, ring gear, and carrier are studied in the fixed coordinate system, the motion of inner/outer planet gears, inner/outer planetary bearing rollers and inner/outer planetary bearing cages are studied in the rotating coordinate systems. The dynamic equations of DPGS are given as follows.

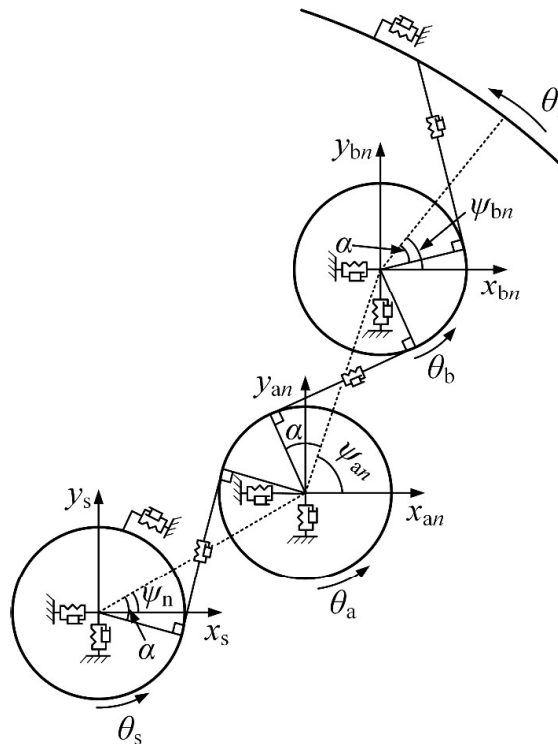


Fig. 2. Gear relative meshing displacements.

The dynamic equations of sun gear are expressed as

$$m_s \ddot{x}_s + c_{sx} \dot{x}_s + k_{sx} x_s + \sum F_{san} \cos \psi_{sn} = m_s x_s \Omega^2 + 2m_s \dot{y}_s \Omega + m_s y_s \dot{\Omega} \quad (1)$$

$$m_s \ddot{y}_s + c_{sy} \dot{y}_s + k_{sy} y_s + \sum F_{san} \sin \psi_{sn} = m_s y_s \Omega^2 - 2m_s \dot{x}_s \Omega - m_s x_s \dot{\Omega} \quad (2)$$

$$\frac{J_s}{r_s} \ddot{\theta}_s + \sum F_{san} = \frac{T_i}{r_s} \quad (3)$$

where,  $r_s$  is the sun gear base radius;  $\Omega$  is the carrier rotation velocity;  $m_s$  is the sun gear mass;  $k_{sx}$  and  $c_{sx}$  are the sun gear bearing stiffness and damping coefficients in  $x$  direction;  $k_{sy}$  and  $c_{sy}$  are the sun gear bearing stiffness and damping coefficients in  $y$  direction;  $J_s$  is the sun gear moment of inertia;  $T_i$  is the DPGS input torque;  $x_s$  and  $y_s$  are the sun gear displacements;  $\theta_s$  is the sun gear angular displacement.  $F_{san}$  is the  $n$ th sun-planet gear mesh force, which is given as

$$F_{san} = k_{san} \delta_{san} + c_{san} \dot{\delta}_{san} \quad (4)$$

where,  $k_{san}$  and  $c_{san}$  are the  $n$ th sun-planet gear mesh stiffness and mesh damping coefficients. The detailed calculation method is given by Ref. [26]. The gear relative meshing displacements is shown in Fig. 2.  $\delta_{san}$  is the  $n$ th sun-planet gear mesh deformation, which is given as

$$\delta_{san} = (x_s - x_{an}) \cos \psi_{sn} + (y_s - y_{an}) \sin \psi_{sn} + r_s \theta_s + r_{an} \theta_{an} - r_{ac} \theta_c \cos \alpha \quad (5)$$

$$\psi_{sn} = \frac{\pi}{2} - \alpha + \psi_n \quad (6)$$

where,  $\psi_n$  is the  $n$ th planet position angle;  $r_{an}$  is the  $n$ th inner planet base radius;  $r_{ac}$  is the inner planet mounting radius;  $\theta_c$  is the carrier angular displacement;  $x_{an}$  and  $y_{an}$  are the inner planet displacements;  $\theta_{an}$  is the inner planet angular displacement;  $\alpha$  is the pressure angle of gear pair.

The dynamic equations of ring gear are expressed as

$$m_r \ddot{x}_r + c_{rx} \dot{x}_r + k_{rx} x_r + \sum F_{rbn} \cos \psi_{rn} = m_r x_r \Omega^2 + 2m_r \dot{y}_r \Omega + m_r y_r \dot{\Omega} \quad (7)$$

$$m_r \ddot{y}_r + c_{ry} \dot{y}_r + k_{ry} y_r - \sum F_{rbn} \sin \psi_{rn} = m_r y_r \Omega^2 - 2m_r \dot{x}_r \Omega - m_r x_r \dot{\Omega} \quad (8)$$

$$\frac{J_r}{r_r} \ddot{\theta}_r + \frac{c_{rt}}{r_r} \dot{\theta}_r + \frac{k_{rt}}{r_r} \theta_r - \sum F_{rbn} = 0 \quad (9)$$

where,  $r_r$  is the ring gear base radius;  $m_r$  is the ring gear mass;  $k_{rx}$  and  $c_{rx}$  are the ring support stiffness and damping coefficients in  $x$  direction;  $k_{ry}$  and  $c_{ry}$  are the ring support stiffness and damping coefficients in  $y$  direction;  $J_r$  is the ring gear moment of inertia;  $x_r$  and  $y_r$  are the ring gear displacements;  $\theta_r$  is the ring gear angular displacement.  $F_{rbn}$  is the  $n$ th ring-planet gear mesh force, which is given as

$$F_{rbn} = k_{rbn} \delta_{rbn} + c_{rbn} \dot{\delta}_{rbn} \quad (10)$$

where,  $k_{rbn}$  and  $c_{rbn}$  are the  $n$ -th ring-planet gear mesh stiffness and mesh damping coefficients.  $\delta_{rbn}$  is the  $n$ th ring-planet gear mesh deformation, which is given as

$$\delta_{rbn} = (x_r - x_{bn}) \cos \psi_{rn} + (y_{bn} - y_r) \sin \psi_{rn} + r_b \theta_{bn} - r_r \theta_{rn} + r_a \theta_c + r_r \theta_c \quad (11)$$

$$\psi_{rn} = \frac{\pi}{2} - \alpha + \psi_{bn} \quad (12)$$

$$\psi_{bn} = \arccos \frac{(Z_s + Z_a)^2 + (Z_r - Z_b)^2 - (Z_a + Z_b)^2}{2(Z_s + Z_a)(Z_r - Z_b)} + \psi_n \quad (13)$$

where,  $r_{bn}$  is the  $n$ th outer planet base radius;  $x_{bn}$  and  $y_{bn}$  are the outer planet displacements;  $\theta_{bn}$  is the outer planet angular displacement;  $Z_s$ ,  $Z_r$ ,  $Z_a$ , and  $Z_b$  are the teeth numbers of sun gear, ring gear, inner planet, and outer planet, respectively.

The dynamic equations of inner planet are expressed as

$$m_{an} \ddot{x}_{an} = m_{an} x_{an} \Omega^2 + 2m_{an} \dot{y}_{an} \Omega + m_{an} y_{an} \dot{\Omega} + m_{an} r_c \Omega^2 \cos \psi_{an} + F_{ax}^o + F_{adx}^o - F_{agx}^o + F_{san} \cos \psi_{sn} - F_{abn} \cos \psi_{abn} \quad (14)$$

$$m_{an} \ddot{y}_{an} = m_{an} y_{an} \Omega^2 - 2m_{an} \dot{x}_{an} \Omega - m_{an} x_{an} \dot{\Omega} + m_{an} r_c \Omega^2 \sin \psi_{an} + F_y^o + F_{dy}^o - F_{agy}^o + F_{san} \sin \psi_{sn} - F_{abn} \sin \psi_{abn} \quad (15)$$

$$\frac{J_{an}}{r_{an}} \ddot{\theta}_{an} + F_{san} - F_{abn} + \frac{M_{cage}}{r_{an}} + \frac{\sum f_{aoj} \frac{d_m}{2}}{r_{an}} = 0 \quad (16)$$

where,  $m_{an}$  is the  $n$ th inner planet mass;  $J_{an}$  is the  $n$ th inner planet moment of inertia;  $d_m$  is the inner planetary bearing pitch diameter;  $M_{cage}$  is the inner planet-cage friction moment;  $f_{aoj}$  is the inner planet-roller friction force;  $F_{abn}$  is the  $n$ th inner-outer planet gear mesh force;  $F_{ax}^o$  and  $F_{ay}^o$  are the resultant forces of inner planet-roller contact force and frictional force in  $x$  and  $y$  directions;  $F_{adx}^o$  and  $F_{ady}^o$  are the inner planet-roller damping forces in  $x$  and  $y$  directions;  $F_{agx}^o$  and  $F_{agy}^o$  are the resultant forces of inner planet-cage impact force and frictional force in  $x$  and  $y$  directions, which are given by

$$F_{abn} = k_{abn} \delta_{abn} + c_{abn} \dot{\delta}_{abn} \quad (17)$$

$$F_{ax}^o = \sum_{j=1}^Z (Q_{ao}(j) \cos \theta_{aj} + f_{aoj} \sin \theta_{aj}) \quad (18)$$

$$F_{ay}^o = \sum_{j=1}^Z (Q_{ao}(j) \sin \theta_{aj} + f_{aoj} \cos \theta_{aj}) \quad (19)$$

$$F_{adx}^o = \sum c_o (\dot{x}_{aej} - \dot{x}_a) \quad (20)$$

$$F_{ady}^o = \sum c_o (\dot{y}_{aej} - \dot{y}_a) \quad (21)$$

where,  $k_{abn}$  and  $c_{abn}$  are the  $n$ th inner-outer planet gear mesh stiffness and mesh damping coefficients;  $c_o$  is the damping coefficient in the roller-planet contact;  $x_{ej}$  and  $y_{ej}$  are the  $j$ th planetary bearing roller displacements;  $\phi_j$  is the  $j$ th roller position angle; the subscript  $a$  represents inner planetary bearing;  $Q_{ao}$  is the inner planet-roller contact force;  $f_{ao}$  is the inner planet-roller friction force;  $\delta_{abn}$  is the mesh deformation of  $n$ th inner-outer planet gear mesh, which are given by

$$\delta_{abn} = (x_{an} - x_{bn}) \cos \psi_{abn} + (y_{an} - y_{bn}) \sin \psi_{abn} - r_a \theta_{an} - r_b \theta_{bn} \quad (22)$$

$$\psi_{abn} = -\frac{\pi}{2} + \alpha + \psi_{an} \quad (23)$$

$$\psi_{an} = \arccos \frac{(Z_r - Z_b) \cos \psi_{bn} - (Z_s + Z_a) \cos \psi_n}{Z_a + Z_b} \quad (24)$$

$$Q_{ao} = \zeta_{ao} K_o \delta_{ao}^{10} \quad (25)$$

$$f_{ao} = \mu Q_{ao} \quad (26)$$

where,  $\zeta_{ao}$  is the inner planet-roller contact factor,  $\zeta_{ao}=1$  when  $\delta_{ao}>0$ ,  $\zeta_{ao}=0$  when  $\delta_{ao}\leq 0$ ;  $\mu$  is the friction coefficient, the lubrication condition is considered to be excellent; and the lubrication state is considered to be elastohydrodynamic lubrication in this work. Thus, the friction coefficient can be considered as a constant value and its value is 0.007 in the work;  $\delta_{ao}$  is the inner planet-roller contact deformation, which is given by

$$\delta_{ao} = (x_{aej} - x_a) \cos \theta_{aj} + (y_{aej} - y_a) \sin \theta_{aj} \quad (27)$$

