

Saturn's Weather-Driven Aurorae Modulate Oscillations in the Magnetic Field and Radio Emissions

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Key Points:

- Keck-NIRSPEC observations of Saturn's northern H_3^+ infrared auroral emission from 2017 are analysed
- First clear picture of how the ionosphere moves in relation to planetary period currents is provided
- Saturn's measured variable rotation rate is driven by twin-vortex flows in the upper atmosphere

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Abstract

The Cassini spacecraft revealed that Saturn’s magnetic field displayed oscillations at a period originally thought to match the planetary rotation rate but later found not to. One of many proposed theories predicts that a polar twin-cell neutral weather system drives this variation, producing observable differences in flows within Saturn’s ionosphere. Here, using spectral observations of auroral H_3^+ emission lines taken by the Keck Observatory’s Near Infrared Echelle Spectrograph (Keck-NIRSPEC) in 2017, we derive ion line-of-sight velocity maps after grouping spectra into rotational quadrants matching phases of the planetary magnetic field. We measure 0.5 km s^{-1} wind systems in the ionosphere consistent with predicted neutral twin-vortex flow patterns. These findings demonstrate that neutral winds in Saturn’s polar regions cause the rotational period, as determined via the magnetic field, to exhibit differences and time variabilities relative to the planet’s true period of rotation in a process never before seen within planetary atmospheres.

Plain Language Summary

We observed Saturn’s northern aurorae in the infrared using the Keck Observatory in Mauna Kea, Hawaii over the course of June, July and August of 2017. Using this data we investigate the motion of an ion, H_3^+ , in the planet’s upper atmosphere. This is done after first placing the data into four groups corresponding to the rotational phase of the planet’s magnetic field. By doing so we are able to detect twin-vortex flows in the upper atmosphere of Saturn, consistent with theories that predict the presence of such a polar feature, thus providing direct evidence that Saturn’s measured variable rotation rate is driven by these flows. These twin-vortex flows are ultimately responsible for the time differences relative to the planet’s true rotation period observed throughout Saturn’s planetary environment.

1 Introduction

An abiding mystery following Cassini’s extended tour of Saturn is also one of the first questions the spacecraft raised about the planet: measurements showed that Saturn’s day appeared to be 6 minutes longer at $\sim 10\text{h } 45\text{m}$ (Gurnett et al., 2005) than that measured by the Voyager 1 and 2 spacecraft (Kaiser et al., 1980). Since it is improbable that the interior of Saturn changed its rotation period over the course of only two decades, it was clear that somehow the magnetic fields above the planet were slipping relative to those generated deep within the interior (Stevenson, 2006). This mystery has remained unresolved despite nearly two decades of Cassini observations at Saturn.

First detected in the radio Saturn kilometric radiation and subsequently throughout the Cassini mission, evidence for a varying planetary rotation rate has been found in numerous parameters including: magnetospheric variations in the plasma (Gurnett et al., 2007), energetic neutrals (Paranicas et al., 2005) and the axisymmetric magnetic field (Giampieri et al., 2006), as well as ultraviolet (Nichols et al., 2008) and infrared (Badman et al., 2012) auroral emissions. The rotation rate also measurably drifts with time (Kurth et al., 2008), with two independent rates in the northern and southern hemispheres (Gurnett et al., 2009), linked to the changing season (Gurnett et al., 2010). These variations are propagated throughout the Saturnian environment by two large-scale planetary period current systems flowing between the ionosphere and magnetosphere (G. J. Hunt et al., 2014).

A wide range of models have been developed in order to explain the source of these planetary period currents with limited observational constraints to substantiate many of them. Some models place the source within the magnetosphere, caused by the Enceladus torus (Gurnett et al., 2007; Goldreich & Farmer, 2007; Burch et al., 2008), inter-

actions with Titan (Winglee et al., 2013), or periodic latitudinal oscillations in the plasma sheet (Carbary et al., 2007). It is worth noting that magnetosphere-driven models are generally unable to robustly explain the distinction between the two independent periodicities observed in each polar region. Other models suggest that the source originates within Saturn’s atmosphere, the result of changing ionospheric conductivity (Gurnett et al., 2007), flows in Saturn’s stratosphere driving Hall drift (Smith, 2014), or the westward drift of ionospheric Rossby waves (Smith, 2018). Conversely, an alternative theory of how these currents are driven is that ions are forced to move through polar magnetic fields by collisions with rotating twin-vortex flows within Saturn’s neutral polar thermosphere (Smith, 2006, 2011; Jia et al., 2012; Southwood & Cowley, 2014) – effectively a form of weather-driven aurora that is itself the result of neutral and ion winds flowing in the planet’s upper atmospheric layers.

