| 1 | North-south Asymmetric Nightside Distorted Transpolar Arcs within A Framework |
|----------------------|--|
| 2 | of Deformed Magnetosphere-Ionosphere Coupling: IMF-By Dependence, |
| 3 4 | Ionospheric Currents, and Magnetotail Reconnection |
| 5 6 7 | Motoharu Nowada ¹ , Qiu-Gang Zong ² , Benoît Hubert ³ , Quan-Qi Shi ¹ , Yong-Fu Wang ² , Jun Yang ¹ , Adrian Grocott ⁴ , Alexander W. Degeling ¹ , An-Min Tian ¹ , Xu-Zhi Zhou ² , and Chao Yue ² |
| 8 9 10 | ¹ Shandong Key Laboratory of Optical Astronomy and Solar-Terrestrial Environment, School of Space Science and Physics, Institute of Space Sciences, Shandong University, Weihai, Shandong, People's Republic of China. |
| 11 12 | ² Institute of Space Physics and Applied Technology, School of Earth and Space Sciences, Peking University, People's Republic of China. |
| 13 14 | ³ Space science, Technologies and Astrophysics Research (STAR) Institute, Université de Liège, Belgium. |
| 15 16 | ⁴ Space and Planetary Physics Group, Department of Physics, Lancaster University, Lancaster, UK. |
| 17 | |
| 18 | Key points: |
| 19 20 | 1. A new morphological type of transpolar arc, characterized by large nightside distortions in the pre- or post-midnight sector, is described. |
| 21 22 | 2. Nightside reconnection and magnetotail deformation by IMF penetration play essential roles in the formation of nightside distorted TPA. |
| 23 24 | 3. Nightside distorted TPAs can be used as a remote-sensing tool to diagnose globally IMF- deformed magnetospheric processes. |
| 25 26 27 28 | Corresponding Authors: Motoharu Nowada (moto.nowada@sdu.edu.cn), Qiu-Gang Zong (qgzong@pku.edu.cn), |

29 Quan-Qi Shi (sqq@sdu.edu.cn)

30 Abstract

The terrestrial magnetosphere is perpetually exposed to, and significantly deformed by the 31 Interplanetary Magnetic Field (IMF) in the solar wind. This deformation is typically detected at 32 33 discrete locations by space- and ground-based observations. Earth's aurora, on the other hand, is a globally distributed phenomenon that may be used to elucidate magnetospheric deformations 34 caused by IMF variations, as well as plasma supply from the deformed magnetotail to the high-35 latitude atmosphere. We report the utilization of an auroral form known as the transpolar arc 36 37 (TPA) to diagnose the plasma dynamics of the globally deformed magnetosphere. Nine TPAs examined in this study have two types of a newly identified morphology, which are designated as 38 "J"- and "L"-shaped TPAs from their shapes, and are shown to have antisymmetric 39 morphologies in the Northern and Southern Hemispheres, depending on the IMF polarity. The 40 TPA-associated ionospheric current profiles suggest that electric currents flowing along the 41 magnetic field lines (Field-Aligned Currents: FACs), connecting the magnetotail and the 42 ionosphere, may be related to the "J"- and "L"-shaped TPA formations. The FACs can be 43 generated by velocity shear between fast plasma flows associated with nightside magnetic 44 45 reconnection and slower background magnetotail plasma flows. Complex large-scale TPA FAC 46 structures, previously unravelled by an Magnetohydrodynamic (MHD) simulation, cannot be elucidated by our observations. However, our interpretation of TPA features in a global context 47 facilitates the usage of TPA as a diagnostic tool to effectively remote-sense globally deformed 48 49 terrestrial and planetary magnetospheric processes in response to the IMF and solar wind plasma 50 conditions.

51

Keywords: Nightside Distorted Transpolar Arc; Solar Wind-Magnetotail-52 Ionosphere/Atmosphere Coupling; Diagnosis; Magnetospheric Magnetotail Magnetic 53 54 Reconnection; Plasma Flow Shear; Field-Aligned Currents

55 Plain Language Summary

In magnetospheric physics, the aurora is one of the most important phenomena in qualitatively 56 and quantitatively understanding the transfer of plasma and energy from the solar wind to the 57 58 high-latitude atmosphere via terrestrial and other planetary magnetospheres. To understand the global picture of the plasma supply from the terrestrial magnetosphere, deformed by the 59 Interplanetary Magnetic Field (IMF) in the solar wind, to the auroral zone, the formation process 60 of a new morphology of auroral transpolar arc (TPA) is investigated in this study. The source of 61 62 these TPAs can be the electric currents flowing along magnetic field lines, induced by the plasma flows in the magnetosphere. The conventional TPA has a straight bar shape, which connects the 63 64 nightside and dayside of the auroral oval. The new TPA morphologies, on the other hand, have significant "distortions" toward pre- and post-midnight at their nightside ends, which may be 65 caused by magnetic field line twisting and magnetosphere deformations due to the action of the 66 IMF. Our results facilitate a paradigm shift in understanding the implications of TPA structure 67 on global scale dynamics in the deformed magnetosphere, and as such, the usage of the auroral 68 TPA shape as a tool to diagnose global-scale magnetospheric effects. 69

70 1. Introduction

The terrestrial magnetosphere, which dynamically changes through interactions with the high-71 speed plasma streams and Interplanetary Magnetic Field (IMF) originating from the Sun, 72 73 effectively shields life on Earth from harmful radiation effects associated with these particles (Black, 1967; Glassmeier et al. 2009, 2010; Shi et al. 2013). The geomagnetic field surrounding 74 the Earth also plays a role in preventing the atmosphere from escaping into space (Wei et al. 75 2014). Therefore, it is important to understand the morphologies and dynamics of our terrestrial 76 77 magnetosphere, in particular, the processes by which plasma is supplied to, and released from, the magnetotail and transferred to the high-latitude atmosphere or ionosphere. 78

Significant global magnetospheric effects are produced not only by changes in the IMF north-79 80 south component (IMF- B_z) but also its dawn-dusk component (IMF- B_y). A series of observational studies (Kaymaz et al. 1995; Nishida et al. 1995, 1998; Pitkänen et al. 2013, 2015, 81 2017) have found that under dominant IMF-B_v conditions, the magnetotail (plasma sheet) 82 becomes increasingly twisted with down-tail distance, caused by the penetration of IMF-B_v into 83 the magnetotail. Magnetotail deformation and IMF penetration to the magnetotail have been 84 attributed to magnetic reconnection under dominant IMF-B_v conditions (Gosling et al. 1990; 85 Cowley, 1981, 1994; Grocott et al., 2007; Tenfjord et al. 2015, 2018), which causes asymmetries 86 87 in the magnetosphere. Inside the deformed magnetosphere, magnetic reconnection can occur and release energized plasma (electrons) earthward and tailward (Petrukovich et al. 1998; Nagai et al. 88 2001; Angelopoulos et al. 2013; Wang et al. 2020, and references therein). The "source" of 89 auroral arcs, which are frequently seen within the polar cap region (sun-aligned arcs), is 90 91 considered to be the currents flowing along the magnetic field, carried by precipitating energetic plasma (electrons) (see the details in a review by Zhu et al. 1997). These field-aligned electron 92 flows originate from the magnetotail. When magnetic reconnection occurs in the nightside 93 magnetosphere, magnetic energy stored in the magnetotail is converted to particle kinetic energy, 94 95 producing accelerated plasma flows out of the reconnection region as earthward and distanttailward high-speed exhaust jets (e.g., Baumjohann et al. 1989, 1990; Angelopoulos et al. 1992, 96 1994). As a result, localized fast plasma flows associated with reconnection are conveyed along 97 the field lines, and embedded within lower velocity plasma flows of magnetospheric origin in the 98 99 magnetotail. Flow shear across field lines between high and low velocity flow regions generates

electric currents that flow parallel to magnetic field lines, known as Field-Aligned Currents
(FACs) (Hasegawa and Sato, 1979; Birn and Hesse, 1991; Fairfield et al. 1999). Evidence of this
process has been compiled by sparse, spatially discrete ground-based, and space-based magnetic
field and particle observations (Angelopoulos et al. 1996; Fairfield et al. 1999, and references
therein). However, the aurorae seen in the Northern and Southern Hemispheres can be used as a
tool to globally diagnose these magnetospheric processes.

A specific auroral form observed under northward IMF-B_z conditions, the Transpolar arc 106 (TPA), occurs at extremely high latitudes. This is identified as a "bar-shaped" emission within 107 108 the polar cap region, extending from the poleward edge of the nightside auroral oval toward the dayside (Frank et al. 1982). Its formation mechanism and features have been explained in terms 109 of magnetospheric convection and its relationship with the IMF orientation (Fear and Milan, 110 2012a, 2012b). TPA locations depend on the extent of clockwise or counter-clockwise plasma 111 112 sheet twisting (viewed from the magnetotail), which is controlled by the $IMF-B_v$ polarity (i.e. either dawnward or duskward, for clockwise or counter-clockwise twisting) (Tsyganenko and 113 114 Fairfield, 2004; Tsyganenko and Stinov, 2005; Tsyganenko et al. 2015; Cumnock et al. 2002).