The dynamic equations of outer planet are expressed as

$$m_{bn}\ddot{x}_{bn} = m_{bn}x_{bn}\Omega^2 + 2m_{bn}\dot{y}_{bn}\Omega + m_{bn}y_{bn}\dot{\Omega} + m_{bn}r_c\Omega^2 \cos\psi_{bn} + F_{bx}^o + F_{bdx}^o - F_{bgx}^o + F_{rbn} \cos\psi_{rn} + F_{abn} \cos\psi_{abn} \quad (28)$$

$$m_{bn}\ddot{y}_{bn} = m_{bn}y_{bn}\Omega^2 - 2m_{bn}\dot{x}_{bn}\Omega - m_{bn}x_{bn}\dot{\Omega} + m_{bn}r_c\Omega^2 \sin\psi_{bn} + F_{by}^o + F_{bdy}^o - F_{bgy}^o - F_{rbn} \sin\psi_{rn} + F_{abn} \sin\psi_{abn} \quad (29)$$

$$\frac{J_{bn}}{r_{bn}}\ddot{\theta}_{bn} + F_{ibn} - F_{abn} + M_{cage} + \sum f_{boj} \frac{d_m}{2} = 0 \quad (30)$$

where,  $m_{bn}$  is the  $n$ th outer planet mass;  $J_{bn}$  is the  $n$ th outer planet moment of inertia;  $x_{bn}$  and  $y_{bn}$  are the  $n$ th outer planet displacements;  $\theta_{bn}$  is  $n$ th outer planet angular displacement;  $F_{bx}^o$  and  $F_{by}^o$  are the resultant forces of outer planet-roller contact force and frictional force in  $x$  and  $y$  directions;  $F_{bdx}^o$  and  $F_{bdy}^o$  are the outer planet-roller damping forces in  $x$  and  $y$  directions;  $F_{bgx}^o$  and  $F_{bgy}^o$  are the resultant forces of outer planet-cage impact force and frictional force in  $x$  and  $y$  directions; the subscript  $b$  represents the outer planetary bearing; the planetary bearing cage is guided by the outer ring and the cage force analysis is shown in Fig. 3. The calculation method of the above variates is given by

$$F_{bx}^o = \sum_{j=1}^z (Q_{bo}(j) \cos\theta_{bj} + f_{boj} \sin\theta_{bj}) \quad (31)$$

$$F_{by}^o = \sum_{j=1}^z (Q_{bo}(j) \sin\theta_{bj} + f_{boj} \cos\theta_{bj}) \quad (32)$$

$$F_{bdx}^o = \sum c_o (\dot{x}_{bcj} - \dot{x}_b) \quad (33)$$

$$F_{bdy}^o = \sum c_o (\dot{y}_{bcj} - \dot{y}_b) \quad (34)$$

$$\begin{bmatrix} F_{gx} \\ F_{gy} \\ M_{cage} \end{bmatrix} = \begin{bmatrix} -\cos\phi_g & -\sin\phi_g & 0 \\ -\sin\phi_g & \cos\phi_g & 0 \\ 0 & 0 & 1 \end{bmatrix} \begin{bmatrix} F_g^n \\ F_g^r \\ M_g \end{bmatrix} \quad (35)$$

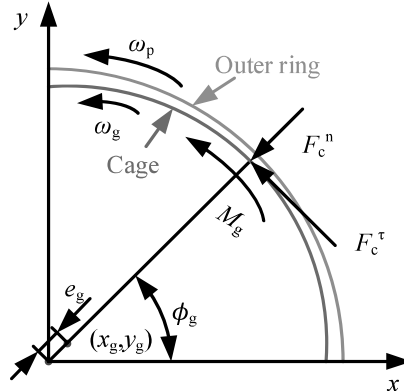


Fig. 3. The force between the cage and guide surface.

where,  $Q_{bo}$  is the outer planet-roller contact force;  $f_{bo}$  is the outer planet-roller friction force;  $F_g^n$  is the outer planet-cage impact force;  $c_o$  is the damping coefficient in the roller-planet contact;  $F_g^r$  is the outer planet-cage friction force;  $M_g$  is the outer planet-cage friction moment, which are given by

$$Q_{bo} = \zeta_{bo} K_o \delta_{bo}^9 \quad (36)$$

$$f_{ao} = \mu Q_{ao} \quad (37)$$

$$F_g^n = K_g (e_g - C_g) + c_g \dot{e}_g \quad (38)$$

$$F_g^r = \mu F_g^n \quad (39)$$

$$M_g = F_g^r R_{cd} \quad (40)$$

where,  $\zeta_{bo}$  is the outer planet-roller contact factor,  $\zeta_{bo}=1$  when  $\delta_{bo}>0$ ,  $\zeta_{bo}=0$  when  $\delta_{bo}\leq 0$ ;  $\delta_{bo}$  is the outer planet-roller contact deformation;  $K_o$  is the contact stiffness between the roller and outer ring (planet); and  $e_g$  is the cage center offset, which are given by

$$\delta_{bo} = (x_{bej} - x_b) \cos \phi_{bj} + (y_{bej} - y_b) \sin \phi_{bj} \quad (41)$$

$$e_g = \sqrt{x_g^2 + y_g^2} \quad (42)$$

The dynamic equations of carrier are expressed as

$$m_c \ddot{x}_c = m_c x_c \Omega^2 + 2m_c \dot{y}_c \Omega + m_c y_c \dot{\Omega} - F_{ax}^i - F_{dax}^i - F_{bx}^i - F_{dbx}^i - c_{cx} \dot{x}_c - k_{cx} x_c \quad (43)$$

$$m_c \ddot{y}_c = m_c y_c \Omega^2 - 2m_c \dot{x}_c \Omega - m_c x_c \dot{\Omega} - F_{ay}^i - F_{day}^i - F_{by}^i - F_{dby}^i - c_{cy} \dot{y}_c - k_{cy} y_c \quad (44)$$

$$J_c \ddot{\theta}_c + F_{ax}^i \sin \psi_{an} r_{ac} + F_{bx}^i \sin \psi_{bn} r_{bc} - F_{ay}^i \cos \psi_{an} r_{ac} - F_{by}^i \cos \psi_{bn} r_{bc} = T_0 \quad (45)$$

where,  $m_c$  is the carrier mass;  $J_c$  is the carrier moment of inertia;  $x_c$  and  $y_c$  are the carrier displacements;  $\theta_c$  is the carrier angular displacement;  $F_{a/bx}^i$  and  $F_{a/by}^i$  are the resultant forces of inner/outer planetary bearing roller-carrier contact force and frictional force in  $x$  and  $y$  directions;  $F_{a/bdx}^i$  and  $F_{a/bdy}^i$  are the inner/outer planetary bearing roller-carrier damping forces in  $x$  and  $y$  directions, which are given by

$$F_{a/bx}^i = \sum_{j=1}^Z (Q_{a/bi}(j) \cos \theta_{a/bj} + f_{a/bij} \sin \theta_{a/bj}) \quad (46)$$

$$F_{a/by}^i = \sum_{j=1}^Z (Q_{a/bi}(j) \sin \theta_{a/bj} + f_{a/bij} \cos \theta_{a/bj}) \quad (47)$$

$$F_{a/bdx}^i = \sum c_i (\dot{x}_c - \dot{x}_{a/bej}) \quad (48)$$

$$F_{a/bdy}^i = \sum c_i (\dot{y}_c - \dot{y}_{a/bej}) \quad (49)$$

where,  $c_i$  is the damping coefficient in roller-carrier contact;  $Q_{a/bi}$  is the inner/outer planetary bearing roller-carrier contact force;  $f_{bo}$  is the inner/outer planetary bearing roller-carrier friction force, which are given by

$$Q_{a/bi} = \zeta_{a/bi} K_i \delta_{a/bi}^{\frac{10}{9}} \quad (50)$$

$$f_{a/bi} = \mu Q_{a/bi} \quad (51)$$

where,  $\zeta_{a/bi}$  is the inner/outer planetary bearing roller-carrier contact factor,  $\zeta_{a/bi}=1$  when  $\delta_{a/bi}>0$ ,  $\zeta_{a/bi}=0$  when  $\delta_{a/bi}\leq 0$ ;  $K_i$  is the contact stiffness between the roller and inner ring (carrier);  $\delta_{a/bi}$  is the inner/outer planetary bearing roller-carrier contact deformation, which is given by

$$\delta_{a/bi} = (x_c - x_{a/be}) \cos \theta_{a/bj} + (y_c - y_{a/be}) \sin \theta_{a/bj} - 0.5C_r \quad (52)$$

where  $C_r$  is the planetary bearing clearance.



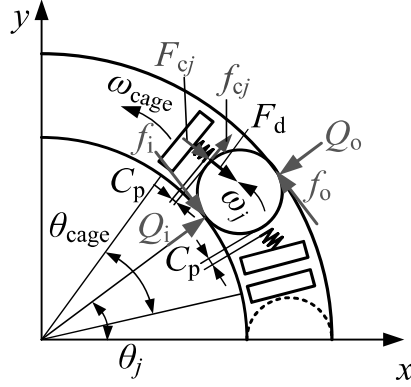


Fig. 4. Force analysis of rollers.

The roller force analysis is given in Fig. 4, the dynamic equations of planetary bearing roller are expressed as

$$m_e \ddot{x}_{ej} = -Q_o(j) \cos \theta_j + Q_i(j) \cos \theta_j - c_o(\dot{x}_{ej} - \dot{x}_p) + c_i(\dot{x}_c - \dot{x}_{ej}) + F_{cj} \sin \theta_j - f_{cj} \sin \theta_j - f_{oj} \sin \theta_j + f_{ij} \sin \theta_j + F_{c1} \cos \phi_c + F_{c2} \cos \theta_j + F_d \sin \theta_j \quad (53)$$

$$m_e \ddot{y}_{ej} = -Q_o(j) \sin \theta_j + Q_i(j) \sin \theta_j - c_o(\dot{y}_{ej} - \dot{y}_p) + c_i(\dot{y}_c - \dot{y}_{ej}) - F_{cj} \cos \theta_j + f_{cj} \sin \theta_j + f_{oj} \cos \theta_j - f_{ij} \cos \theta_j + F_{c1} \sin \phi_c + F_{c2} \sin \theta_j - F_d \cos \theta_j \quad (54)$$

$$I_e \ddot{\phi}_{ej} = (F_r^o + F_r^i - f_{cj}) \frac{D_w}{2} \quad (55)$$

$$I_{ec} \ddot{\theta}_{ej} = F_r^o R_o - F_r^i R_i - F_{cj} \frac{d_m}{2} - F_d \frac{d_m}{2} \quad (56)$$

where,  $m_e$  is the planetary bearing roller mass;  $I_e$  is the planetary bearing roller moment of inertia around roller axis;  $I_{ec}$  is the planetary bearing roller moment of inertia about the planetary bearing axis;  $x_e$  and  $y_e$  are the roller displacements;  $\phi_e$  and  $\theta_e$  are the roller angular displacements around its axis and planetary bearing axis;  $F_{c1}$  and  $F_{c2}$  are the centrifugal forces rotating about the planetary bearing axis and planetary gear set axis as given in Fig. 5. The calculation method is given by

$$F_{c1} = m_e \dot{\theta}_c^2 \overline{oo_r} \quad (57)$$

$$F_{c2} = m_e \dot{\theta}_{ej}^2 \frac{d_m}{2} \quad (58)$$

where,  $\overline{oo_r}$  is the planetary gear set-roller center distance, which is given by

$$\overline{oo_r} = \sqrt{\left( \frac{d_m}{2} \sin \theta_{ej} + \frac{m(Z_s + Z_p)}{2} \sin \varphi_p \right)^2 + \left( \frac{d_m}{2} \cos \theta_{ej} + \frac{m(Z_s + Z_p)}{2} \cos \varphi_p \right)^2} \quad (59)$$

$$\varphi_c = \arctan \left( \frac{d_m \cos \theta_{ej} + m(Z_s + Z_p) \cos \varphi_p}{d_m \sin \theta_{ej} + m(Z_s + Z_p) \sin \varphi_p} \right) \quad (60)$$

$$\varphi_p = \theta_c + \frac{2(n-1)\pi}{n} \quad (61)$$

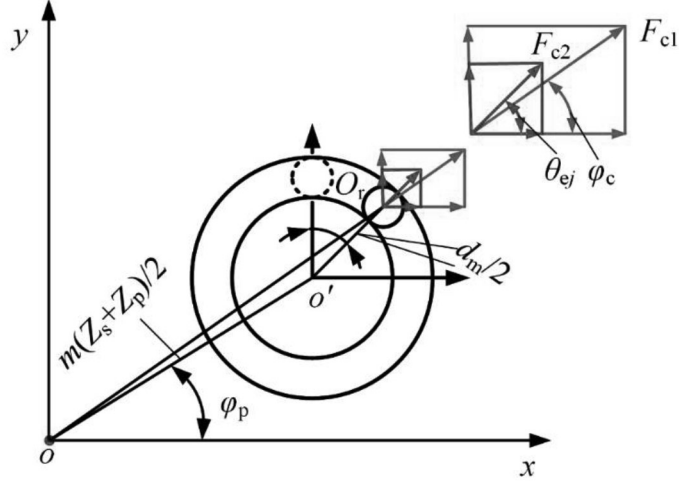


Fig. 5. Roller centrifugal forces.