Tests of these theories using magnetospheric observations are difficult since both the magnetospheric currents and generated aurora are nearly identical whether they originate in the magnetosphere or the atmosphere. The only observational evidence for an atmospheric source comes from an apparent 1° modulation in the location of the planetary period auroral current layer, with maximum and minimum modulations occurring at southern magnetic phases of 270° and 90° , respectively, where local magnetic phases are analogous to magnetic longitude (G. J. Hunt et al., 2014). However, since the model predicting that these currents are driven by a thermospheric twin-cell vortex explicitly requires that the ionosphere flows in the opposite direction from other models, the motion of ions within the ionosphere provides a unique diagnostic of the source of planetary period currents (Smith, 2014).

Figure 1 shows these expected ion flows as would be seen by an observer at Earth with planetary dawn to the left and dusk to the right. In the case of a magnetospheric driver, at a noon northern magnetic phase of 0° (Ψ_{0°) as shown in Figure 1(a), ion-neutral collisions in the ionosphere lead to an anti-sunward flow (red-shifted) away from Earth-based observers over the central polar region with a return flow (blue-shifted) towards observers at lower latitudes. For an atmospheric driver at Ψ_{0° , seen in Figure 1(b), the neutral wind drives parallel plasma flow in the ionosphere, resulting in a sunward flow (blue-shifted) towards the observer over the polar region and an anti-sunward return flow (red-shifted) away from the observer at lower latitudes (G. J. Hunt et al., 2014; Jia et al., 2012; Southwood & Cowley, 2014; G. Hunt et al., 2015). In addition, these patterns rotate with phase at the measured period of the magnetic field oscillations.

2 Data Analysis

To test the thermospheric twin-cell vortex hypothesis, we used the Keck Observatory’s Near Infrared Echelle Spectrograph (Keck-NIRSPEC) (McLean et al., 1998) to scan the auroral region in a manner similar to a study at Jupiter (Johnson et al., 2017), measuring emission from H_3^+ – a dominant molecular ion species in Saturn’s ionosphere. The peak emission altitude of H_3^+ is at around 1,155 km above the atmospheric 1 bar level (T. S. Stallard et al., 2012) and is more generally found to exist between 1,000 to 2,000 km of the same altitudinal level (T. Stallard et al., 2010) which is, as demonstrated by Moore et al. (2010) and Shebanits et al. (2020), in the electrical dynamo region of the ionosphere. In a previous publication by T. S. Stallard et al. (2019), techniques were established to measure the line-of-sight velocity of H_3^+ emission lines on individual nights, providing a view of the varying ionospheric velocity in the Earth’s reference frame. Similar work by O’Donoghue et al. (2016) and Chowdhury et al. (2019) have also provided insights into the H_3^+ aurorae at Saturn in recent years.

Here, we measure the ion velocities after mapping data into four rotational bins matching the phase (Ψ_N) of the magnetic field oscillation resulting from the rotating current system (Provan et al., 2018): Ψ_{0° ($315^\circ - 45^\circ$), Ψ_{90° ($45^\circ - 135^\circ$), Ψ_{180° ($135^\circ -$