The TPA formation model proposed by Milan et al. (2005) is one of the most representative 115 TPA formation models based on nightside magnetic reconnection, and has been applied to 116 explain the developments of many TPAs (Fear and Milan, 2012a, b; Kullen et al. 2005; Nowada 117 et al. 2018). In this model, the TPA growth is attributed to the continual formation of newly 118 closed field lines by magnetotail reconnection, whose location retreats tailward. Several "non-119 120 straight" TPAs were also identified in previous statistical studies (Fear and Milan, 2012a; Kullen et al., 2015), which contrast with the "bar"-shaped TPA (hereafter, referred to as a "regular 121 TPA") previously discussed (Fear and Milan, 2012a, b; Kullen et al. 2005; Nowada et al. 2018). 122 However, neither the physical mechanism for these TPAs, nor their implications on the IMF-123 deformed magnetospheric dynamics have been discussed. 124

In this paper, we first identify a new morphological type of nightside distorted TPA, which is distinct from the "regular" TPA. Utilizing space-borne images and in-situ magnetotail observations, together with ground-based geomagnetic field and high-frequency (HF) radar observations, we obtain a global picture of the plasma supply from the deformed magnetotail to the high-latitude atmosphere (auroral zone) by considering the implications of these observations in the nightside distorted TPA formation. In so doing, we demonstrate that the nightside distorted
 TPAs can be used as a remote-sensing diagnostic tool for global magnetospheric effects.

132

133 **2. Instrumentation**

134 New morphological TPAs discussed in this paper were identified using a large database spanning 5 years of auroral observations from 2000 to 2005 by the Wideband Imaging Camera 135 (WIC), which is part of the Far Ultraviolet (FUV) instrument (Mende et al. 2000a, b, c) onboard 136 Imager for Magnetopause-to-Aurora Global Exploration (IMAGE), launched in March, 2000. 137 IMAGE FUV-WIC imaged the aurora in a broad wavelength range from 140 nm to 190 nm, with 138 a cadence of 2 minutes. From this database, we chose 9 nightside distorted TPAs based on visual 139 140 inspection, which were clearly imaged in the plots of the IMAGE FUV-WIC data after the removal of dayglow and background contamination, as described below. 141

The IMAGE FUV-WIC data frequently includes optical contamination, such as sunlight 142 (dayglow) and instrumental optical noise. These non-auroral signals are removed as much as 143 possible from the original WIC images by least squares fitting techniques. The image is 144 separated into two parts along the terminator, but still has an overlap between the dayside and the 145 nightside parts (i.e. the nightside part extends somewhat over the dayside and vice versa). The 146 dayside part of the image is fitted using a two-dimensional Fourier series while the nightside part 147 148 is fitted using a two-dimensional polynomial. The auroral emission is excluded from the fitting process, and the overlap region is used to produce a smooth merging of both parts. The fitted 149 150 glow and background are interpolated over the auroral region (including over the transpolar arc region, which is excluded from the fitting process as well). Subtraction of the fitted signal from 151 the images taken and observed over the whole Earth extracts only the auroral signals, albeit with 152 some unavoidable noise contamination. The light from stars, which occasionally appear over the 153 154 limb of the Earth, can be somewhat scattered and leave their traces on the images. Note that 155 these optical effects are hard to remove, and must not be confused with a real emission from the upper atmosphere. Because non-auroral signals generated by the bright dayglow can only 156 approximately be represented by this method, the optical contamination cannot be completely 157 cleaned from the image. In this study, we discuss the characteristics of unique TPA 158 morphologies, identified based on significant auroral signals, which were extracted through these 159

161

162 **3. Results**

163 **3.1 Overview of Nightside Distorted TPAs**

"Regular" TPAs generally have a straight shape connecting the nightside and dayside auroral 164 oval. In contrast, all TPAs discussed in this paper have a significant "distortion" at the nightside 165 ends (hereafter, referred to as "nightside distorted TPAs"). Figure 1 shows false color images of 166 8 representative nightside distorted TPAs, which were identified from IMAGE-FUV-WIC 167 observations. The top (bottom) row of panels correspond to cases of IMF- $B_v < 0$ (IMF- $B_v > 0$), 168 and the first three columns show Northern Hemisphere (NH) observations, while the last column 169 displays Southern Hemisphere (SH) observations. Each panel is oriented such that the top, right, 170 bottom and left sides, corresponding to noon (12 MLT), dawn (6 MLT), midnight (24 MLT), and 171 dusk (18 MLT), respectively. The color scale is expressed in Analogic-Digital Units (ADU), 172 which is proportional to the observed auroral brightness (Mende et al. 2000b). The upper panels 173 (a) to (c) display dawnside TPAs with the nightside ends distorted toward midnight or pre-174 midnight, observed in the NH. Hereafter, we identify these as "J"-shaped TPAs based on their 175 resemblance to the letter "J". In all observed TPAs, the "J"-shaped TPAs in the NH occur during 176 a negative (dawnward) IMF-B_v interval. The bottom panels (e) to (g) show nightside distorted 177 TPAs with the opposite chirality that occurred on the duskside, in which the nightside ends get 178 distorted toward midnight or post-midnight. We identify these as "L"-shaped TPAs based on 179 their resemblance to the letter "L". Panels (d) and (h) show observations in the SH during 180 negative and positive IMF-B_v intervals, respectively. Interestingly, these two panels appear to 181 show the opposite chirality to their NH counterparts under the same IMF conditions, with an 182 "L"-shaped TPA (panel d), and a "J"-shaped TPA (panel h). The detailed growth of these 183 representative four nightside distorted TPAs and corresponding solar wind conditions are shown 184 185 in the Supporting Information (Figure S1).

186

187 **3.2 In-situ Duskside Magnetotail Observations during the Nightside Distorted TPA interval**

All of the "J (L)"-shaped TPAs identified in our study, shown in Figure 1, originate in the

nightside main auroral oval and protrude toward the dayside, indicating that nightside magnetic 189 reconnection plays a significant role in the formation of these TPAs (see the detailed series of 190 figures shown in Figure S1). In-situ magnetotail observations were examined during the 191 nightside distorted TPA intervals. Figure 2 shows a summary plot of the solar wind (observed by 192 Advanced Composition Explorer: ACE), and the magnetotail (observed by Geotail) on March 193 12th, 2002, when the "L"-shaped TPA was detected by IMAGE FUV-WIC. The panels from top 194 to bottom show: the IMF-B_v and -B_z components in GSM coordinates, the solar wind dynamic 195 pressure, the Geotail measurements of the sun-earth (GSM-X), dawn-dusk (GSM-Y) and north-196 south (GSM-Z) magnetic field components in the duskside magnetotail, the associated magnetic 197 field elevation angle, and the ion flow velocity in GSM and Mean Field Coordinates (MFC), 198 which has axes parallel and perpendicular to local magnetic field lines over the 1 hour 40 199 200 minutes time interval between 00:10 UT and 01:50 UT. During this interval, the "L"-shaped TPA intensifications were clearly identified from 00:31:34 UT to 00:58:12 UT and from 201 01:10:29 UT to 01:37:07 UT, which are bracketed by two gold broken lines, and labelled 'LS'. 202 The GSM locations of Geotail when the "L"-shaped TPAs were seen are indicated below the last 203 204 panel of Figure 2(a). The IMF-B_v and -B_z components were oriented roughly duskward (positive) and northward (positive) during both TPA intervals. Associated solar wind dynamic 205 206 pressure showed no significant changes. The large abrupt decreases and increases twice seen in the Geotail-B_x component indicate multiple crossings (four times) of the magnetotail current 207 208 sheet from the Northern to Southern, and from the Southern to Northern Hemispheres, respectively. The variations of associated B_v and B_z components were anti-correlated with that of 209 the B_x component. Particular enhancements of the B_z component and elevation angle, seen in 210 both LS intervals, suggest that the nightside magnetospheric configuration becomes more 211 212 "dipole-like", presumably resulting from a pile-up of the magnetic flux transported from the distant magnetotail. Before the B_z enhancements, the V_x component shows earthward "bursty" 213 enhancements, indicating the occurrence of magnetotail magnetic reconnection at the onset and 214 the initial stage of the two "L"-shaped TPAs. Taking a look at the x-directional components of 215 plasma flow speed parallel and perpendicular to the field lines (V_{parax}, V_{perpx}), this flow burst had 216 a much more dominant field-aligned component (V_{parax}) than the perpendicular flow velocity 217 (V_{perpx}). The second flow bursts seen during the second LS interval also had a strong field-218 aligned velocity. These earthward flow burst profiles also suggest the tailward retreat of the 219