The dynamic equations of planetary bearing cage are expressed as

$$m_g \ddot{x}_g = -F_{cj} \sin \theta_j - f_{cj} \cos \theta_j + F_{cx} \quad (62)$$

$$m_g \ddot{y}_g = F_{cj} \cos \theta_j - f_{cj} \sin \theta_j + F_{cy} - G \quad (63)$$

$$I_g \ddot{\theta}_g = M_{cage} + F_{cj} \frac{d_m}{2} \quad (64)$$

where,  $m_g$  is the cage mass;  $I_g$  is the cage moment of inertia around its own axis;  $x_g$  and  $y_g$  are the cage displacements;  $\theta_g$  is the angular displacement around its axis;  $G$  is the cage gravity;  $F_{cj}$  is the roller-cage impact force;  $f_{cj}$  is the roller-cage friction force, which are given by

$$\begin{cases} F_{cj} = K_{cage} \left[ \left( \theta_j - \theta_{cage} \right) \frac{d_m}{2} - C_p \right] & \left( \theta_j - \theta_{cage} \right) > 0 \\ F_{cj} = K_{cage} \left[ \left( \theta_j - \theta_{cage} \right) \frac{d_m}{2} + C_p \right] & \left( \theta_j - \theta_{cage} \right) < 0 \end{cases} \quad (65)$$

$$f_{cj} = \mu F_{cj} \quad (66)$$

where,  $K_{cage}$  is the pocket contact stiffness between the roller and cage, and  $C_p$  is the cage pocket clearance.

#### 4. Results and discussion

In this section, a DPGS is used to analyze the vibrations of planetary bearing. The gear parameters of DPGS are shown in Table 1. The planetary bearing parameters of DPGS are shown in Table 2. The carrier rotation speed is 400r/min in the following analysis.

Table 1. Gear parameters of DPGS.

	Sun gear	Ring gear	Inner planet	Outer planet	Carrier
Teeth number	34	82	21	22	-
Module (mm)	4	4	4	4	-
Pressure angle (°)	25	25	25	25	-
Face width (mm)	38	38	38	8	-
Mass (kg)	2.27	7.008	1.196	1.36	7.008
Bearing stiffness (N/m)	$1 \times 10^8$	$1 \times 10^8$	-	-	$1 \times 10^8$

Table 2. Planetary bearing parameters of DPGS.

Parameters	Value	Parameters	Value
Roller number	22	Roller length(mm)	28.5
Inner ring diameter(mm)	38	Clearance( $\mu$ m)	10

Outer ring diameter(mm)	46	Pocket clearance( $\mu\text{m}$ )	20
Roller diameter(mm)	4	Cage guidance clearance(mm)	0.52

#### 4.1 Effect of the sun gear rotation speed on the planetary bearing contact force

The dynamic model is solved by the fourth-order Runge-Kutta method. The initial displacements are  $10^{-6}$  m. The initial velocities are 0 m/s. The initial angular displacements are 0 rad. The initial roller angular velocities around its axis are set to 0.9 times of that in the case of pure rolling. The other angular velocities are set to the values in the case of pure rolling.

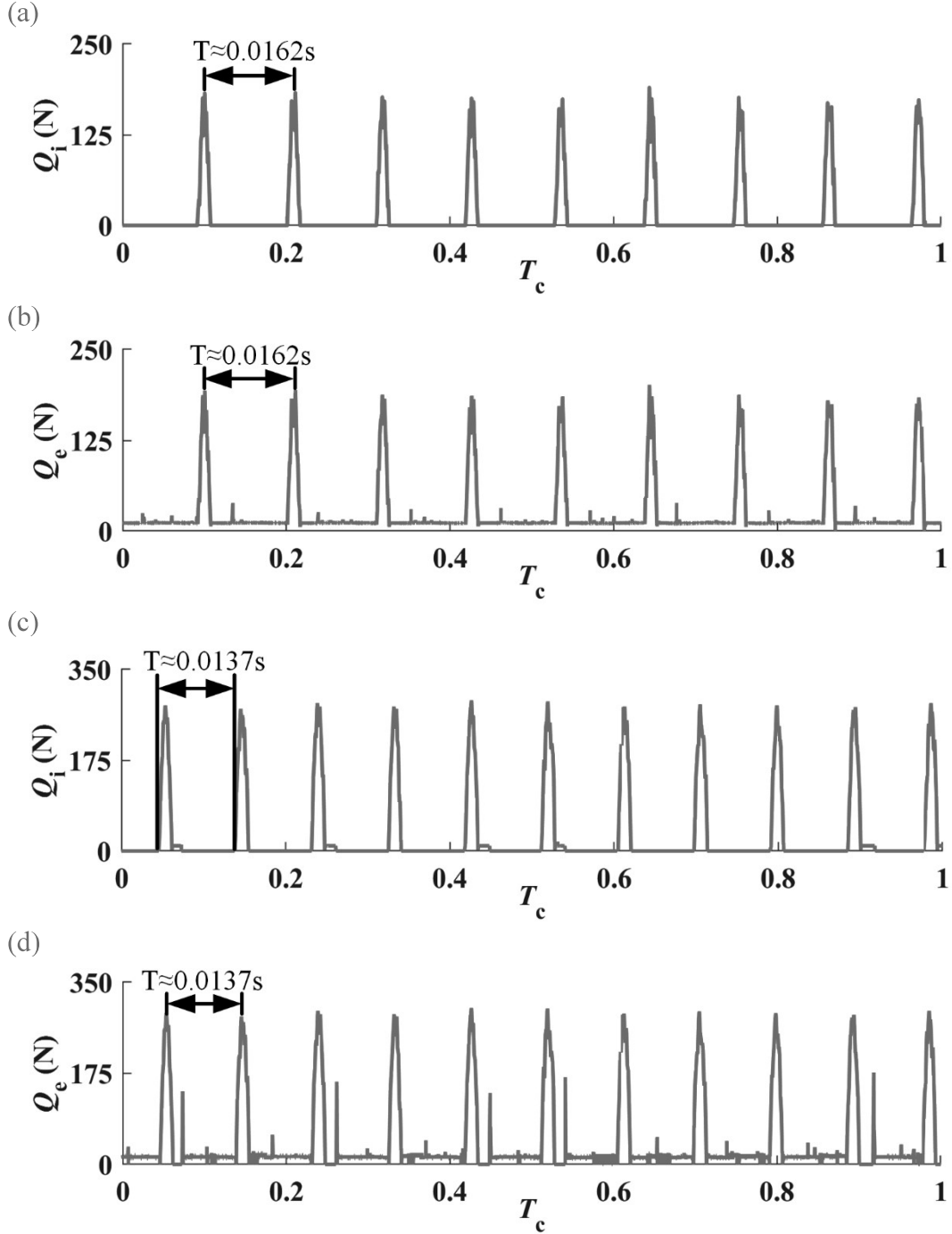


Fig. 6. (a) The inner ring-roller contact force of inner planetary bearing; (b) the outer ring-roller contact force of inner planetary bearing; (c) the inner ring-roller contact force of outer planetary bearing contact force; and (d) the outer ring-roller contact force of outer planetary bearing contact force.

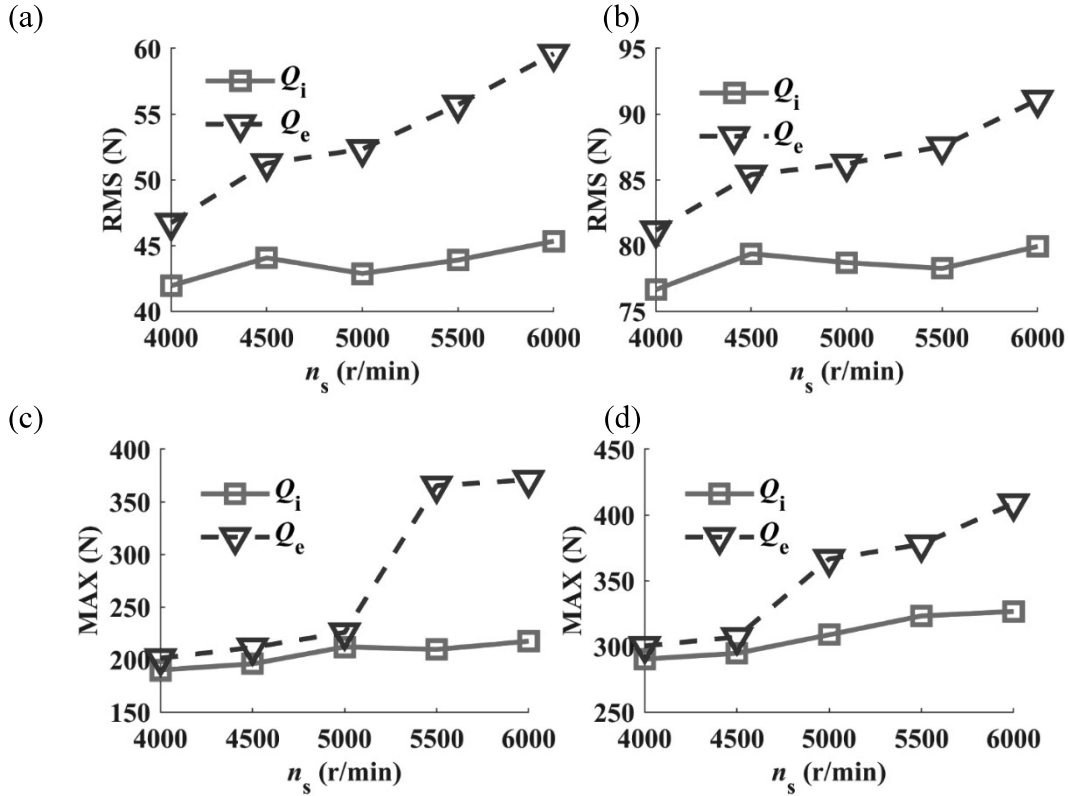


Fig. 7. Effect of the sun gear rotation speed on the (a) RMS value of inner planetary bearing contact force; (b) RMS value of inner planetary bearing contact force; (c) MAX value of outer planetary bearing contact force; and (d) MAX value of outer planetary bearing contact force.

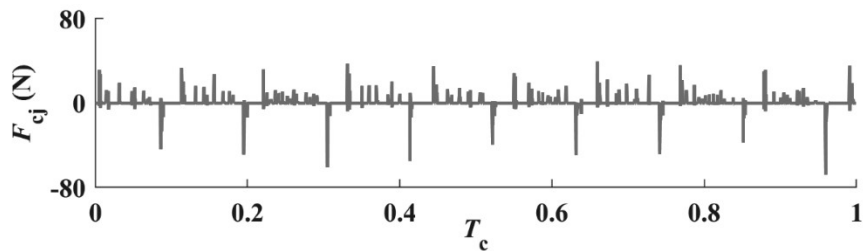
Fig. 6 gives the time-domain planetary bearing contact force when the sun gear rotation speed is 4000rpm. The inner/outer planetary bearing contact force exhibit obvious periodic variations. The period of inner planetary bearing contact force is about 0.0162s and that of outer planetary bearing contact force has a period of 0.0137s. Thus, the periodic frequency of inner planetary bearing contact force is about 61.73Hz and that of outer planetary bearing contact force is about 71.72Hz. The inner planetary bearing cage rotation period is 0.0146s without the roller slip; and the outer planetary bearing cage rotation period is 0.0153s without the roller slip. Thus, the periodic frequency of inner planetary bearing cage rotation period ( $f_{ga}$ ) is about 68.49Hz; and that of outer planetary bearing cage rotation period ( $f_{gb}$ ) is about 65.36Hz. The carrier rotation periodic frequency ( $f_c$ ) is about 6.67 Hz. The contact force periodic frequency of inner planetary bearing approximately equal to  $f_{ga}-f_c$ ; and that of outer planetary bearing approximately equal to  $f_{gb}+f_c$ . This is because the inner planetary bearing cage and the carrier rotate in the same direction, while the outer planetary bearing cage and carrier rotate in the opposite direction. The effect of the sun gear rotation speed on the planetary bearing contact force is given in Fig. 7. In Fig. 7, the inner/outer planetary bearing contact force root mean square (RMS) and maximum (MAX) values of different sun gear rotation speeds are given. The RMS and MAX values of inner/outer planetary bearing contact force generally increase with the increment of sun gear rotation speed. Figure 7(a) illustrates the effect of the sun gear rotation speed on the RMS value of inner planetary bearing contact force. When the sun gear rotation speed is from 4000rpm to 6000rpm, the RMS value of inner planetary bearing inner ring-roller contact force is from 41.95N to 45.33N; and the RMS value of inner planetary bearing outer ring-roller contact force is from 46.73N to 59.59N. Figure 7(b) illustrates the effect of the sun gear rotation speed on the RMS value of outer planetary bearing contact force. When the sun gear rotation speed is from 4000rpm to 6000rpm, the RMS value of outer planetary bearing

inner ring-roller contact force is from 76.65N to 79.93N, and the RMS value of outer planetary bearing outer ring-roller contact force is from 81.16N to 91.08N. Figure 7(c) illustrates the effect of the sun gear rotation speed on the MAX value of inner planetary bearing contact force. When the sun gear rotation speed is from 4000rpm to 6000rpm, the MAX value of inner planetary bearing inner ring-roller contact force is from 190.13N to 217.42N; and the MAX value of inner planetary bearing outer ring-roller contact force is from 201.33N to 370.85N. Figure 7(d) illustrates the effect of the sun gear rotation speed on the MAX value of outer planetary bearing contact force. When the sun gear rotation speed is from 4000rpm to 6000rpm, the MAX value of outer planetary bearing inner ring-roller contact force is from 290.36N to 326.59N, and the MAX value of outer planetary bearing outer ring-roller contact force is from 300.42N to 408.8N.