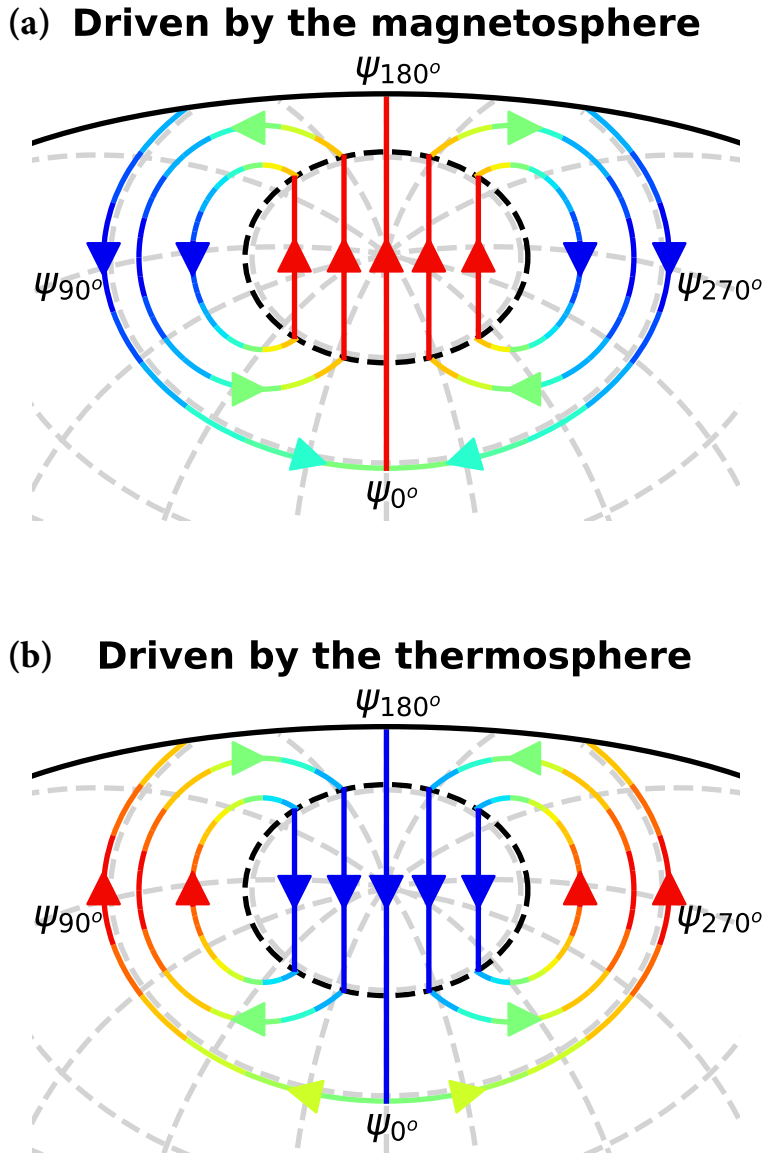


Figure 1. Predicted ion wind flows in Saturn’s northern thermosphere (as observed from Earth) based on a magnetospheric (a) and thermospheric (b) source with planetary dawn on the left and dusk to the right. We adapt the predicted ion flows for a magnetospheric (G. J. Hunt et al., 2014) and atmospheric (G. Hunt et al., 2015) driver into what would be observed in the line-of-sight from Earth when the central meridian planetary phase (Ψ_N) is 0° . Colors indicate in broad-scale the expected line-of-sight blue- and red-Doppler shifts. Magnetic fields are not shown here for the sake of the clarity, but can be found in the original works.

121 225°) and Ψ_{270° (225° – 315°), in order to reveal the underlying ion winds that are as-
 122 sociated with the rotating current system. Tables S1 and S2 in the Supporting Informa-
 123 tion outline the corresponding night-by-night and quadrant-by-quadrant breakdown of
 124 spectra into the four aforementioned groupings of northern planetary magnetic phase.

125 Saturn’s auroral currents and thus our line-of-sight velocity maps consist of sub-
 126 corotational flows associated with the outer magnetosphere. Both the co-rotational flows
 127 and the outer magnetosphere are thought to be broadly fixed in local-time with the for-
 128 mer rotating at the planetary period. In this work, we assume the local-time currents
 129 are the same no matter the rotational phase, while the rotating ion flows (containing both
 130 local-time and rotational phase components) are approximately equal in magnitude but
 131 opposite in direction for opposing phase quadrants. As such, subtracting the ion winds
 132 measured in two rotational phases in anti-phase (e.g. Ψ_{0° and Ψ_{180°) removes all the local-
 133 time fixed flows, thus isolating the rotating component of the flow which has subsequently
 134 doubled as a result of the subtraction; the same technique has previously been used with
 135 magnetospheric currents (G. J. Hunt et al., 2014). After dividing these data products
 136 by two in order to give an average velocity of rotational phase-associated ion flows, we
 137 are able to produce two difference maps of the ion winds ($\Psi_{0^\circ-180^\circ}$ and $\Psi_{90^\circ-270^\circ}$) at
 138 planetary phases that are perpendicular to one another.