reconnection locations; the V_x component in the first interval had already started to decrease at 220 the onset of the "L"-shaped TPA, and the flow burst velocity during the second "L"-shaped TPA 221 interval was lower than that in first TPA interval (considering that there was little difference in 222 the satellite positions between first and second TPA intervals). The V_z components at the two 223 earthward flow bursts were negative, suggesting that the plasma in the lobe region was flowing 224 225 into the reconnection region in the plasma sheet. Further energized plasma was associated with the fast plasma flows because the temperature abruptly enhanced at the time of the first flow 226 burst, however, a significant temperature enhancement was not seen at the second fast flow event. 227 The magnetic pressure $(B_t^2/2\mu_0; B_t$ is the magnetic field intensity) was higher than the plasma 228 pressure (N_ikT_i) during the two flow burst intervals, indicating that the regions where the two 229 flow burst events occurred may be plasma sheet boundary layer (PSBL). Before the fast flow 230 231 burst, Geotail was situated in the lobe region in the Northern Hemisphere, but detected the fast plasma flow just after its entry to the PSBL. This is because the plasma pressure began to 232 gradually enhance against a slight decrease of the magnetic pressure. After the flow burst, the 233 plasma pressure was higher than the magnetic pressure due to the migration of Geotail to the 234 235 inner plasma sheet (central plasma sheet). The satellite experienced multiple crossings of the current sheet. During the second LS interval, Geotail transiently went out of the plasma sheet, 236 237 and recorded a weaker second flow burst in the PSBL. After the detection of the second flow burst, the satellite returned to the inner plasma sheet. Baumjohann et al. (1988) reported that 238 239 faster plasma flows in the PSBL tend to have dominant field-aligned components, that is, away from the magnetic equatorial plane, which is consistent with our interpretation in which the two 240 flow bursts seen here occurred in the PSBL. Furthermore, these fast plasma flows seem not to be 241 associated with a "plasma bubble" (Chen and Wolf, 1993). If "plasma bubble" structures were 242 243 formed by magnetotail reconnection and resultant fast plasma flows were driven, the plasma flows should have flow velocity components dominantly perpendicular to the magnetic field 244 lines (see figure 3 in Chen and Wolf, 1993). Chen and Wolf (1993) also pointed out that when 245 fast bursts are caused by a "plasma bubble", the ion temperature and the plasma pressure are 246 gradually increased from the onset of fast flow. However, during the presented interval, the 247 observed plasma flow bursts were predominantly field-aligned. The associated temperatures 248 explosively increased (did not increase) in case of the first (second) flow burst, and the plasma 249 pressure enhancements were not seen at both flow onsets. After the first flow burst, the 250

enhancements of the plasma pressure were found because of the satellite entry to the plasmasheet.

Panels (b) and (c) of Figure 2 show zoomed-in plots of the plasma flow velocity and ground-253 254 based magnetic field perturbations measured at two ground observatories close to the TPA, for the first and second plasma flow bursts, respectively. The top two panels in each case show 255 plasma flow velocity components in GSM coordinates and the x-directional components of 256 plasma flow speed parallel and perpendicular to the local magnetic field, and the bottom two 257 panels show ground magnetic field perturbations in the B_N (local magnetic north-south) and B_E 258 259 (local magnetic east-west) components measured at two representative ground magnetic observatories close to the TPA. Detailed information for the ground stations is listed in Table S2 260 in order of geographic latitude. In the first plasma flow burst (panel b), the peaks of the V_x and 261 V_{parax} components, and those in the ΔB_N components are seen at the same time, suggesting that 262 the fast flows associated with magnetotail reconnection may trigger electric currents, and cause 263 the variations of geomagnetic field. During the second flow burst interval, the geomagnetic field 264 265 peaks were not seen as shown in panel (c). Therefore, at this stage of our analysis, it remains unclear whether or not electric currents which would disturb the geomagnetic field were induced 266 by reconnection-associated fast plasma flows in this case. The summary and zoomed-in plasma 267 velocity plots from Geotail observations of the opposite dawnside magnetotail are shown in the 268 Supporting Information (Figure S2). Panel (d) shows the footpoints of the Geotail trajectory 269 during the same time interval as panel (a) (1h 40m from 0:10 UT to 1:50 UT), which were 270 calculated based on the Tsyganenko 96 empirical magnetic field model (Tsyganenko and Stern, 271 1996), and projected onto IMAGE FUV-WIC data on 1:20 UT. The asterisk and diamond denote 272 the start (0:10 UT) and end (1:50 UT) times of the Geotail footpoint trajectory. During the time 273 interval of interest, the Geotail footpoints were located in the region of 74 degrees ~ 75 degrees 274 MLat at ~ 22 hrs MLT, and were close to the "straight bar" part of TPA in the nightside. 275 Therefore, it is expected that the fast plasma flows associated with magnetotail magnetic 276 reconnection, which were observed during this time interval, may play a role in the formation of 277 the nightside distorted TPA. 278

279

3.3 Ionospheric Electric Currents Inferred from The Ground and Direct Evidence for FACs

When a shear in flow velocity exists between reconnection-associated earthward fast flows and 282 slow magnetotail background flows is present, electric currents flowing along the geomagnetic 283 field lines (Field-Aligned Currents: FACs) can be driven (Hasegawa and Sato, 1979; Birn and 284 Hesse, 1991; Fairfield et al. 1999). These FACs are closely related to the auroral phenomena in 285 the high-latitude atmosphere. In order to investigate this current system, an electric current map 286 in the ionosphere is made based on the geomagnetic field variations beneath and in close 287 proximity to the regions of growth of the nightside distorted ("L"-shaped) TPAs. These 288 measurements were conducted using ground-based magnetic field observations from the 289 SuperMAG ground observatory network (Gjerloev, 2012). Figure 3 shows the equivalent 290 ionospheric current (EIC) distributions, projected onto the IMAGE FUV-WIC data in 291 geomagnetic coordinates during time intervals spanning the two earthward flow bursts. The 292 293 electric current maps are derived from the local magnetic north – south (B_N) and east – west (B_E) components of the geomagnetic field perturbations, which were measured at the ground 294 295 magnetic observatories beneath and in close proximity to the growth regions of the nightside distorted TPA. It is well-known that these ground magnetic disturbances are generated by the 296 297 horizontal components of electric currents in the ionosphere (EIC) (Glassmeier et al. 1989). The orientation and scale of FACs can also be estimated based on the EIC distributions (Glassmeier 298 299 et al. 1989; Morretto et al. 1997; Motoba et al. 2003, and references therein). The geomagnetic field perturbations were taken from 50-minute-high-pass filtered B_N and B_E components. Electric 300 301 current orientations were estimated by rotating these geomagnetic field fluctuation components 90 degrees clockwise (e.g., Glassmeier et al. 1989; Moretto et al. 1997). On the maps during both 302 the first (A) and second (B) flow burst intervals, counter-clockwise current vortices were found, 303 as indicated with magenta circular arrows. This counter-clockwise vortex-structured current 304 305 suggests that FACs, oriented from the ionosphere toward the magnetotail, are caused by electron precipitation associated with reconnection-triggered plasma flow bursts, which were observed by 306 Geotail. This result also suggests that the energized plasma (electrons) were conveyed by the 307 magnetotail fast flows from the magnetotail to the ionosphere. The vortex spatial scale appears to 308 be different between first and second interval. In panels (A), a "large-scale" vortex-like current 309 310 structure is discerned by the electric current vectors measured at most observatories, which are mainly located in the dusk sector (westside) of the nightside distorted TPA, while "small-scale" 311 current vortices with a similar rotational sense are indicated on the nightside part of the TPA 312

during the second interval (panels B). Neither vortex current structure showed any poleward
(high-latitude) migration as the "L"-shaped TPA grew to the dayside.