#### 4.2 Effect of the sun gear rotation speed on the planetary bearing roller-cage impact force

Fig. 8 gives the time-domain planetary bearing roller-cage impact force when the sun gear rotation speed is 4000rpm. The inner/outer planetary bearing roller-cage impact force exhibit obvious periodic variations. The effect of the sun gear rotation speed on the planetary bearing roller-cage impact force is given in Fig. 9. In Fig. 9, the RMS and MAX values of planetary bearing roller-cage impact force for different sun gear rotation speeds are given. The RMS and MAX values of inner/outer planetary bearing roller-cage impact force generally increase with the increment of sun gear rotation speed. Fig. 9(a) illustrates the effect of the sun gear rotation speed on the RMS value of inner planetary bearing roller-cage impact force. When the sun gear rotation speed is from 4000rpm to 6000rpm, the RMS value of inner planetary bearing roller-cage impact force is from 1.82N to 2.86N. Fig. 9(b) illustrates the effect of the sun gear rotation speed on the RMS value of outer planetary bearing roller-cage impact force. When the sun gear rotation speed is from 4000rpm to 6000rpm, the RMS value of outer planetary bearing roller-cage impact force is from 1.39N to 2.97N. Fig. 9(c) illustrates the effect of the sun gear rotation speed on the MAX value of inner planetary bearing roller-cage impact force. When the sun gear rotation speed is from 4000rpm to 6000rpm, the MAX value of inner planetary bearing roller-cage impact force is from 61.89N to 98.02N. Fig. 9(d) illustrates the effect of the sun gear rotation speed on the MAX value of outer planetary bearing roller-cage impact force. When the sun gear rotation speed is from 4000rpm to 6000rpm, the MAX value of outer planetary bearing roller-cage impact force is from 56.9N to 88.45N.

(a)



(b)

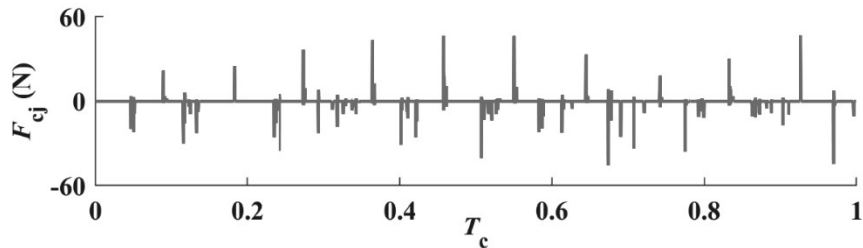


Fig. 8. Time-domain (a) inner planetary bearing roller-cage impact force; and (b) outer planetary bearing roller-cage impact force.

(a)

(b)

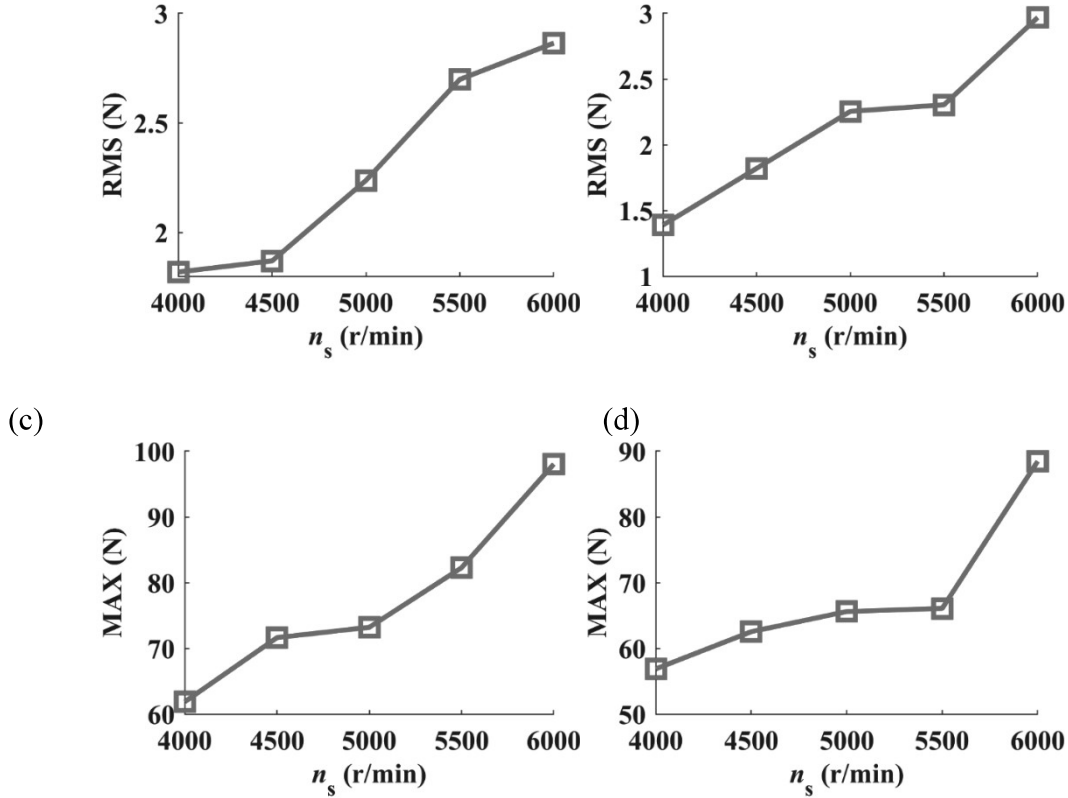


Fig. 9. Effect of the sun gear rotation speed on the (a) RMS value of inner planetary bearing roller-cage impact force; (b) RMS value of inner planetary bearing roller-cage impact force; (c) MAX value of outer planetary bearing roller-cage impact force; and (d) MAX value of outer planetary bearing roller-cage impact force.

#### 4.2 Effect of the sun gear rotation speed on the planetary bearing vibrations

Fig. 10 gives the time-domain planetary bearing roller acceleration when the sun gear rotation speed is 4000rpm. The inner/outer planetary bearing roller acceleration exhibit obvious periodic variations. The effect of the sun gear rotation speed on the planetary bearing roller acceleration is shown in Fig. 11. In Fig. 11, the RMS and MAX values of inner/outer planetary bearing roller acceleration for different sun gear rotation speeds are given. The RMS and MAX values of inner/outer planetary bearing roller acceleration generally increase with the increment of sun gear rotation speed. Fig. 11(a) illustrates the effect of the sun gear rotation speed on the RMS value of inner planetary bearing roller acceleration. When the sun gear rotation speed is from 4000rpm to 6000rpm, the RMS value of inner planetary bearing roller acceleration in  $x$  direction is from  $1970.53\text{m/s}^2$  to  $3253.35\text{m/s}^2$ ; and the RMS value of inner planetary bearing roller acceleration in  $y$  direction is from  $2012.65\text{m/s}^2$  to  $3203.17\text{m/s}^2$ . Fig. 11(b) illustrates the effect of the sun gear rotation speed on the RMS value of outer planetary bearing roller acceleration. When the sun gear rotation speed is from 4000rpm to 6000rpm, the RMS value of outer planetary bearing roller acceleration in  $x$  direction is from  $3978.8\text{m/s}^2$  to  $4407.24\text{m/s}^2$ ; and the RMS value of outer planetary bearing roller acceleration in  $y$  direction is from  $3953.7\text{m/s}^2$  to  $4716.4\text{m/s}^2$ . Fig. 11(c) illustrates the effect of the sun gear rotation speed on the MAX value of inner planetary bearing roller acceleration. When the sun gear rotation speed is from 4000rpm to 6000rpm, the MAX value of inner planetary bearing roller acceleration in  $x$  direction is from  $1.416 \times 10^4\text{m/s}^2$  to  $1.231 \times 10^5\text{m/s}^2$ ; and the MAX value of inner planetary bearing roller acceleration in  $y$  direction is from  $1.285 \times 10^4\text{m/s}^2$  to  $1.05 \times 10^5\text{m/s}^2$ . Fig. 11(d) illustrates the effect of the sun gear rotation speed on the MAX value of outer planetary bearing roller acceleration. When the sun gear rotation speed is from 4000rpm to 6000rpm, the MAX

value of outer planetary bearing roller acceleration in  $x$  direction is from  $4.489 \times 10^4 \text{m/s}^2$  to  $1.179 \times 10^5 \text{m/s}^2$ ; and the MAX value of outer planetary bearing roller acceleration in  $y$  direction is from  $5.779 \times 10^4 \text{m/s}^2$  to  $1.176 \times 10^5 \text{m/s}^2$ .

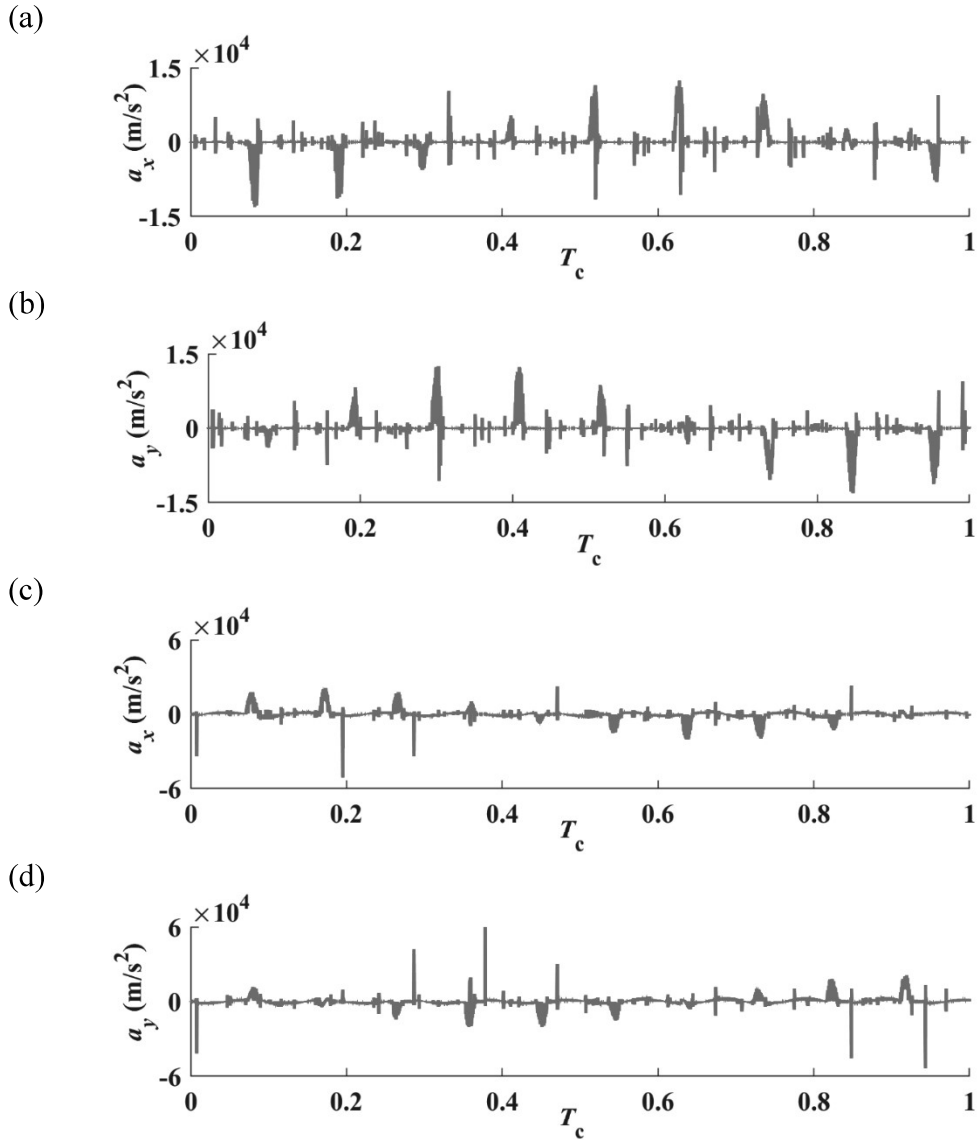
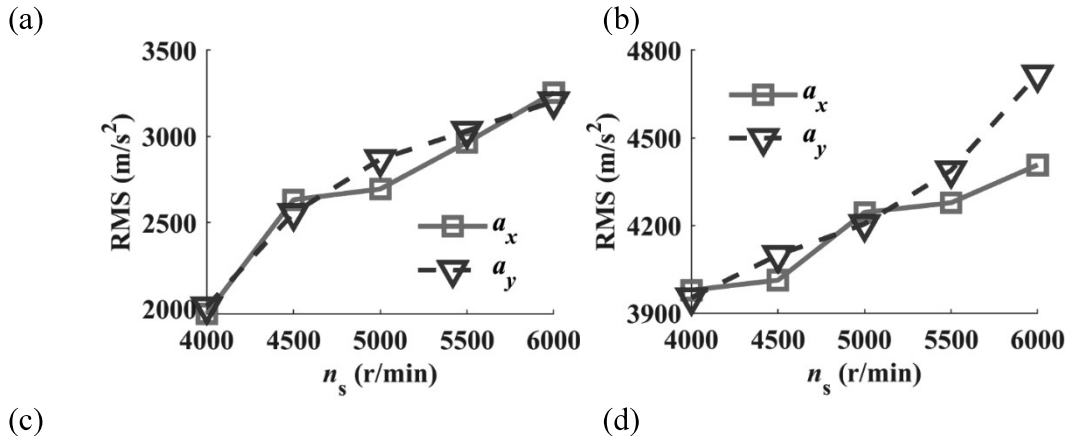