139 Please refer to the Supporting Information for a more explicit account of the arith-
 140 metical considerations involved in this process.

141 3 Results

142 Figure 2(c) shows profile plots of the average intensity ($\Psi_{0^\circ+180^\circ}$) and the ion dif-
 143 ference winds within the rotating phase ($\Psi_{0^\circ-180^\circ}$), computed as the means of the re-
 144 gion delineated by dashed horizontal lines in Figure 2(a) and (b). This shows: blue shifted
 145 +0.5 km s⁻¹ ion winds across the pole, inside (poleward of) the main auroral oval, well
 146 above the errors; a stagnant region close to the main auroral emission; and, a red-shifted
 147 -0.5 km s⁻¹ ion wind along the flanks (at lower latitudes), equatorward of the main au-
 148 roral emission. We would expect auroral intensity to be symmetric about the pole, but
 149 instead observe a brightening on the dusk side (right side) of the aurora; this indicates
 150 that an upward current in this region that is stronger at dusk may be affected by local-
 151 time ion winds. Figure 2(f) shows the perpendicular average intensity ($\Psi_{90^\circ+270^\circ}$) and
 152 ion difference velocity ($\Psi_{90^\circ-270^\circ}$). These are more symmetric in intensity but include
 153 a brightening near noon similar to a local-time cusp brightening observed in previous in-
 154 tensities (Badman et al., 2011); the velocities are harder to interpret until placed in the
 155 context of Figure 3.

156 In Figure 3 we combine the two perpendicular line-of-sight ion wind difference maps
 157 from Figure 2(b) and (e) – $\Psi_{0^\circ-180^\circ}$ flowing from midnight (top) to noon (bottom) and
 158 $\Psi_{90^\circ-270^\circ}$ flowing from dusk (right) to dawn (left) – into a two-dimensional map of iono-
 159 spheric flows as seen from Earth at a noon rotational phase of Ψ_{0° , enabling compar-
 160 ison with the schematics presented in Figure 1. This clearly shows a strong blue-shifted
 161 flow over the pole, and a red-shifted return flow along the dawn side of the lower lati-
 162 tude sub-auroral region, as well as some evidence for a second return flow on the dusk
 163 side of the polar region. This accords well with models that make use of a thermospheric
 164 driver to explain the observed behavior in Saturn’s polar auroral regions (Smith, 2011).

165 Ions across the pole can only flow towards us in this configuration if they are blown
 166 by a strong coincident thermospheric neutral wind, as other sources would cause the flow
 167 to move away from us (Smith, 2014). Magnetospheric sources of the planetary period
 168 oscillations cannot explain the observed flows, as the magnetospheric currents would drive
 169 ions to flow away from the observer over the pole, irrespective of the sources of this os-