The vortex-like ionospheric current structures, deduced from the geomagnetic field fluctuations, 315 indicate that upward (from the ionosphere toward the magnetotail) FACs play an essential role in 316 the formation of nightside distorted TPAs. To obtain clearer evidence for the presence of FACs 317 associated with the TPAs, we investigated whether or not the DMSP (Defense Meteorological 318 Satellite Program) satellites crossed the TPAs, and could measure the associated magnetic field 319 to extract the current density along the magnetic field lines (FACs). From our 9 TPA events, we 320 found that the DMSP-F13, -F14 and -F16 satellites crossed the dayside straightforward bar-321 shaped part of the "L"-shaped TPA, as seen on 28th October 2003. 322

Figure 4 shows the temporal variations of the current density parallel to the magnetic field lines 323 324 (J_{para}) derived by the DMSP magnetic field data, which are plotted against universal time (UT), and the DMSP tracking information, such as MLat, magnetic longitude (MLON) and MLT. The 325 326 current density can be computed by applying Ampère's law to the magnetic field perturbations, measured just before and after the DMSP-F13 (panel a), -F14 (panel b) and -F16 (panel c) 327 328 crossings of the TPA. More detailed theory and techniques to derive the current density from the magnetic field data are described by Wang et al. (2005) and Lühr et al. (2016). During the DMSP 329 330 crossing interval of each TPA (in each case less than1 minute), bracketed by two magenta broken lines, negative J_{para} values were found. This indicates that upward FACs were flowing out of the 331 332 TPA (Wang et al., 2005). The geomagnetic field measurements on ground showed large- and small-scale counter-clockwise vortex-like current structures beneath and in close proximity to 333 the TPA, and negative J_{para} bays were found during the DMSP TPA crossings. These results 334 indicate that upward FACs are a dominant source of the nightside distorted TPAs. 335

336

337 3.4 Retreat of Reconnection Points

The electric current vortices suggest that FACs may be essential to formation of the nightside distorted TPA. Here, we consider the "growth" of the TPA. According to the conventional model to explain the TPA formation based on nightside reconnection (Milan et al. 2005), which does not take into account the influence of FACs in the TPA formation, the reconnection points should retreat tailward as the TPA grows to the dayside. A summary plot of the Geotail observations shown in Figure 2 has already suggested the tailward retreat of the reconnection

point. To further support this scenario, we examine the geomagnetic field variations associated 344 with the nightside distorted TPAs using ground-based observations. Figure 5(A) shows 345 geomagnetic field observations at several ground magnetic observatories corresponding to the 346 locations beneath or in close proximity to the regions of growth of a nightside distorted TPA 347 ("L"-shaped TPA observed on 12th March 2002). All magnetic field data for the ground 348 observatories were taken from the SuperMAG network (Gjerloev, 2012). Several magenta points 349 labelled with numbers in the IMAGE FUV-WIC plots in panel (b) correspond to similarly 350 labelled locations in geographical map (panel c). Panel (a) in figure 5(A) shows a stack plot of 351 fluctuations in the local (magnetic) north-south geomagnetic field component (ΔB_N) at these 352 observatories, which are shown by blue. The magnetic fluctuations are obtained by the 353 subtraction of the average magnetic field over the time interval of interest from the raw magnetic 354 355 field values. The fluctuation component at each station is plotted upon their averages as indicated by horizontal grey broken lines, and its peak during the "L"-shaped TPA intensification intervals, 356 bracketed by two gold broken lines, is marked by magenta open circle. The plots are sorted in 357 decreasing order of latitude. The magnetic field fluctuation component at the time of panel (b) 358 359 (00:39:45 UT) is indicated by a vertical cyan solid line in the panel (a). The color code of the IMAGE FUV-WIC data in panel (b) is assigned according to ADU. 360

Figure 5(B) shows a scatter plot of the time-delay of the fluctuation peaks in the local 361 (magnetic) north-south magnetic field component (ΔB_N) from the onset times of 5 nightside 362 distorted TPAs at several ground magnetic observatories from geographical low- to high-363 latitudes. The detailed geomagnetic field plots and information on the ground magnetic 364 observatories in the other four TPA events, except for the 12th March 2002 event, are shown in 365 the Supporting Information (Figure S3). All peaks seen in the magnetic field fluctuation 366 components were positive, implying enhancements of FACs flowing out of the ionosphere, that 367 is, downflowing of electrons from the magnetotail. For three of the TPAs (2000/09/22, 368 2001/12/31 and 2002/03/02), the magnetic peaks are clearly seen at later times for observatories 369 with higher latitude, suggesting that the reconnection points (the source regions of the energetic 370 electrons) were retreating further down-magnetotail, associated with the growth of the TPA to 371 372 the dayside. This result supports not only the tail reconnection occurrence but also the retreat of the reconnection points. The average velocity of the reconnection point retreat can roughly be 373 estimated based on the slope of a line of geographical latitude versus the time delay between the 374

magnetic peaks and the TPA onsets. We adopted a value of 1 degree = 110.95 km to convert a 375 unit of geographic latitude (degree) to equatorial distance (km). The estimated reconnection 376 point retreat velocity is summarized in the table in the top-right of the figure. The three TPAs 377 mentioned above, with very apparent reconnection point retreats, had reconnection point retreat 378 velocities within a range between about 1.2 km/s and 3.0 km/s. The others (2000/11/05 and 379 2002/03/12) showed a much faster retreat speed (7.3 km/s and 12.3 km/s) because their magnetic 380 field peaks appeared with much lower time lags, irrespective of the latitudes of the observatory 381 locations. 382

383 **3.5 Persistence of Magnetotail Reconnection During the Northward IMF Interval**

384 We discuss the plasma flows and their patterns in the polar cap region measured by Super Dual Auroral Radar Network High Frequency (SuperDARN HF) radars (Greenwald et al. 1995; 385 Chisham et al. 2007) during the nightside distorted TPA intervals, in order to obtain evidence for 386 the persistence of magnetotail magnetic reconnection even under northward IMF conditions. The 387 SuperDARN radars, which are located in the high-latitude regions in both Northern and Southern 388 Hemispheres, provide line-of-sight ionospheric plasma flow velocity over much of the polar and 389 390 auroral regions. These measurements, particularly obtained from nine SuperDARN radars in the 391 Northern Hemisphere, have been used to produce high-latitude convection maps based on the "Map Potential" technique (Ruohoniemi and Baker, 1998). The line-of-sight velocity vectors are 392 projected onto geomagnetic grids, and fitted to electrostatic potential solutions, which are 393 394 described by a sixth order spherical harmonic expansion. Complementary flow data from a 395 statistical model characterised by upstream IMF conditions (Ruohoniemi and Greenwald, 1996) is used to constrain the construction of the large-scale flow pattern in regions where the radars 396 provide no measurements (Ruohoniemi and Baker, 1998). 397

Figure 6 presents 6 selected 2 minutes integrations of the northern hemispheric plasma flow streamlines and drift velocity vectors during the interval of a nightside distorted TPA ("J"shaped TPA) observed on 31st December 2001. We overlay these flow velocity profiles onto the corresponding IMAGE FUV-WIC auroral imager data. Black regions indicate higher auroral luminosity, and the IMAGE observation time is shown at the top in each panel. The left, bottom and right sides in each panel correspond to 18h, 24h, and 6h in magnetic local time, respectively. The dotted semicircles indicate the magnetic latitude (MLat) range between 60 degrees and 80

degrees. During the growth of the "J"-shaped TPA, westward plasma flows, ranging between 405 0.35 km/s and 0.85 km/s, were locally (although non-continuously) observed at the poleward 406 edge of the midnight-sector main auroral oval, highlighted by magenta ovals. These flows were 407 originally oriented toward the equator, but rotated toward the west at the poleward edge of the 408 main auroral oval. They are highly suggestive of magnetic reconnection in the magnetotail, 409 identified as "Tail Reconnection during IMF Northward and Non-substorm Intervals (TRINNIs)" 410 (Grocott et al. 2003, 2004) under dawnward IMF-By conditions (see the IMF condition shown in 411 Figure S1c) (Milan et al. 2005; Grocott et al. 2003, 2004). Therefore, at least, nightside 412 reconnection was ongoing during the growth of the "J"-shaped TPA even under the northward 413 IMF conditions, and should play a significant role in the nightside distorted TPA formation. 414

415

416 4. Discussion

417 4.1 A Possible Formation Scenario of the Nightside Distorted TPA

The conventional TPA formation model proposed by Milan et al. (2005) is based on the 418 magnetospheric convection of closed magnetic fluxes formed by magnetotail reconnection. The 419 ground-based observations revealed that the reconnection points retreated tailward with the 420 poleward growth of the TPAs. Furthermore, the SuperDARN HF radar detected TRINNIs, which 421 are remote-sensing evidence for persistent magnetotail reconnection under the northward IMF 422 423 conditions, being consistent with the framework of the conventional TPA formation model (Milan et al. 2005). However, our observations show that FACs can be generated by a plasma 424 flow shear between the fast plasma flows triggered by nightside magnetic reconnection and 425 background magnetospheric slow plasma flows, and appear to play an essential role in the 426 formation of nightside distorted TPAs. In Figure 3, counter-clockwise vortex-like ionospheric 427 current structures are detected by ground-based magnetic field observations beneath and in close 428 429 proximity to the growth regions of the nightside distorted TPAs during the plasma flow bursts 430 seen in the magnetotail. The current density component along magnetic field lines derived by the magnetic field perturbations during the DMSP satellite crossings of the TPA show significant 431 negative bays in Figure 4. These observations suggest the presence of upward FACs associated 432 with nightside distorted TPAs. 433