Fig. 10. Time-domain (a) inner planetary bearing roller vibrations in  $x$  direction; (b) inner planetary bearing roller vibrations in  $y$  direction; (c) outer planetary bearing roller vibrations in  $x$  direction; and (d) outer planetary bearing roller vibrations in  $y$  direction.



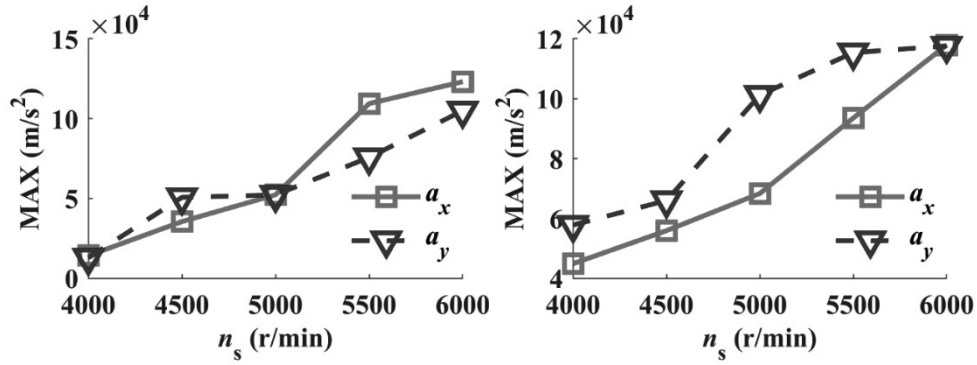
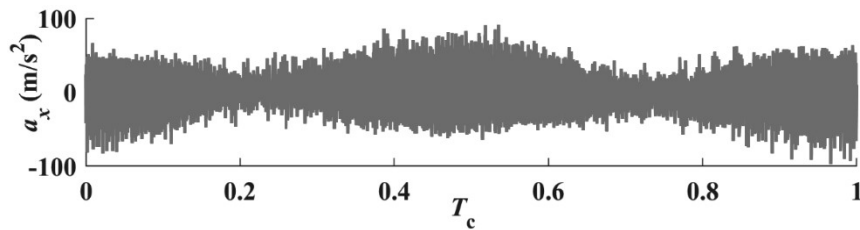
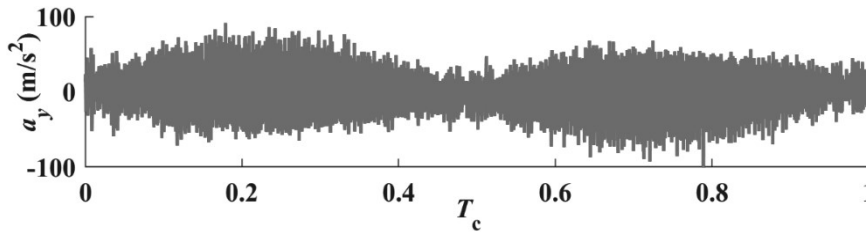


Fig. 11. Effect of the sun gear rotation speed on the (a) RMS value of inner planetary bearing roller; (b) RMS value of outer planetary bearing roller; (c) MAX value of inner planetary bearing roller; and (d) MAX value of outer planetary bearing roller.

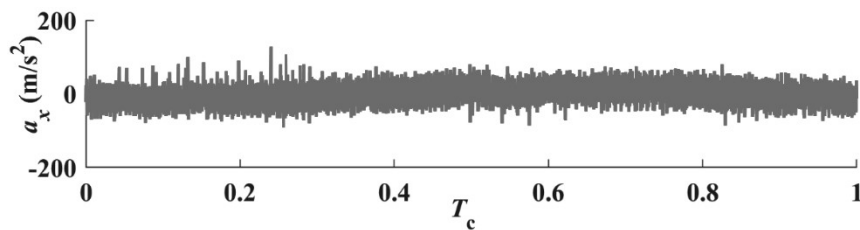
(a)



(b)



(c)



(d)

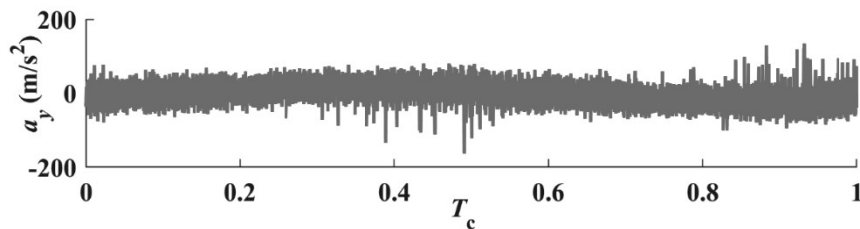


Fig. 12. Time-domain (a) inner planet vibrations in  $x$  direction; (b) inner planet vibrations in  $y$  direction; (c) outer planet vibrations in  $x$  direction; and (d) outer planet vibrations in  $y$  direction.

Fig. 12 gives the time-domain planet acceleration when the sun gear rotation speed is 4000rpm. The inner/outer planet acceleration exhibit obvious periodic variations. The effect of sun gear rotation speed on the planet acceleration is given in Fig. 13. In Fig. 13, the RMS and MAX values of inner/outer planet acceleration for different sun gear rotation speeds are given. The RMS and MAX values of inner/outer planet acceleration generally increase with the increment of sun gear rotation speed. Fig. 13(a) illustrates the effect of the sun gear rotation



speed on the RMS value of inner planet acceleration. When the sun gear rotation speed is from 4000rpm to 6000rpm, the RMS value of inner planet acceleration in  $x$  direction is from  $25.73\text{m/s}^2$  to  $43.82\text{m/s}^2$ ; and the RMS value of inner planet acceleration in  $y$  direction is from  $25.57\text{m/s}^2$  to  $43.88\text{m/s}^2$ . Fig. 13(b) illustrates the effect of the sun gear rotation speed on the RMS value of outer planet acceleration. When the sun gear rotation speed is from 4000rpm to 6000rpm, the RMS value of outer planet acceleration in  $x$  direction is from  $24.99\text{m/s}^2$  to  $42.21\text{m/s}^2$ ; and the RMS value of outer planet acceleration in  $y$  direction is from  $26.55\text{m/s}^2$  to  $42.86\text{m/s}^2$ . Fig. 13(c) illustrates the effect of sun gear rotation speed on the MAX value of inner planet acceleration. When the sun gear rotation speed is from 4000rpm to 6000rpm, the inner planet acceleration MAX value in  $x$  direction is from  $97.04\text{m/s}^2$  to  $243.53\text{m/s}^2$ ; and the MAX value of inner planet acceleration in  $y$  direction is from  $104.23\text{m/s}^2$  to  $283.67\text{m/s}^2$ . Fig. 13(d) illustrates the effect of the sun gear rotation speed on the MAX value of outer planet acceleration. When the sun gear rotation speed is from 4000rpm to 6000rpm, the MAX value of outer planet acceleration in  $x$  direction is from  $128.24\text{m/s}^2$  to  $316.15\text{m/s}^2$ ; and the MAX value of outer planet acceleration in  $y$  direction is from  $162.69\text{m/s}^2$  to  $354.89\text{m/s}^2$ .

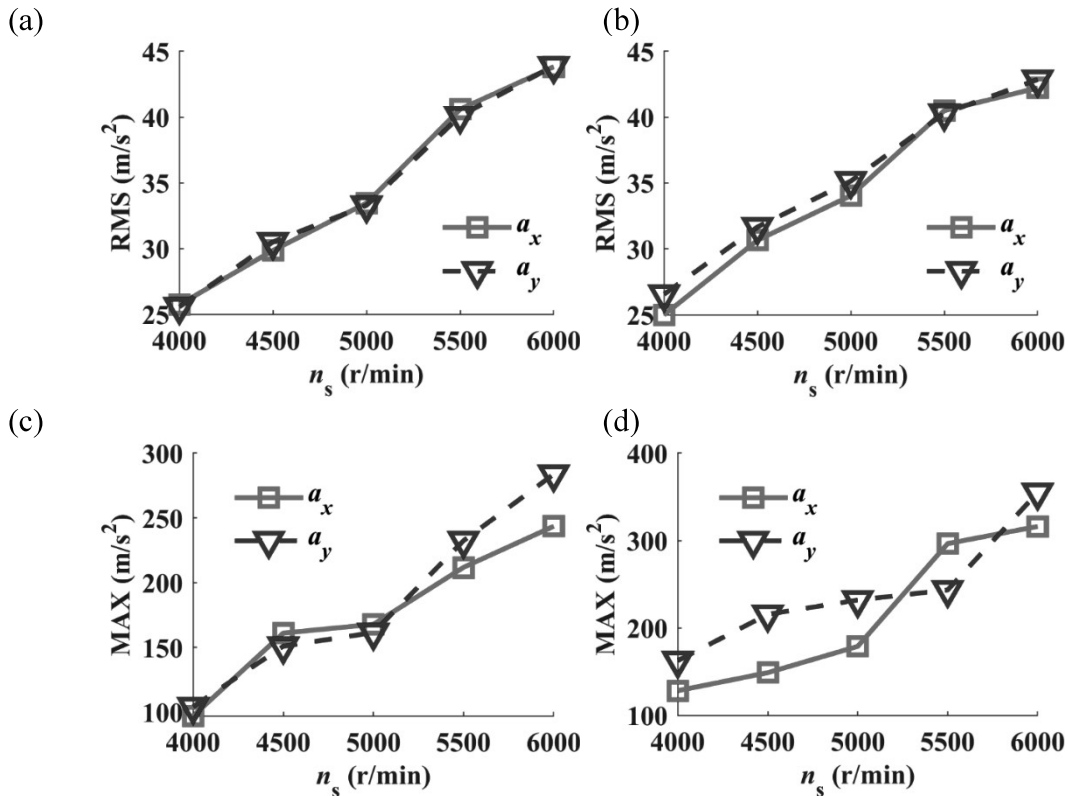


Fig. 13. Effect of the sun gear rotation speed on the (a) RMS value of inner planet; (b) RMS value of outer planet; (c) MAX value of inner planet; and (d) MAX value of outer planet.