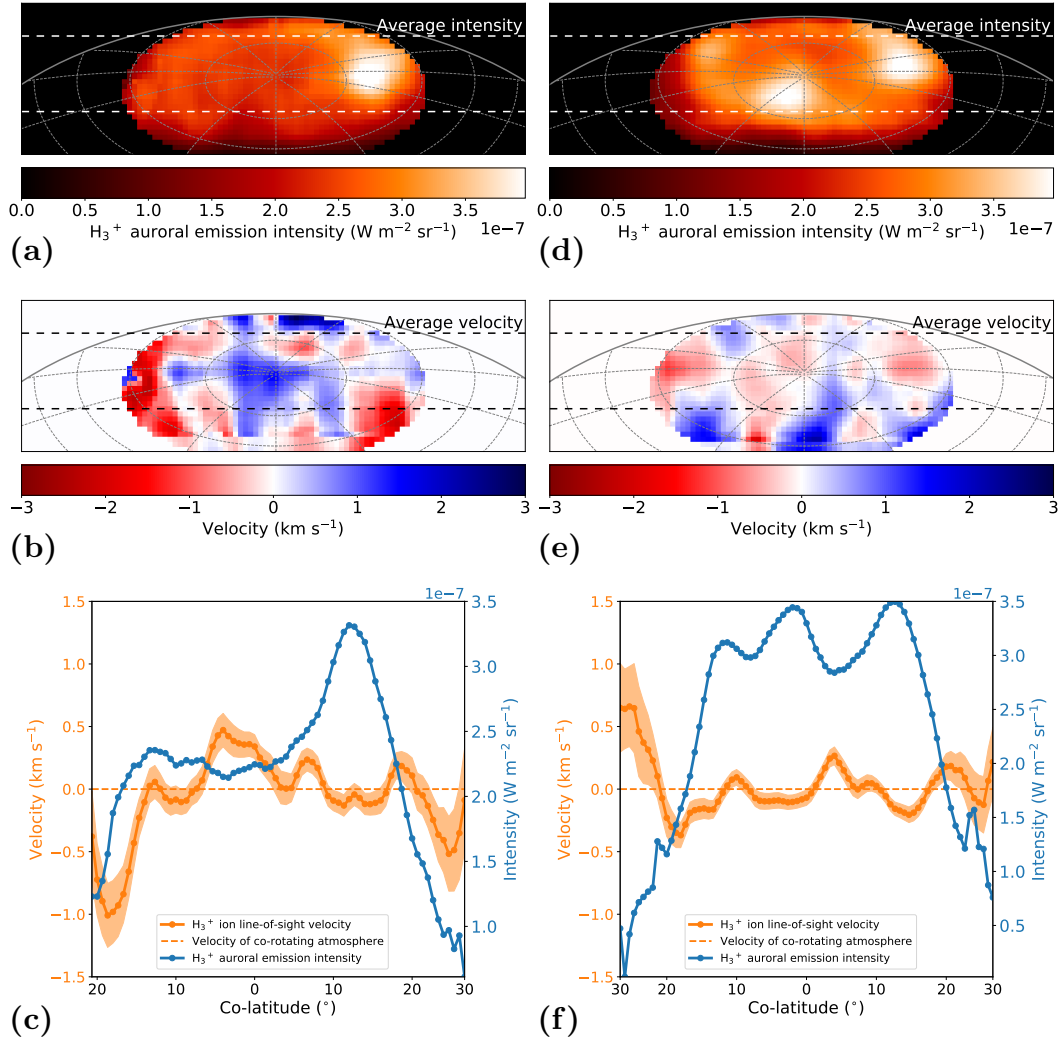


Figure 2. Auroral emission intensity maps and ion wind velocity difference maps for $\Psi_{0^\circ-180^\circ}$ and $\Psi_{90^\circ-270^\circ}$ with planetary dawn to the left and dusk to the right. We make use of the individual emission intensity and ion line-of-sight velocity maps for each of the four quadrant groupings of rotational phase (shown in Figure S1 in the Supplementary Information) to produce average emission intensity maps for Ψ_{0° and Ψ_{180° ($\Psi_{0^\circ+180^\circ}$), shown in panel (a), and for Ψ_{90° and Ψ_{270° ($\Psi_{90^\circ+270^\circ}$), shown in panel (d). We also subtract the observed ion velocity at Ψ_{180° from Ψ_{0° and the velocity at Ψ_{270° from Ψ_{90° to create ion wind difference maps for $\Psi_{0^\circ-180^\circ}$ and $\Psi_{90^\circ-270^\circ}$ displayed in panels (b) and (e), respectively. Overall H_3^+ emission structures shown in (a) and (d) reveal a bright spot on the dusk of $\Psi_{0^\circ+180^\circ}$ and broadly symmetric emission in $\Psi_{90^\circ+270^\circ}$. The line-of-sight velocity difference maps, seen in (b) and (e), show clear structures with $\Psi_{0^\circ-180^\circ}$ dominated by a blue-shift over the pole and red-shift on the flanks, and $\Psi_{90^\circ-270^\circ}$ highlighting more complexity. The average emission intensities and ion winds for each (taken from between the dotted horizontal lines in the maps) are illustrated in (c) and (f). These show intensities of $\sim 3.5 W m^{-2} sr^{-1}$ across the auroral region and ion wind difference flows up to $\sim \pm 0.5 km s^{-1}$, well above the errors.