434

Taking into account these observations, we construct a model to illustrate nightside distorted

TPA (in particular, "L"-shaped TPA) formation. Figure 7 displays a schematic diagram of the 435 possible formation process of an "L"-shaped TPA under positive (duskward) IMF-B_v conditions. 436 The main "bar-like" emissions of the nightside distorted TPAs are located on the dusk side under 437 positive IMF-B_v conditions as seen in Figure 1. The location of the "L"- ("J")-shaped TPAs 438 strongly depends on the IMF-B_v sign; the relation between the location of the main TPA part and 439 the IMF-B_v polarity is the same as that for the "regular" TPA (Comnock et al. 2002; Kullen et al. 440 2002) (see the plots of the OMNI and Geotail-measured solar wind data in Figure S1). This 441 model is depicted in terms of the configuration changes of magnetic field lines due to 442 magnetospheric convection, FACs, reconnection-associated plasma flows, and the reconnection 443 point retreat. The closed field lines formed by nightside reconnection are illustrated by thick blue 444 solid curves, and the orange curves indicate the electric currents induced by the plasma flow 445 446 shear between the background slow plasma flows and fast flows originating from magnetotail magnetic reconnection (blue arrows). FACs flowing out of the ionosphere toward the 447 magnetotail constitute the "source" of the nightside distorted TPAs, being consistent with large-448 and small-scale electric current vortices beneath and in close proximity to the growth regions of 449 450 the nightside distorted TPAs, and significant negative bays of the current density component along the magnetic field lines (J_{para}) across the TPA. Magnetotail reconnection continues at the 451 452 point denoted by red dots until the TPA completely forms, and associated closed field lines convect earthward. The reconnection location retreats further tailward from T₀ to T₃, which are 453 454 highlighted by the thick red arrows and the pink-shaded area, as the tip of the TPA approaches the dayside. This is because higher latitude field lines within the TPA have their nightside 455 (equatorial crossing) positions further down-tail. 456

As the reconnection points retreat tailward, the TPA-associated closed flux tubes are 457 contemporaneously twisted clockwise (counter-clockwise), depending on the dawnward 458 (duskward) IMF-By component. Meanwhile, the nightside plasma sheet undergoes an oppositely-459 oriented deformation (Tsyganenko et al. 2015; Tsyganenko and Fairfield, 2004), indicated by 460 inclined red bar in Figure 7. The closed flux tube twisting is caused by the IMF-B_v penetration, 461 which produces "asymmetry" for the magnetic fields in the Northern and Southern Hemisphere, 462 exerting "torque rotation" due to the electromagnetic force (Gosling et al. 1990; Cowley, 1981, 463 1994). This results in the "L"- and "J"-shaped TPAs, corresponding to the ionospheric footpoints 464 of these field lines in the Northern and Southern Hemispheres. 465

Before and during all nightside distorted TPAs examined in this study (listed in Table S1), the IMF-B_z had been dominantly northward, however magnetotail reconnection appears to occur and, at least, persist during the TPA interval. This result is supported by significant enhancements in geomagnetic activity even under strong and persistent northward IMF-B_z conditions (Shi et al. 2012), and indicates that solar wind energy can enter the magnetosphere during the northward IMF intervals.

Zhu et al. (1997) suggests that the FACs associated with polar cap arcs (TPAs) indicate the 472 presence of upward and downward current pairs. Chen and Wolf (1993) proposed a model of 473 474 closure of upward and downward FACs in the dawn and dusk sectors, which are linked with the inertial currents in the magnetotail and the currents perpendicular to the magnetic field line in the 475 ionosphere. In this model, it is considered that the magnetotail-ionosphere FACs were generated 476 by the reconnection fast flows driven by a "plasma bubble". However, in our model, TPA-477 478 associated magnetic field lines are closed by magnetotail reconnection, and FACs, which are the 479 source of the TPA, may be caused by the flow shear due to the reconnection-associated fast 480 plasma flows. This model simply explains that the nightside distorted TPA is comprised of only closed field lines that have been recently generated by nightside magnetic reconnection, and does 481 not include the fate of other regions of closed fluxes which do not significantly contribute to the 482 formation of the nightside distorted TPA. 483

Because the contribution of a "plasma bubble" for the observed fast flows seems to be small or 484 insignificant, as shown in Figure 2, the bubble-associated current closure scenario is not well 485 486 supported. In this study, sufficient data is unavailable to make an ionospheric current map that would reveal the global FAC profile in the Northern (Southern) Hemisphere. A series of studies 487 based on a global MHD simulation (e.g., Tanaka et al. 2004; Watanabe et al. 2013) showed and 488 discussed a large-scale profile of FAC distributions associated with TPA formation. Upward 489 (downward) FACs can be developed in the sector opposite to the downward (upward) FACs, so 490 that closed current systems are formed, but the development processes during the TPA growth 491 are complicated. In particular, Watanabe et al. (2013) showed that multiple current closures, 492 consisting of multiple upward and downward FACs, can be formed during the TPA growth. At 493 this stage, the global FAC structure associated with the nightside distorted TPAs is not yet 494 revealed with in-situ geomagnetic field measurements. This is a problem to be clarified in the 495 future. In Figure S4, the SuperDARN radar data during this "L"-shaped TPA detected counter-496

497 clockwise ionospheric plasma flows in the dawnside in the Northern Hemisphere. These plasma 498 flow patterns (vortex-like plasma flows) may indicate that the clockwise ionospheric currents, 499 that is, downward FACs, can be generated (e.g., Moretto et al. 1997; Motoba et al. 2003). 500 However, in order to reveal the complete current system associated with the nightside distorted 501 TPAs with greater certainty, a more extensive set of geomagnetic field observations is required.

502 4.2 Scale of Electric Current Vortex Associated with Nightside Distorted TPA

The electric current vortices provide indicative evidence of FACs flowing out of the ionosphere 503 504 to the magnetotail. However, their scales were found to be different between two nightside distorted TPA intervals: a large-scale vortex was seen during the first interval, whereas during 505 506 the second "L"-shaped TPA interval, local small-scale vortices were found at the observatories near the TPA. Since these FACs were induced by the plasma flow shear, the velocity of the 507 plasma flows associated with magnetotail reconnection would be a key physical parameter to 508 determine the current vortex scale on the ground. Therefore, the current vortex scale might be 509 roughly proportional to the plasma flow speed. The electric current vortex scale should become 510 smaller, if the energy of plasma (electrons) released by magnetic reconnection in the magnetotail 511 512 was dissipated upon the ionosphere (Tanskanen et al. 2002).

513 4.3 Formation of the Distortions at TPA Nightside Ends

After the onset of nightside reconnection, the reconnection locations retreated tailward as the 514 515 tips of the TPAs (in Northern and Southern Hemispheres) approach the dayside, and apparently become "stagnant points", which are unaffected by magnetospheric convection. Furthermore, the 516 517 closed flux tubes within the nightside distorted TPAs, which are generated by persistent nightside reconnection even under northward IMF conditions, are twisted, associated with the 518 magnetotail deformation. During the growth of nightside distorted TPA under the significant 519 IMF-B_v conditions as the reconnection site moves further tailward, the tail deformation becomes 520 521 larger and associated field lines are also twisted more strongly (Tsyganenko et al. 2015; 522 Tsyganenko and Fairfield, 2004). Significantly, this twisting of field lines, caused by the IMF- B_{y} penetration (Gosling et al. 1990; Cowley, 1981, 1994), gives opposite chirality to the "J"- and 523 "L"-shaped TPAs seen in the Northern and Southern Hemispheres, even though magnetotail 524 magnetic reconnections occur at the "same" locations in the Northern and Southern Hemispheres 525 (see Figure 7). In a previous study (Milan et al. 2005), it was considered that the nightside 526

527 magnetospheric deformation and field line twisting are only important in determining the TPA 528 growth point in the nightside main auroral oval. Our scenario, however, emphasizes that they 529 play an important role in determining not only the TPA morphology but also how the plasma 530 (electrons) released by magnetotail reconnection are supplied to the ionosphere.

- 531
- 532
- 533

534 **5. Conclusions**

In this study, we have demonstrated that investigations of TPA morphology are important in 535 assessing how the energy stored in the deformed magnetotail is released and supplied to the 536 high-latitude atmosphere or ionosphere. In particular, we have shown that the nightside distorted 537 TPA is a good remote-sensing diagnostic tool for monitoring global magnetospheric effects. The 538 fundamental characteristics and the formation scenario of nightside distorted TPAs obtained 539 540 through this study have clear potential for application to other planets. Namely, this study contributes to understanding the roles of the IMF and solar wind plasma in auroral processes, 541 542 which can also occur at other planets of solar system. Hereafter, more detailed observations of the solar wind-magnetosphere-ionosphere coupling are required to better understand the process 543 544 of nightside distorted TPA formation.