Fig. 14 gives the time-domain carrier acceleration when the sun gear rotation speed is 4000rpm. The effect of the sun gear rotation speed on the planet acceleration is given in Fig. 15. In Fig. 15, the RMS and MAX values of carrier acceleration for different sun gear rotation speeds are given. The RMS and MAX values of carrier acceleration generally increase with the increment of sun gear rotation speed. Fig. 15(a) illustrates the effect of the sun gear rotation speed on the RMS value of carrier acceleration. When the sun gear rotation speed is from 4000rpm to 6000rpm, the RMS value of inner planet acceleration in  $x$  direction is from  $4.02\text{m/s}^2$  to  $7.08\text{m/s}^2$ ; and the RMS value of inner planet acceleration in  $y$  direction is from  $4.12\text{m/s}^2$  to  $7.14\text{m/s}^2$ . Fig. 15(b) illustrates the effect of the sun gear rotation speed on the MAX value of carrier acceleration. When the sun gear rotation speed is from 4000rpm to 6000rpm, the MAX

value of carrier acceleration in  $x$  direction is from  $13.56\text{m/s}^2$  to  $27.71\text{m/s}^2$ ; and the MAX value of carrier acceleration in  $y$  direction is from  $15.12\text{m/s}^2$  to  $27.22\text{m/s}^2$ .

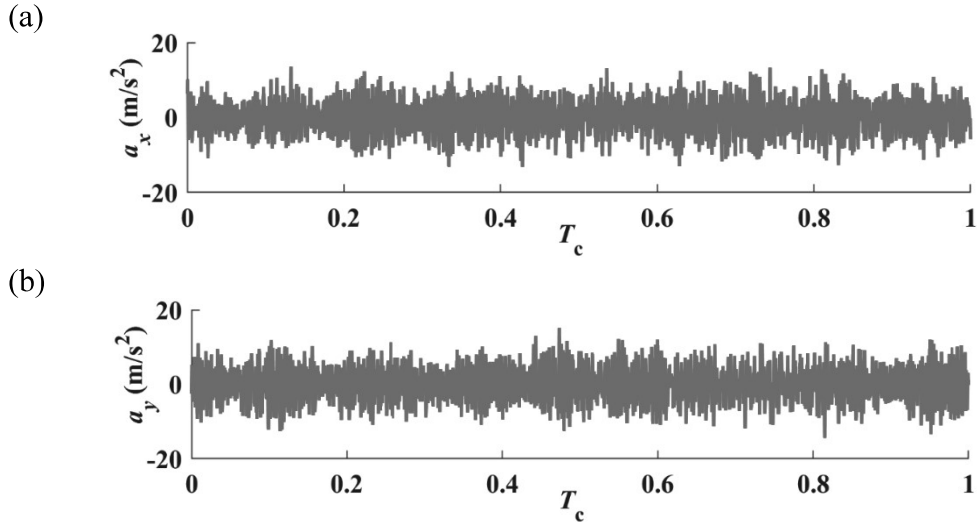


Fig. 14. Time-domain vibrations of carrier in (a)  $x$  and (b)  $y$  directions.

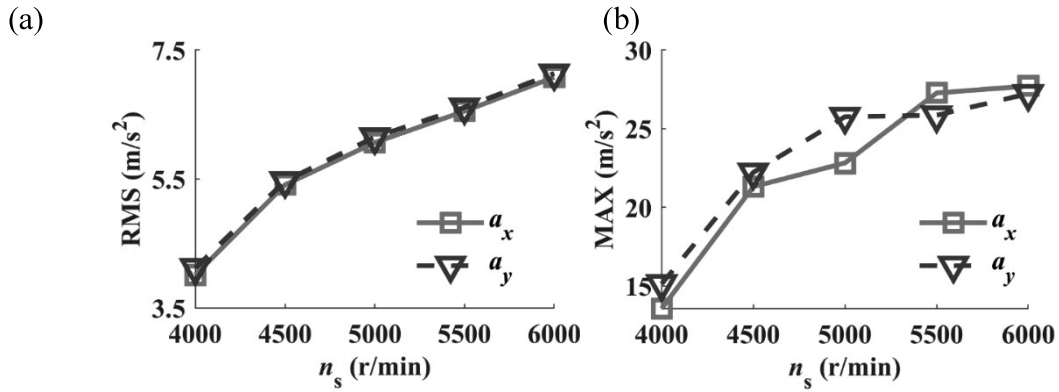
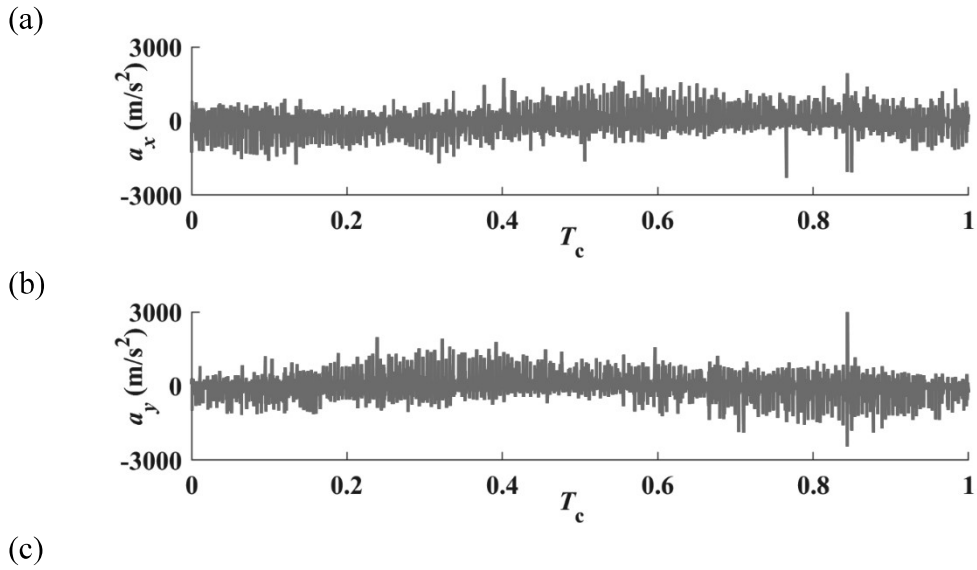


Fig. 15. Effect of the sun gear rotation speed on the (a) RMS and (b) MAX values of carrier accelerations.



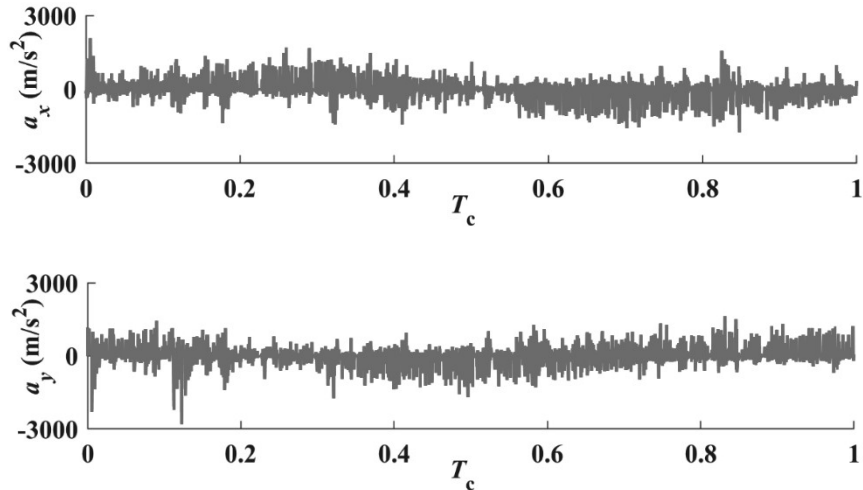


Fig. 16. Time-domain accelerations of the (a) inner planetary bearing cage in x direction; (b) inner planetary bearing cage in y direction; (c) outer planetary bearing cage in x direction; and (d) outer planetary bearing cage in y direction.

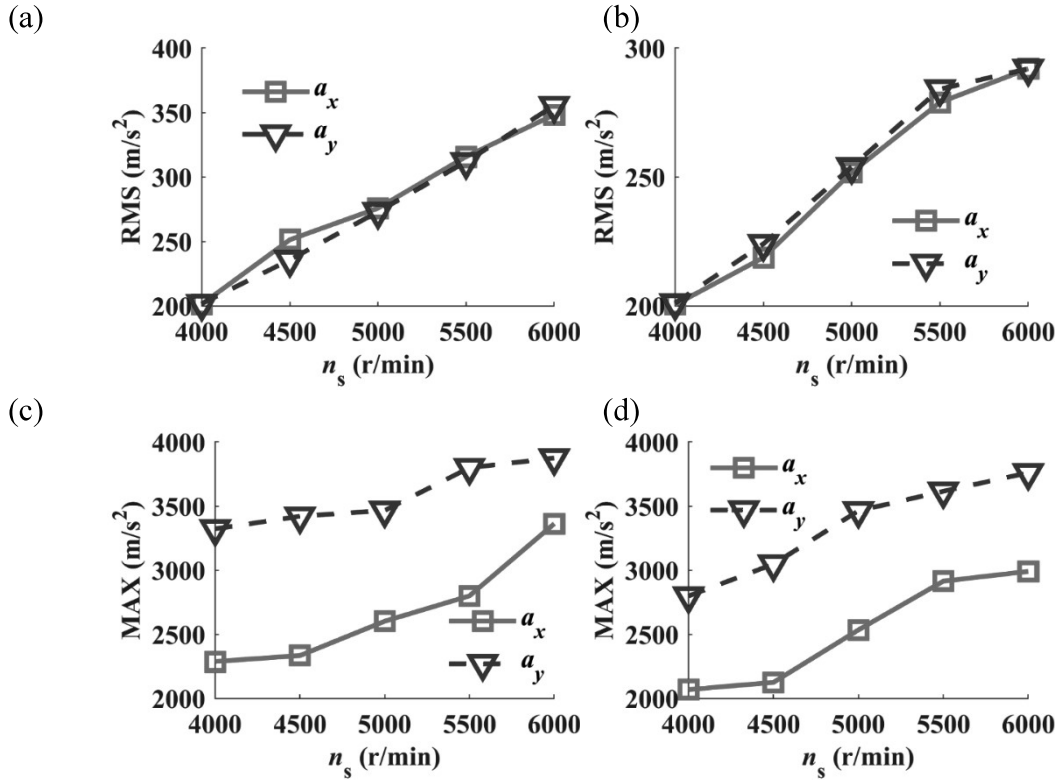


Fig. 17. Effect of the sun gear rotation speed on the (a) RMS value of inner planetary bearing cage; (b) RMS value of outer planetary bearing cage; (c) MAX value of inner planetary bearing cage; and (d) MAX value of outer planetary bearing cage.

Fig. 16 gives the time-domain planetary bearing cage acceleration when the sun gear rotation speed is 4000rpm. The inner/outer planetary bearing cage acceleration exhibit obvious periodic variations. The effect of the sun gear rotation speed on the planetary bearing cage acceleration is given in Fig. 17. In Fig. 17, the RMS and MAX values of inner/outer planetary bearing cage acceleration for different sun gear rotation speeds are given. The RMS and MAX values of inner/outer planetary bearing cage acceleration generally increase with the increment of sun gear rotation speed. Fig. 17(a) illustrates the effect of the sun gear rotation speed on the RMS value of inner planetary bearing cage acceleration. When the sun gear rotation speed is

from 4000rpm to 6000rpm, the RMS value of inner planetary bearing cage acceleration in  $x$  direction is from  $201.3\text{m/s}^2$  to  $348.2\text{m/s}^2$ ; and the RMS value of inner planetary bearing cage acceleration in  $y$  direction is from  $201.8\text{m/s}^2$  to  $355.4\text{m/s}^2$ . Fig. 17(b) illustrates the effect of the sun gear rotation speed on the RMS value of outer planetary bearing cage acceleration. When the sun gear rotation speed is from 4000rpm to 6000rpm, the RMS value of outer planetary bearing cage acceleration in  $x$  direction is from  $200.5\text{m/s}^2$  to  $291.9\text{m/s}^2$ ; and the RMS value of outer planetary bearing cage acceleration in  $y$  direction is from  $201.1\text{m/s}^2$  to  $292\text{m/s}^2$ . Fig. 17(c) illustrates the effect of the sun gear rotation speed on the MAX value of inner planetary bearing cage acceleration. When the sun gear rotation speed is from 4000rpm to 6000rpm, the MAX value of inner planetary bearing cage acceleration in  $x$  direction is from  $2287.7\text{m/s}^2$  to  $3361.9\text{m/s}^2$ ; and the MAX value of inner planetary bearing cage acceleration in  $y$  direction is from  $3322.4\text{m/s}^2$  to  $3874.2\text{m/s}^2$ . Fig. 17(d) illustrates the effect of the sun gear rotation speed on the MAX value of outer planetary bearing cage acceleration. When the sun gear rotation speed is from 4000rpm to 6000rpm, the MAX value of outer planetary bearing cage acceleration in  $x$  direction is increased  $2068.8\text{m/s}^2$  to  $2990.5\text{m/s}^2$ ; and the MAX value of outer planetary bearing cage acceleration in  $y$  direction is from  $2798.1\text{m/s}^2$  to  $3759.4\text{m/s}^2$ .