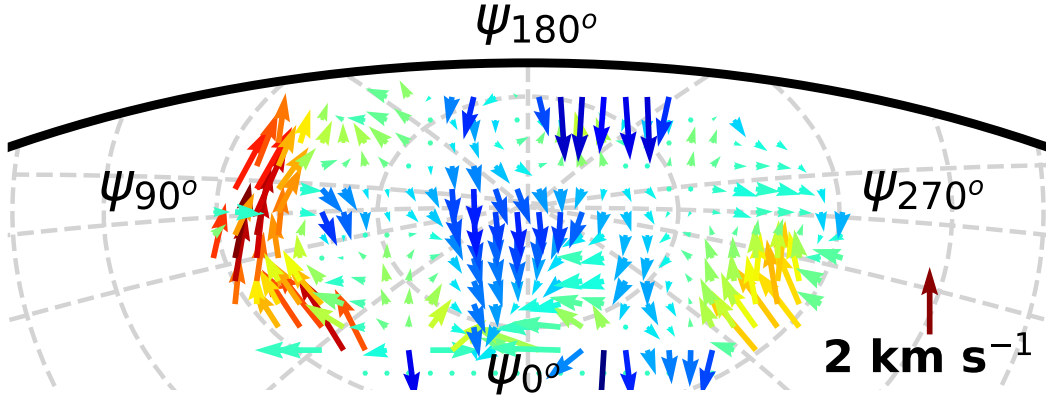


Figure 3. Our combined observed ion flows as seen from Earth with Saturnian dawn to the left and dusk to the right. Velocities are again shown at a noon magnetic phase of Ψ_{0° , as in Figure 1. The vectors represent flows from Figure 2(b) and (e), with blue-shifted $\Psi_{0^\circ-180^\circ}$ flowing from midnight to midday (top to bottom), and $\Psi_{90^\circ-270^\circ}$ flowing from dusk to dawn (right to left). The magnitude of the arrows is scaled with velocity up to $\sim 2.4 \text{ km s}^{-1}$, and the blue-to-red color represents the magnitude in the $\Psi_{0^\circ-180^\circ}$ direction. Please note that the magnitudes of the arrows shown here correspond exactly with velocities presented in Figure 2(b) and (e). A clear blue-shift across the pole and red-shift along the lower latitude flanks matches well with Figure 1(b), making it consistent with a thermospheric origin for the ion winds. It is as yet unclear why the dawn (left) side vortex is more intense than the much weaker corresponding dusk (right) side vortex. One possibility may be that the current emerging from the dawn side vortex (Ψ_{90°) is higher than the one which emerges from the dusk side vortex (Ψ_{270°). Furthermore, it is not unreasonable to surmise that the dawn side vortex will weaken as it rotates into noon (Ψ_{0°) and through to the dusk side while the dusk side vortex will strengthen as it rotates into midnight (Ψ_{180°) through to the dawn side.

170 cillation (Gurnett et al., 2007; Goldreich & Farmer, 2007; Burch et al., 2008; Winglee
171 et al., 2013; Carbary et al., 2007).

172 4 Discussion & Conclusions

173 Having excluded all other theories of how the planetary period current system is
174 generated, our measurements of a sunward (midnight to noon) flow across the pole in
175 the northern ionosphere provides direct evidence that the variable measured rotation rate
176 of Saturn is produced by a source within the planet’s atmosphere. The observed flows
177 match very well with predicted flows from models using a previously undetected twin-
178 vortex within the thermosphere of the planet, driving currents out into the magnetosphere
179 and producing a radio pulse. The temporal variations in rotational period seen within
180 the radio emission, and throughout Saturn’s magnetosphere, must result from long-term
181 localised variations in the rotation speed of this atmospheric twin-vortex.

182 The atmospheric vortex model proposed by (Jia et al., 2012) for a single source in
183 one hemisphere and for dual sources in two hemispheres (Jia & Kivelson, 2012), which
184 successfully reproduced many of the observed magnetospheric periodicities, invoked vor-
185 tical flows in the ionosphere with speeds ranging between 0.3 and 3 km s⁻¹. Unfortu-
186 nately, this does not enable us to conclude anything substantial relating to the altitu-
187 dinal location of the driving vortices because not enough is known about the dynamics
188 of the thermosphere or indeed the underlying neutral atmospheric weather that drives
189 this ion flow and the resultant magnetospheric current system. Models of how thermo-
190 spheric flows might generate these ion winds resulted in a comparable twin-cell flow by
191 placing a heat source located at auroral latitudes between Ψ_{30° and Ψ_{210° (Smith, 2011).
192 This places the source of heating close to the region of brightest aurora, perhaps hint-
193 ing that these flows may be self-sustaining. However, such modelling predicts much smaller
194 ion winds than those we observe.