545 Acknowledgments

This work is supported by grants of the National Natural Science Foundation of China (NSFC 546 41961130382, 41974189, and 41404131). B.H. is supported by the Belgian National Fund for 547 548 Scientific Research (FNRS). A.W.D is supported by NSFC grant (41774172). A.G. is supported by STFC grant (ST/R000816/1) and NERC grants (NE/P001556/1 and NE/T000937/1). M.N. 549 thanks Anthony T. Y. Lui for constructive and insightful discussion on our obtained results, and 550 for modeling of the nightside distorted TPAs, and Chen-Yao Han for helping to draw Figure 7. 551 552 Also, he thanks Yukinaga Miyashita for helping the MFC/FAC coordinate transformation of the Geotail plasma data, and the calculations of the Geotail footpoints on ionosphere. We thank the 553 554 PIs of the SuperDARN radars for provision of the ionosphere flow data. SuperDARN is funded by the research agencies of Australia, China, Canada, France, Italy, Japan, South Africa, the U. 555 K. and the U. S. For the ground magnetometer data we gratefully acknowledge: Intermagnet; 556 USGS, Jeffrey J. Love; CARISMA, PI Ian Mann; CANMOS, Geomagnetism Unit of the 557 Geological Survey of Canada; The S-RAMP Database, PI K. Yumoto and Dr. K. Shiokawa; The 558 SPIDR database; AARI, PI Oleg Troshichev; The MACCS program, PI M. Engebretson; GIMA; 559 MEASURE, UCLA IGPP and Florida Institute of Technology; SAMBA, PI Eftyhia Zesta; 210 560 561 Chain, PI K. Yumoto; SAMNET, PI Farideh Honary; The IMAGE magnetometer network, PI L. Juusola; AUTUMN, PI Martin Connors; DTU Space, PI Anna Willer; South Pole and McMurdo 562 Magnetometer, PI's Louis J. Lanzarotti and Alan T. Weatherwax; ICESTAR; RAPIDMAG; 563 British Artarctic Survey; McMac, PI Dr. Peter Chi; BGS, PI Dr. Susan Macmillan; Pushkov 564 565 Institute of Terrestrial Magnetism, Ionosphere and Radio Wave Propagation (IZMIRAN); GFZ, PI Dr. Juergen Matzka; MFGI, PI B. Heilig; IGFPAS, PI J. Reda; University of L'Aquila, PI M. 566 Vellante; BCMT, V. Lesur and A. Chambodut; Data obtained in cooperation with Geoscience 567 Australia, PI Marina Costelloe: AALPIP, co-PIs Bob Clauer and Michael Hartinger; SuperMAG, 568 PI Jesper W. Gjerloev; Sodankylä Geophysical Observatory, PI Tero Raita; Polar Geophysical 569 Institute, Alexander Yahnin and Yarolav Sakharov; Geological Survey of Sweden, Gerhard 570 Schwartz; Swedish Institute of Space Physics, Mastoshi Yamauchi; UiT the Arctic University of 571 Norway, Magnar G. Johnsen; Finish Meteorological Institute, PI Kirsti Kauristie. 572

573

574

575 Data Availability

IMAGE FUV-WIC data can be obtained by contacting the corresponding authors (M.N. and 576 B.H.) and can also be accessed from http://image.gsfc.nasa.gov. SuperDARN data is freely 577 provided for scientific research purposes and can be obtained from the SuperDARN data mirror 578 (http://bslsuperdarnc.nerc-bas.ac.uk:8093/docs/) or by contacting any of the SuperDARN PI 579 research groups (http://www.superdarn.ac.uk). All SuperDARN radar data are processed by the 580 software of fitacf v1.2 and make_grid v1.14.er which are part of the Radar Software Toolkit 581 582 (RST v4.2 https://zenodo.org/record/1403226#.Xy0u7y3MxTY). OMNI (ACE) IMF and solar 583 wind plasma were obtained from Coordinated Data Analysis Web (https://cdaweb.sci.gsfc.nasa.gov/index.html/), provided by NASA Goddard Flight Space Flight 584 Center (GSFCs) Space Physics Data Facility. The Geotail MGF and CPI data can be obtained 585 from Data ARchives and Transmission System (DARTS), provided by the Center for Science-586 satellite Operation and Data Archive (C-SODA) at ISAS/JAXA 587 (http://darts.isas.jaxa.jp/about.html.en). The ground magnetic field data used in this paper can be 588 589 downloaded from the SuperMAG website (http://supermag.jhuapl.edu/). We also thank the World Data Centre for Geomagnetism, Kyoto University for accessing the data of AU and AL 590 indices from http://wdc.kugi.kyoto-u.ac.jp/index.html. The triaxial fluxgate magnetometer data 591 of DMSP (Defense Meteorological Satellite Program) with 1 second temporal resolution are 592 593 accessible from the website of the database of the Coupling, Energetics and Dynamics of 594 Atmospheric Regions (CEDAR)/Madrigal (http://cedar.openmadrigal.org/list/ and https://dmsp.bc.edu/html2/dmspssm.html). 595

596 **References**

- Angelopoulos, V., Baumjohann, W., Kennel, C. F., Coroniti, F. V., Kivelson, M. G., Pellat, R., et
 al. (1992), Bursty bulk flows in the inner central plasma sheet. Journal of Geophysical
- 599 Research: Space Physics, 97, 4027 4039.
- Angelopoulos, V., Kennel, C. F., Coroniti, F. V., Pellat, R., Kivelson, M. G., Walker, R. J.,
- Russell, C. T., Baumjohann, W., Feldman, W. C., and Gosling, J. T. (1994), Statistical
- 602 characteristics of bursty bulk flow events, Journal of Geophysical Research: Space Physics, 99(
- 603 A11), 21257–21280, doi:10.1029/94JA01263.
- Angelopoulos, V., et al. (1996), Multipoint analysis of a bursty bulk flow event on April 11,
 1985, Journal of Geophysical Research: Space Physics, 101(A3), 4967–4989,
 doi:10.1029/95JA02722.
- Angelopoulos, V., Runov, A., Zhou, X.-Z., Turner, D. L., Kiehas, S. A., Li, S.-S., Shinohara, I.
 (2013). Electromagnetic energy conversion at reconnection fronts. Science, 341(6153), 1478 –
 1482. https://doi.org/10.1126/science.1236992.
- Baumjohann, W., Paschmann, G., and Cattell, C. A. (1989), Average plasma properties in the
 central plasma sheet, Journal of Geophysical Research: Space Physics, 94(A6), 6597–6606,
 doi:10.1029/JA094iA06p06597.
- Baumjohann, W., Paschmann, G., and Lühr, H. (1990), Characteristics of high-speed ion flows
- in the plasma sheet, Journal of Geophysical Research: Space Physics, 95(A4), 3801– 3809,
 doi:10.1029/JA095iA04p03801.
- Baumjohann, W., Paschmann, G., Sckopke, N., Cattell, C. A., and Carlson, C. W. (1988),
 Average ion moments in the plasma sheet boundary layer, Journal of Geophysical Research:
 Space Physics, 93(A10), 11507 11520, doi:10.1029/JA093iA10p11507.
- Birn, J., and Hesse, M., The substorm current wedge and field-aligned currents in MHD
 simulations of magnetotail reconnection, J. Geophys. Res.: Space physics, 96, A2, 1611 1618,
 (1991).
- Black, D. I., Cosmic ray effects and faunal extinctions at geomagnetic field reversals, Earth
 Planet. Sci. Lett. 3, 225–236, (1967).