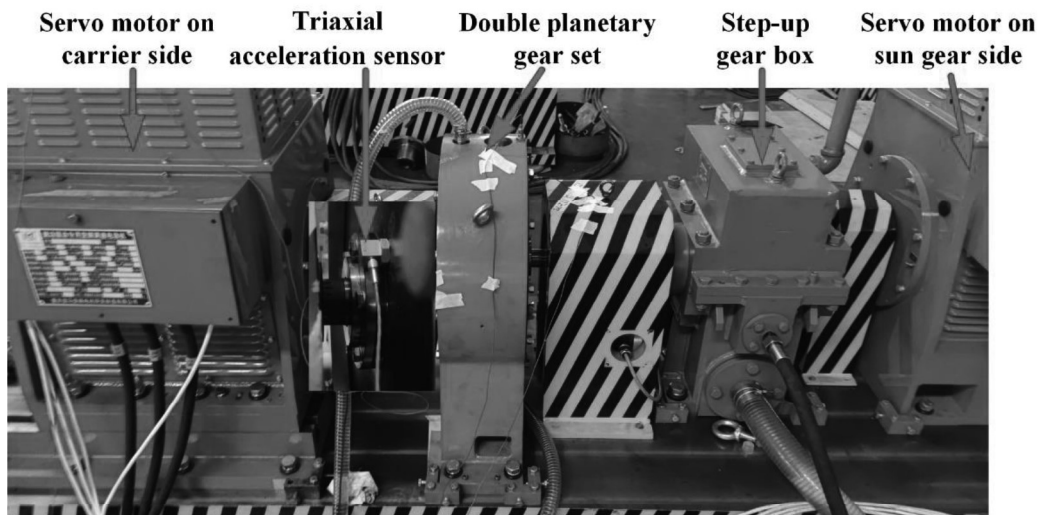
#### 4.3 Experimental verification and comparison

The DPGS vibration test rig is plotted in Fig. 18. Two servo motors are included in the test rig. One motor is used to adjust the carrier rotation speed. The other one is used to adjust the sun gear rotation speed. The servo motors are controlled by the control cubicle. The rotation speed and torque signal will be transmitted to the data acquisition system. Moreover, the DPGS vibrations are collected by a triaxial acceleration sensor installed on the carrier bearing housing and near the carrier axis. The acceleration signals from experiment and simulation are compared to prove the correctness of proposed model. The acceleration signals are obtained when the sun gear rotation speed is 4000rpm. The simulated acceleration signals are the acceleration of each part in the DPGS in the cartesian coordinate system. However, the triaxial acceleration sensor installed on the carrier bearing housing and near the carrier axial. The experimental acceleration signals are not the acceleration of any part in the DPGS. Thus, the simulated acceleration signals cannot be directly compared with the experimental ones. Considering the transmission path effect, the simulated acceleration signals can be given by [28]

$$a_p(t) = C_s a_s(t) + C_c a_c(t) + C_r a_r(t) + \sum_{i=1}^z (C_a a_a(t) + C_b a_b(t)) \quad (67)$$

where,  $C_s$ ,  $C_c$ ,  $C_r$ ,  $C_a$ , and  $C_b$  are the sun gear, carrier, ring gear, inner planet, and outer planet vibration contribution coefficient, respectively ( $C_s=0.57$ ,  $C_c=0.14$ ,  $C_r=0.269$ ,  $C_a=0.06(1-\cos(\theta_t))$ , and  $C_b=0.06$  in this work), detailed calculation method is given in Appendix A;  $a_s$ ,  $a_c$ ,  $a_r$ ,  $a_a$ , and  $a_b$  are the acceleration signals of sun gear, carrier, ring gear, inner planet, and outer planet, respectively.

(a)



(b)

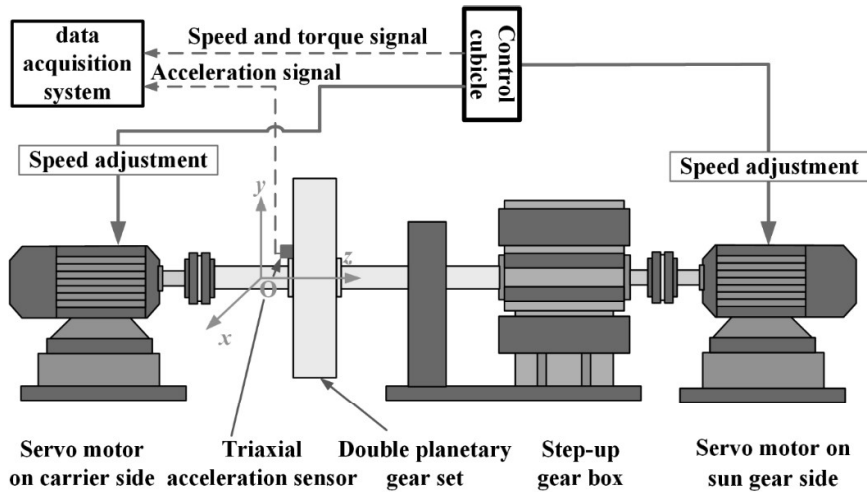
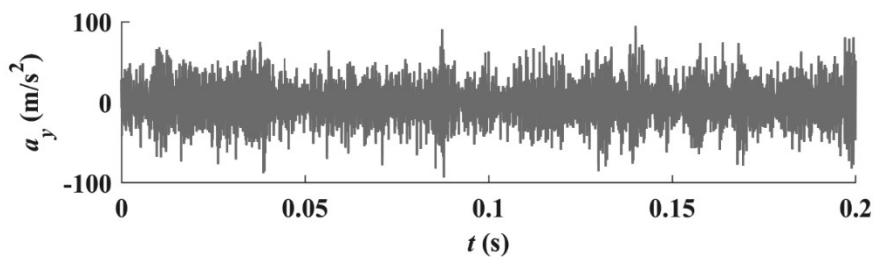


Fig. 18. A vibration test rig for DPGS. (a) Photograph and (b) Schematic diagram of test rig.

(a)



(b)

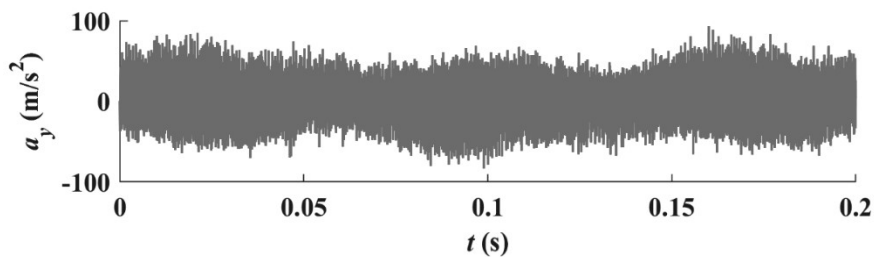


Fig. 19. Time-domain (a) experimental and (b) simulated acceleration signals of DPGS.

Fig. 19 gives the time-domain acceleration signals of DPGS. The experimental time-domain acceleration signals are shown in Fig. 19(a). In Fig. 19(a), the RMS value of experimental acceleration is  $26.31 \text{ m/s}^2$ ; and the MAX value of experimental acceleration is  $94.44 \text{ m/s}^2$ . The

simulated time-domain acceleration signals are shown in Fig. 19(b). In Fig. 19(b), the RMS value of simulated acceleration is  $22.49\text{m/s}^2$ ; and the MAX value of experimental acceleration is  $93.42\text{m/s}^2$ . The difference of RMS values between the experimental and simulated accelerations is 14.54%. The difference of MAX values between the experimental and simulated accelerations is 1.1%.

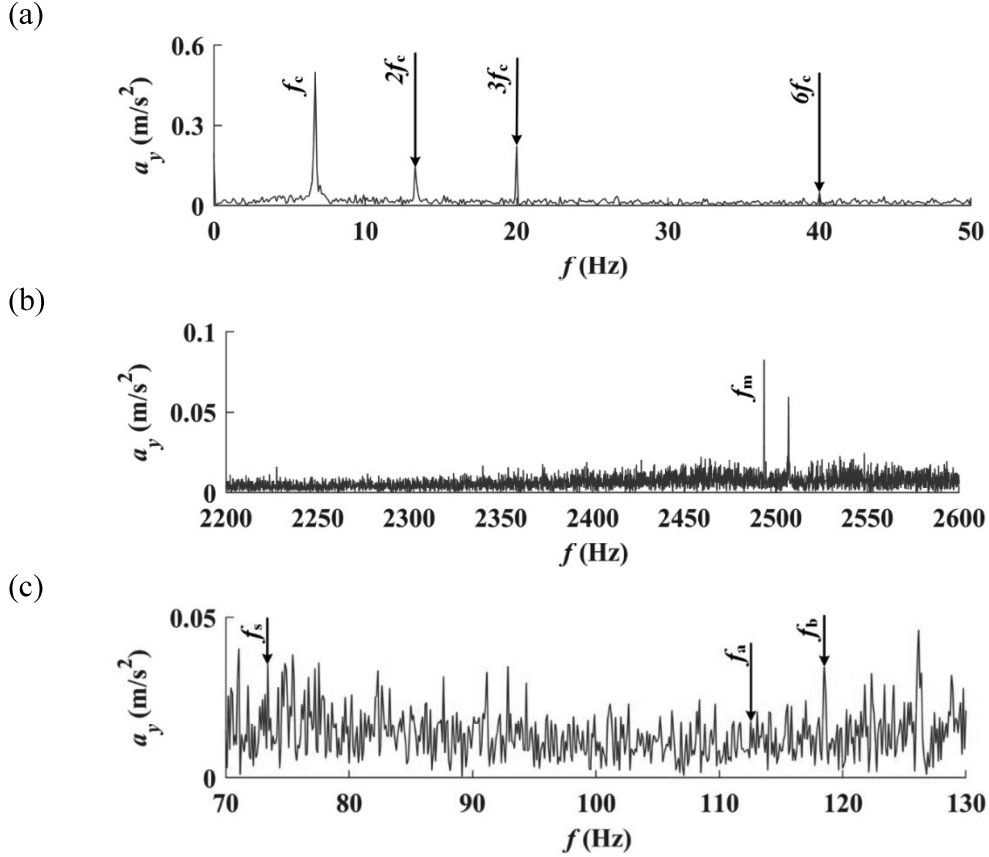
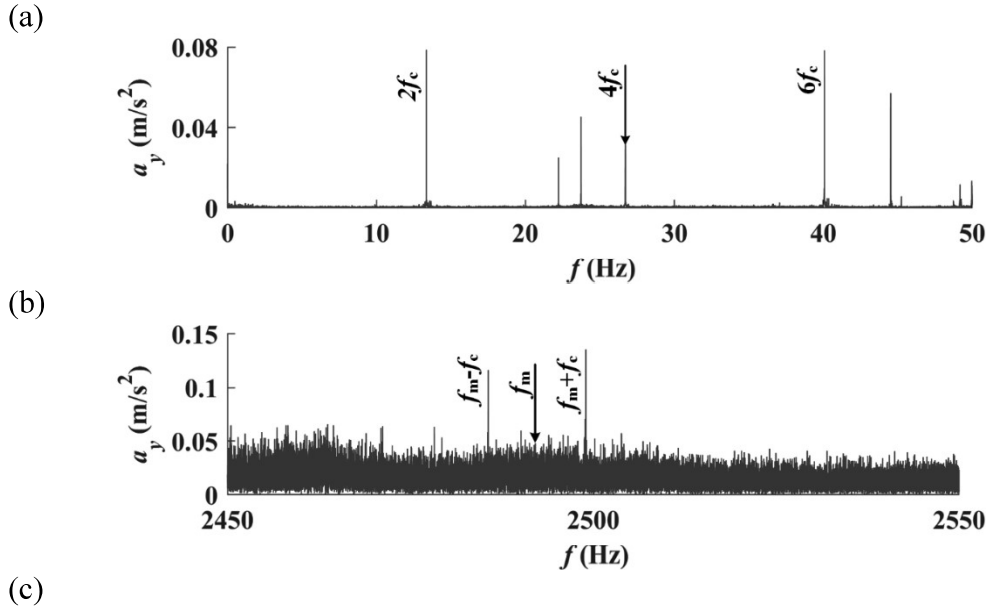


Fig. 20. Spectra of simulated DPGS acceleration signals. (a) 0-50Hz (b) 2200-2600Hz; and (c) 70-130Hz.