195 Subsequent modelling predicted significant heating when driving flows at the mag-
196 nitude observed here. Future observations should constrain the neutral flows and how
197 they drive this ionospheric twin-cell vortex, and help resolve whether localised heating
198 places the thermal source of this twin-cell in the thermosphere itself, or if it is driven by
199 other deeper atmospheric processes (Smith, 2014). Additionally, with favorable viewing
200 geometries, observations of the southern hemisphere will also allow the other half of this
201 picture to be pieced together.

202 Our findings provide direct evidence that it is the thermosphere driving the sys-
203 tem and helps to answer the long-standing question of why the period of magnetic field
204 oscillations resulting from the current system varies at Saturn. We have also presented
205 the first known example of weather-driven aurora as neutral winds in the thermosphere
206 drive the currents which go on to exert forces in the magnetosphere. If atmospheric weather
207 at Saturn can drive intensity perturbations in auroral emissions then it stands to high-
208 light the potential importance of giant planets possessing significantly strong feedback
209 systems whereby neutral atmospheres can drive the magnetospheres; a major factor in
210 the context of the Saturn field-aligned current.

211 Recent observations at Earth have evidenced non-auroral forcing of the equatorial
212 magnetospheric regions, as neutral atmospheric dynamics appear to drive ionospheric
213 flows at altitudes of 600 km above the surface (Immel et al., 2021). A similar example
214 of equatorial neutral driving is suggested for the azimuthal magnetic field structures ob-
215 served in Saturn’s equatorial regions, where tropospheric winds have been evoked to ex-
216 plain magnetic field variations (Khurana et al., 2018; Agiwal et al., 2021). Both Uranus
217 and Neptune have highly inclined magnetic poles, placing them in these same equato-
218 rial regions, suggesting even modest flows similar to either Earth or Saturn could drive
219 dramatic new weather-driven aurora at those planets.

220 Observations have shown that planets close to their star experience intense atmo-
 221 spheric heating, resulting in strong neutral winds. For example, HD 189733b is a Hot
 222 Jupiter that has a lower atmosphere super-rotating at 3 km s^{-1} that becomes a verti-
 223 cal 40 km s^{-1} neutral outflow at the top of the upper atmosphere (Seidel et al., 2020).
 224 These authors evoke the super-rotation of ions in the upper atmosphere across the mag-
 225 netic field to explain this outward atmospheric flow. Such an intense neutral wind in-
 226 teraction is vastly stronger than the ion winds observed at Saturn, and would result in
 227 commensurately stronger auroral currents. These flows may help explain the detection
 228 of aurora (Pineda, 2020) from a proposed close-in Hot Super-Earth planet GJ 1151b (Mahadevan
 229 et al., 2021), the production of which does not match with present theories of auroral
 230 generation at these worlds, perhaps indicating that close-in planetary aurora are dom-
 231 inated by atmospherically driven interactions, rather than currents from the surround-
 232 ing space environment.

233 While planetary radio measurements remain a powerful tool for determining plan-
 234 etary rotation rates, interpreting such measurements requires detailed knowledge of ion-
 235 eutral coupling in the ionosphere-thermosphere. This work provides observational evi-
 236 dence to support such arguments. As such, ion wind measurements could be used to ex-
 237 plore the impact of atmospheric effects on planetary periods, particularly in cases where
 238 there exists an axisymmetric planetary magnetic field and a system with a tilted dipole
 239 masking the effects of such setups, both within our own solar system and farther afield.

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252 Open Research: Keck-NIRSPEC data used in this study are publicly available on
 253 the Keck Observatory Archive (KOA) [<https://www2.keck.hawaii.edu/koa/public/koa.php>].
 254 Similarly, the planetary period oscillation phases obtained from NASA Cassini magnetic
 255 field data that have been employed in this study are available from the University of Le-
 256 icester Research Archive [<http://hdl.handle.net/2381/42436>]. Further details are out-
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