- 624
- Chen, C. X., and Wolf, R. A. (1993), Interpretation of high speed flows in the plasma sheet,
 Journal of Geophysical Research: Space Physics, 98(A12), 21409 21419,
 doi:10.1029/93JA02080.
- 628 Chisham, G., Lester, M., Milan, S. E., Freeman, M. P., Bristow, W. A., Grocott, A., McWilliams,
- 629 K. A., Ruohoniemi, J. M., Yeoman, T. K., Dyson, P. L., Greenwald, R. A., Kikuchi, T.,
- Pinnock, M., Rash, J. P. S., Sato, N., Sofko, G. J., Villain, J.-P., and Walker, A. D. M., A
 decade of the Super Dual Auroral Radar Network (SuperDARN): Scientific achievements, new
- techniques and future directions. Surveys in Geophysics, 28(1), 33 109, (2007).
- Cowley, S. W. H., Magnetospheric asymmetries associated with the y-component of the IMF,
 Planet. Space Sci., 29, 79 96, (1981).
- Cowley, S. W. H., Earth's plasma environment: magnetic reconnection and its effect on
 magnetospheric fields and flows, Philosophical Transaction: Physical Sciences and Engineering,
 349, 1690, The Solar-System: A Review of Results from Space Mission (Nov. 15), 237 247,
 (1994).
- Cumnock, J. A., Sharber, J. R., Heelis, R. A., Blomberg, L. G., Germany, G. A., Spann, J. F., and
 Coley, W. R., Interplanetary magnetic field control of theta aurora development, J. Geophys.
 Res.: Space physics, 107(A7), 1108, (2002).
- Fairfield, D. H., Mukai, T., Brittnacher, M., Reeves, G. D., Kokubun, S., Parks, G. K., Nagai, T.,
 Matsumoto, H., Hashimoto, K., Gurnett, D. A., and Yamamoto, T., Earthward flow bursts in
 the inner magnetotail and their relation to auroral brightenings, AKR intensifications,
 geosynchronous particle injections and magnetic activity, J. Geophys. Res.: Space physics, 104,
 A1, 355 370, (1999).
- Fear, R. C. and Milan, S. E., The IMF dependence of the local time of transpolar arcs:
 Implications for formation mechanism, J. Geophys. Res.: Space physics, 117(A03213),
 (2012a).
- Fear, R. C. and Milan, S. E. Ionospheric flows relating to transpolar arc formation, J. Geophys.
 Res.: Space physics, 117, A09230, (2012b).

- Frank, L. A., Craven, J. D., Burch, J. L., and Winningham, J. D., Polar views of the Earth's
 aurora with Dynamics Explorer, Geophys. Res. Lett., 9(9), 1001 1004, (1982).
- Gjerloev, J. W., The SuperMAG data processing technique, J. Geophys. Res.: Space physics,
 117, A09213, (2012).
- 656 Glassmeier, K. -H., Hönisch, M., and Untiedt, J., Ground-based and spacecraft observations of
- traveling magnetospheric convection twin vortices, Journal of Geophyocal Research: Space
- 658 physics, 94, 2520–2528, (1989).
- Glassmeier, K. -H., Richter, O., Vogt, J., Möbus, P., and Schwalb, A., The Sun, geomagnetic
- 660 polarity transitions, and possible biospheric effects: review and illustrating model. Int. J.
- 661 Astrobiol., 8, 147–159, (2009).
- Glassmeier, K. -H., and Vogt, J., Magnetic polarity transitions and biospheric effects, Space Sci.
 Rev., 155, 1-4, 387 410, (2010).
- Gosling, J. T., Thomsen, M. F., Bame, S. J., Elphic, R. C., and Russell, C. T., Plasma flow
 reversals at the dayside magnetopause and the origin of asymmetric polar cap convection, J.
 Geophys. Res.: Space physics, 95(A6), 8073 8084, (1990).
- Greenwald, R. A., Baker, K. B., Dudeney, J. R., Pinnock, M., Jones, T. B., Thomas, E. C.,
- Villain, J. -P., Cerisier, J. -C., Senior, C., Hanuise, C., Hunsucker, R. D., Sofko, G., Koehler, J.,
- Nielsen, E., Pellinen, R., Walker, A. D. M., Sato, N., and Yamagishi, H., DARN/SuperDARN:
- A global view of high latitude convection, Space Sci. Rev., 71(1-4), 761 796, (1995).
- Grocott, A., Cowley, S. W. H., and Sigwarth, J. B., Ionospheric flow during extended intervals
 of northward but B_Y-dominated IMF, Ann. Geophys., 21(2), 509 538, (2003).
- Grocott, A., Badman, S. V., Cowley, S. W. H., Yeoman, T. K., and Cripps, P. J., The influence
- of IMF B_Y on the nature of the nightside high-latitude ionospheric flow during intervals of
- 675 positive IMF B_Z, Ann. Geophys., 22(5), 1755 1764, (2004).
- 676 Grocott, A., Yeoman, T. K., Milan, S. E., Amm. O., Frey, H. U., Juusola, L. et al., Multi-scale
- observations of magnetotail flux transport during IMF-northward non-substorm intervals. Ann.
- 678 Geophys., 25, 1709-1720. doi:10.5194/angeo-25-1709-2007, (2007).
- Hasegawa, A., and Sato, T., Generation of Field Aligned Current During Substorm. In: Akasofu
- 680 SI. (eds) Dynamics of the Magnetosphere. Astrophysics and Space Science Library (A Series of

- Books on the Recent Developments of Space Science and of General Geophysics and
 Astrophysics Published in Connection with the Journal Space Science Reviews), 78. Springer,
 Dordrecht, (1979).
- Heppner, J. P., and Maynard, N. C., Empirical high-latitude electric field models, J. Geophys.
 Res.: Space physics, 92, 4467 4489, (1987).
- Kaymaz, Z., Siscoe, G., Luhmann, J. G., Fedder, J. A., and Lyon, J. G., Interplanetary magnetic
- field control of magnetotail field: IMP 8 data and MHD model compared, J. Geophys. Res.:
- 688 Space physics, 100, A9, 17,163 17,172, (1995).
- 689 Kullen, A., Brittnacher, M., Cumnock, J. A., and Blomberg, L. G., Solar wind dependence of the
- 690 occurrence and motion of polar auroral arcs: A statistical study, J. Geophys. Res.: Space 691 physics, 107(A11), 1362, (2002).
- Kullen, A., Fear, R. C., Milan, S. E., Carter, J. A., and Karlsson, T., The statistical difference
 between bending arcs and regular polar arcs, J. Geophys. Res.: Space physics, 120, (2015).
- 694 Lühr, H., Huang, T., Wing, S., Kervalishvili, G., Rauberg, J., and Korth, H. (2016), Filamentary field-aligned currents at the polar cap region during northward interplanetary magnetic field 695 derived with the Swarm constellation, Ann. Geophys., 34, 901 915, 696 https://doi.org/10.5194/angeo-34-901-2016. 697
- Milan, S. E., Hubert, B., and Grocott, A., Formation and motion of a transpolar arc in response to
 dayside and nightside reconnection, J. Geophys. Res.: Space physics, 110, A01212, (2005).
- Mende, S. B., Heetderks, H., Frey, H. U., Lampton, M., Geller, S. P., Habraken, S., Renotte, E.,
- Jamar, C., Rochus, P., Spann, J., Fuselier, S. A., Gerard, J. -C., Gladstone, R., Murphree, S.,
- and Cogger, L., Far ultraviolet imaging from the IMAGE spacecraft: 1. System design, Space

703 Sci. Rev., 91, 243 – 270, (2000).

- Mende, S. B., Heetderks, H., Frey, H. U., Lampton, M., Geller, S. P., Abiad, R., Siegmund, O.
- H. W., Tremsin, A. S., Spann, J., Dougani, H., Fuselier, S. A., Magoncelli, A. L., Bumala, M.
- B., Murphree, S., and Trondsen, T., Far ultraviolet imaging from the IMAGE spacecraft: 2.
- 707 Wideband FUV imaging, Space Sci. Rev., 91, 271 285, (2000).
- Mende, S. B., Heetderks, H., Frey, H. U., Stock, J. M., Lampton, M., Geller, S. P., Abiad, R.,
- ⁷⁰⁹ Siegmund, O. H. W., Habraken, S., Renotte, E., Jamar, C., Rochus, P., Gerard, J. -C., Sigler, R.,