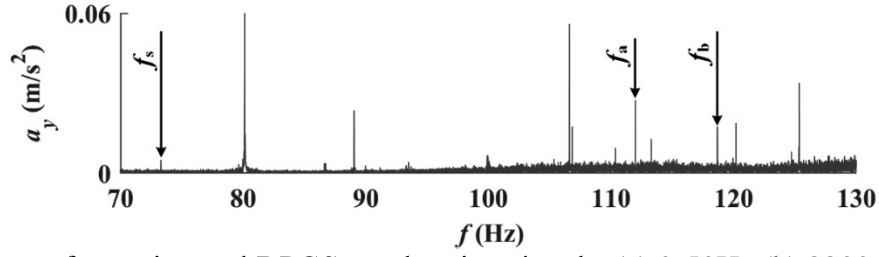


Fig. 21. Spectra of experimental DPGS acceleration signals. (a) 0-50Hz (b) 2200-2600Hz; and (c) 70-130Hz.

The spectra of simulated DPGS acceleration signals are given in Fig. 20. The spectra of experimental DPGS acceleration signals are given in Fig. 21. In Fig. 20 and Fig. 21,  $f_c$ ,  $f_m$ ,  $f_s$ ,  $f_a$ , and  $f_b$  are the carrier rotation frequency, gear meshing frequency, sun gear rotation frequency, inner planet rotation frequency, and outer planet rotation frequency, respectively. The carrier rotation frequency in the simulated acceleration spectrum is 6.70Hz, while its experimental value is 6.68Hz. The gear meshing frequency in the simulated acceleration spectrum is 2493.64Hz, while its experimental value is 2493.20Hz. Besides, because the test rig has planet axis error, the amplitudes are obviously larger in  $f_m \pm f_c$ . The sun gear rotation frequency in the simulated acceleration spectrum is 73.41Hz, while its experimental value is 73.35Hz. The inner planet rotation frequency in the simulated acceleration spectrum is 112.55Hz, while its experimental value is 113.30Hz. The outer planet rotation frequency in the simulated acceleration spectrum is 118.51Hz, while its experimental value is 118.70Hz. The simulated results are similar as the experimental results. Thus, the proposed dynamic model is correct.

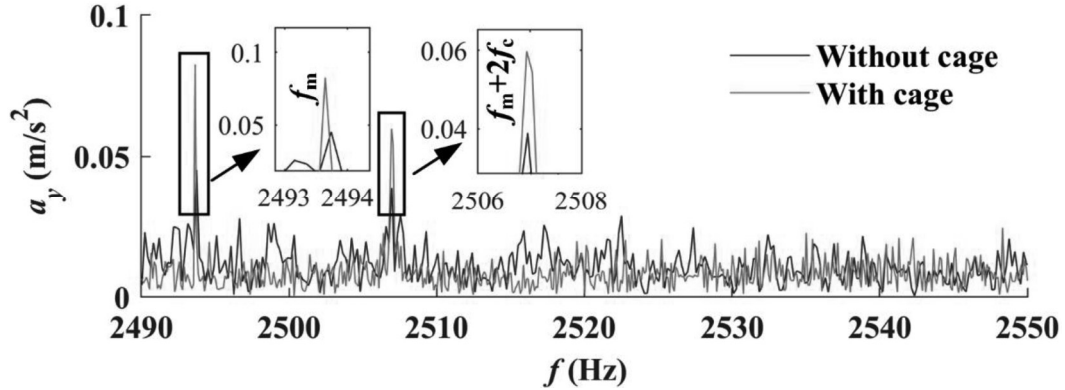
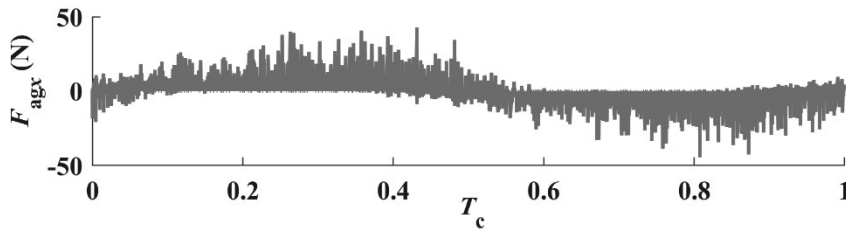
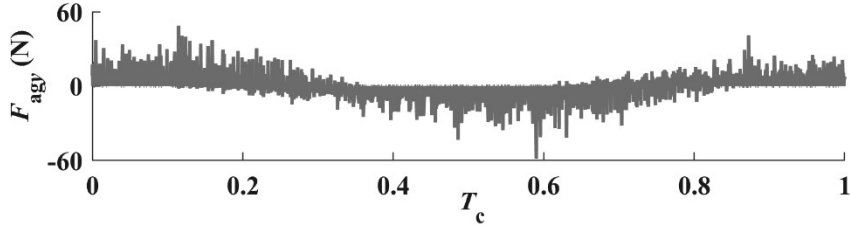


Fig. 22. Comparisons of the model with the cage (proposed model) and the one without the cage in the spectrum.

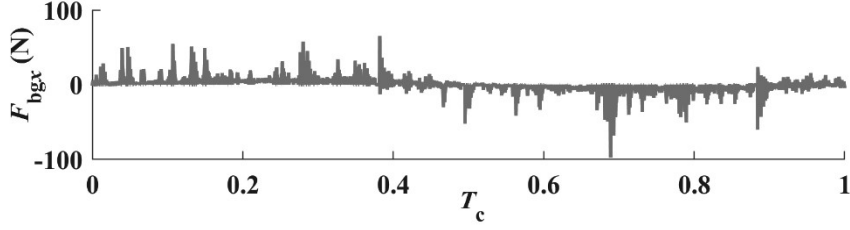
(a)



(b)



(c)



(d)

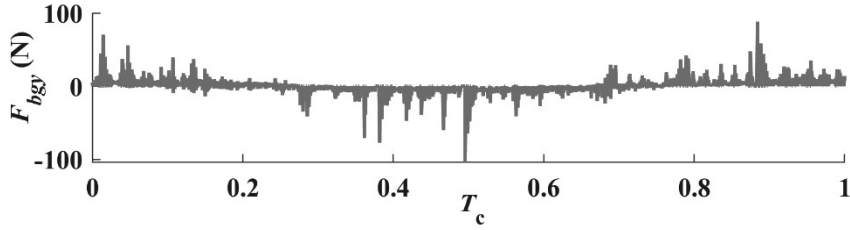


Fig. 23. The resultant forces of (a) inner planet-cage impact force and frictional force in  $x$  directions; (b) inner planet-cage impact force and frictional force in  $y$  directions; (c) outer planet-cage impact force and frictional force in  $x$  directions; and (d) outer planet-cage impact force and frictional force in  $y$  directions.

Fig. 22 gives the comparison of model with the cage (proposed model) and that without the cage in the spectrum. In Fig. 22,  $f_m$  and  $f_m+2f_c$  can be found for those with and without the cage models. The reason for the phenomenon is that the relative meshing phases of the planets are different. Thus, the meshing stiffness in the same angle position are different for different planet. Only when the planet rotates  $180^\circ$ , the DPGS will generate the same vibrations with the initial condition. On the other hand, the DPGS vibration spectrum will be modulated by  $2f_c$ . Thus, the frequency  $f_m+2f_c$  become more obvious in the spectrum. The values of  $f_m$  and  $f_m+2f_c$  are 2493.64 Hz; and that of proposed model is 2503.94 Hz; while those are 2493.74 Hz and 2506.96 Hz for the model without the cage. The experimental value of  $f_m$  is 2493.2 Hz, as shown in Fig. 21(b). The value of  $f_m$  obtained by the proposed model is closer to the experimental value than that obtained by the model without the cage. Moreover, the amplitudes of  $f_m$  and  $f_m+2f_c$  of proposed model are larger than those of model without the cage. The amplitudes of  $f_m$  and  $f_m+2f_c$  of proposed model are  $0.0827 \text{ m/s}^2$  and  $0.0598 \text{ m/s}^2$ ; while those of model without the cage are  $0.0454 \text{ m/s}^2$  and  $0.039 \text{ m/s}^2$ . The amplitude of  $f_m$  for proposed model is 38.3% larger than that of model without the cage; and the amplitude of  $f_m+2f_c$  for proposed model is 16.4% larger than that of model without the cage. This is because the proposed model considers the roller-cage impact force and cage vibrations, which will affect the gear rotation velocities and vibrations. Moreover, as shown in Fig. 23 and Fig. 6, the resultant forces of planet-cage impact force and



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frictional force are still not small values compared with the contact forces. Thus, the planet-cage impact force and frictional force will have a non-negligible effect on the DPGS vibrations. In addition, the planet-cage frictional force will have a significant effect on the planetary bearing sliding rate, which will also generate effect on the DPGS vibrations. These indicate that the cage has a significant effect on the DPGS vibrations. On the other hand, these indicate that the proposed model is more accurate.

## 5. Conclusions

This work proposes a dynamic model for vibration analysis of the planetary bearing in DPGS. The proposed model considers the planetary bearing roller and cage dynamic characteristics. Moreover, the gear interactions, planetary gear-bearing-carrier interactions, and roller-cage interactions are contained in the proposed model. In comparison with the previous studies, this work establishes an improved dynamic model for a DPGS containing all components (sun gear, ring gear, carrier, inner planet, outer planet, planetary bearing roller, and planetary bearing cage). The components' excitations in this work are more comprehensive. An experiment is used to prove the correctness of the proposed model. The conclusions are summarized as follows:

(1) The planetary bearing contact force, roller-cage impact force, planetary bearing roller acceleration, planet acceleration, carrier acceleration, and planetary bearing cage acceleration exhibit obvious periodic variations. The contact force periodic frequency of inner planetary bearing approximately equal to  $f_{ga}-f_c$  and that of outer planetary bearing approximately equal to  $f_{ga}+f_c$ .

(2) The RMS and MAX values of planetary bearing contact force and roller-cage impact force generally increase with the increment of sun gear rotation speed.

(3) The RMS and MAX values of planetary bearing roller acceleration, planet roller acceleration, carrier roller acceleration, and planetary bearing cage roller acceleration generally increase with the increment of sun gear rotation speed.

(4) The simulated results are similar as the experimental results both in the time-domain and frequency-domain. The value of  $f_m$  obtained by the proposed model is closer to the experimental value than that obtained by the model without the cage. Moreover, the amplitudes of  $f_m$  and  $f_m+2f_c$  of proposed model are larger than those of model without the cage. Thus, the proposed dynamic model is correct.

Although the proposed model is more comprehensive than other models, it also has some aspects need to be improved. The planetary bearing cage has a big centrifugal force rotating about the planetary gear set axis. It will generate significant flexible deformation due to the cage-planet impact, which will affect the cage-planet contact stiffness and cause the change of DPGS vibrations. Moreover, the effect of the gear backlash on the DPGS vibrations are ignored in this work. However, the gear backlash is an important excitation source for DPGS and has a significant effect on the DPGS vibrations. Therefore, the cage flexibility and gear backlash need to be considered in the future work.

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## Conflict of Interest

The author declared that he has no conflicts of interest.

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#### Appendix A: Calculation method for the vibration contribution coefficients

$$y_p = \sum_{k=1}^{N_t} \left( \text{RMS}(a_{pk}) + \text{MAX}(a_{pk}) \right) \quad (1)$$

where,  $a_{pk}$  is the simulated acceleration signals of  $k$ -th work condition,  $N_t$  is the number of working conditions of calculation group.

$$y_t = \sum_{k=1}^{N_t} \left( \text{RMS}(a_{tk}) + \text{MAX}(a_{tk}) \right) \quad (2)$$

where,  $a_{pk}$  is the experimental acceleration signals of  $k$ -th work condition.

$$f = |y_p - y_t| \quad (3)$$

The vibration contribution coefficients can be obtained by finding the minimum value of  $f$ , the problem can be expressed by

$$\min_x f(x) \text{ such that } \begin{cases} x(1) - x(2) = 0 \\ 0 \leq x \leq 1 \end{cases} \quad (4)$$

$$x = [C_a \quad C_b \quad C_c \quad C_s \quad C_r] \quad (5)$$

Then using the validation group signals to verify the correctness of the vibration contribution coefficients

$$\hat{y}_p = \sum_{k=1}^{N_c} \left( \text{RMS}(a_{pk}) + \text{MAX}(a_{pk}) \right) \quad (6)$$

$$\hat{y}_t = \sum_{k=1}^{N_c} \left( \text{RMS}(a_{tk}) + \text{MAX}(a_{tk}) \right) \quad (7)$$

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where,  $N_c$  is the number of working conditions of the validation group.

$$\text{rel} = \frac{|\hat{y}_p - \hat{y}_t|}{\hat{y}_t} \quad (8)$$

If  $\text{rel} < 0.15$ , the obtained vibration contribution coefficients can be recognized as correct.