- and Lauche, H., Far ultraviolet imaging from the IMAGE spacecraft: 3. Spectral imaging of Lyman- α and OI 135.6 nm, Space Sci. Rev., 91, 287 318, (2000).
- Motoba, T., Kikuchi, T., Okuzawa, T., and Yumoto, K. (2003), Dynamical response of the
 magnetosphere-ionosphere system to a solar wind dynamic pressure oscillation, Journal of
 Geophysical Research: Space Physics, 108(A5), 1206, doi:10.1029/2002JA009696.
- Moretto, T., Friis-Christensen, E., Lühr, H., and Zesta, E. (1997), Global perspective of
 ionospheric traveling convection vortices: Case studies of two Geospace Environmental
 Modeling events, Journal of Geophysical Research: Space Physics, 102(A6), 11597–11610,
 doi:10.1029/97JA00324.
- Nagai, T., Shinohara, I., Fujimoto, M., Hoshino, M., Saito, Y., Machida, S., and Mukai, T.
 (2001), Geotail observations of the Hall current system: Evidence of magnetic reconnection in
 the magnetotail, Journal of Geophysical Research: Space Physics, 106(A11), 25929 25949,
 doi:10.1029/2001JA900038.
- Nishida, A., Mukai, T., Yamamoto, T., Saito, Y., Kokubun, S., and Maezawa, K., GEOTAIL
 observation of magnetospheric convection in the distant tail at 200 R_E in quiet times, J.
 Geophys. Res.: Space physics, 100, A12, 23,663 23,675, (1995).
- Nishida, A., Mukai, T., Yamamoto, T., Kokubun, S., and Maezawa, K., A unified model of the
 magnetotail convection in geomagnetically quiet and active times J. Geophys. Res.: Space
 physics, 103, A3, 4409 4418, (1998).
- Nowada, M., Fear, R. C., Grocott, A., Shi, Q. -Q., Yang, J., Zong, Q. -G., Wei, Y., Fu, S. -Y., Pu,
 Z. -Y., Mailyan, B., and Zhang, H., Subsidence of ionospheric flows triggered by magnetotail
 magnetic reconnection during transpolar arc brightening, J. Geophys. Res.: Space physics, 123,
 (2018).
- Petrukovich, A. A., Sergeev, V. A., Zelenyi, L. M., Mukai, T., Yamamoto, T., et al. (1998). Two
 spacecraft observations of a reconnection pulse during an auroral breakup. Journal of
 Geophysical Research, 103(A1), 47–59. https://doi.org/10.1029/97JA02296.
- Pitkänen, T., Hamrin, M., Norqvist, P., Karlsson, T., and Nilsson, H., IMF dependence of the
 azimuthal direction of earthward magnetotail fast flows, Geophys. Res. Lett., 40, 5598, (2013).
- 738 Pitkänen, T., Hamrin, M., Norqvist, P., Karlsson, T., Nilsson, H., Kullen, A., Imber, S. M., and

- Milan, S. E., Azimuthal velocity shear within an earthward fast flow: Further evidence for
 magnetotail untwisting, Ann. Geophys., 33, 245, (2015).
- Pitkänen, T., Hamrin, M., Karlsson, T., Nilsson, H., and Kullen, A., On IMF By-induced dawndusk asymmetries in earthward convective fast flows. In S. Haaland, A. Runov, & C. Forsyth
 (Eds.), Dawn-dusk asymmetries in planetary plasma environments, Geophyscial Monograph
- (Easi), Easin ausi asymmetres in planeary plasma environments, Ceophyselar Mono
- ⁷⁴⁴ Series, 95–106, Hoboken, NJ: John Wiley, (2017).
- Ruohoniemi, J. M., and Baker, K. B., Large-scale imaging of high-latitude convection with
 Super Dual Auroral Radar Network HF radar observations, J. Geophys. Res.: Space physics,
 103(A9), 20797 20811, (1998).
- Ruohoniemi, J. M., and Greenwald, R. A., Statistical patterns of high-latitude convection
 obtained from Goose Bay HF radar observations, J. Geophys. Res.: Space physics, 101(A10),
 21743 21763, (1996).
- 751 Shi, Q. -Q., Zong, Q. -G., Fu, S. -Y., Dunlop, M. W., Pu, Z. -Y., Parks, G. K., Wei, Y., Li, W. -
- H., Zhang, H., Nowada, M., Wang, Y. B., Sun, W. -J., Xiao, T., Rème, H., Carr, C., Fazakerley,
 A. N., and Lucek, E., Solar wind entry into the high-latitude terrestrial magnetosphere during
 geomagnetically quiet times. Nat. Commun., 4, 1466, (2013).
- Shi, X. -F., Zong, Q. -G., and Wang, Y. -F., Comparison between the ring current energy
 injection and decay under southward and northward IMF B_z conditions during geomagnetic
 storms, Sci. China Tech, Sci., 55, 10, 2769 2777, (2012).
- Tanaka, T., T. Obara, and M. Kunitake (2004), Formation of the theta aurora by a transient
 convection during northward interplanetary magnetic field, Journal of Geophysical Research:
 Space Physics, 109, A09201, doi:10.1029/2003JA010271.
- Tanskanen, E., Pulkkinen, T. I., Koskinen, H. E. J., and Slavin, J.A., Substorm energy budget
 during low and high solar activity: 1997 and 1999 compared, J. Geophys. Res.: Space Physics,
 107(A6), 1086, (2002).
- Tenfjord, P., Østgaard, N., Snekvik, K., Laundal, K. M., Reistad, J. P., Haaland, S., and Milan, S.
 E., How the IMF B_y induces a B_y component in the closed magnetosphere and how it leads to
 asymmetric currents and convection patterns in the two hemispheres, J. Geophys. Res.: Space
- 767 Physics, 120, 9368 9384, (2015).

- 768
- Tenfjord, P., Østgaard, N., Haaland, S., Snekvik, K., Laundal, K. M., Reistad, J. P., Strangeway,
 R., Milan, S. E., Hesse, M., and Ohma, A., How the IMF B_y induces a local By component
 during northward IMF B_z and characteristic timescales, J. Geophys. Res.: Space physics, 123,
 (2018).
- 773 Tsyganenko, N. A., and Stern, D. P. (1996), Modeling the Global Magnetic Field of the Large-
- Scale Birkeland Current Systems, J. Geophys.Res.: Space Physics, 101, 27187 27198.
- Tsyganenko, N. A., Andreeva, V. A., and Gordeev, E. I. (2015), Internally and externally
 induced deformations of the magnetospheric equatorial current as inferred from spacecraft data,
 Ann. Geophys., 33, 1 11.
- Tsyganenko, N. A., and Fairfield, D. H. (2004), Global shape of the magnetotail current sheet as
 derived from Geotail and Polar data, J. Geophys. Res.: Space physics, 109(A03218).
- Tsyganenko, N. A., and Sitnov, M. I. (2005), Modeling the dynamics of the inner magnetosphere
 during strong geomagnetic storms, J. Geophys. Res.: Space physics, 110 (A3), A03208.
- Wang, C. -P., Liu, Y. -H., Xing, X., Runov, A., Artemyev, A., and Zhang, X., (2020). An event
 study of simultaneous earthward and tailward reconnection exhaust flows in the Earth's midtail,
 Journal of Geophysical Research: Space Physics, 125, e2019JA027406,
 https://doi.org/10.1029/2019JA027406.
- Wang, H., H. Lühr, and S. -Y. Ma (2005), Solar zenith angle and merging electric field control
 of field-aligned currents: A statistical study of the Southern Hemisphere, Journal of
 Geophysical Research: Space Physics, 110, A03306, doi:10.1029/2004JA010530.
- Watanabe, M., S. Sakito, T. Tanaka, H. Shinagawa, and K. T. Murata (2014), Global MHD
 modeling of ionospheric convection and field-aligned currents associated with IMF By
 triggered theta auroras, Journal of Geophysical Research: Space Physics, 119, 6145–6166,
 doi:10.1002/2013JA019480.
- Wei, Y., Pu, Z. -Y., Zong, Q. -G., Wan, W.-X., Ren, Z. -P., Fraenz, M., Dubinin, E., Tian, F., Shi,
 Q. -Q., Fu, S. -Y., Hong, M. -H., Oxygen escape from the Earth during geomagnetic reversals:
 Implications to mass extinction, Earth Planet. Sci. Lett., 394, 94-98, (2014).

- 796 Zhu, L., Schunk, R. W., and Sojka, J. J. (1997), Polar cap arcs: A review, J. Atmos. Sol.-Terr.
- 797 Phys, 59, 10, 1087 1126.

798 **Figures and Captions**

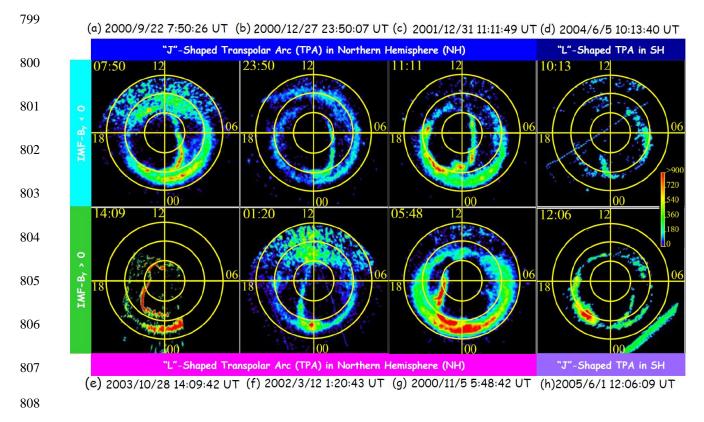


Figure 1: IMAGE-FUV-WIC data plots of selected 8 nightside distorted TPAs are shown. The 809 upper panels (a) to (c) display the "J"-shaped TPAs whose nightside ends are distorted toward 810 midnight or pre-midnight, observed in the Northern Hemisphere under negative (dawnward) 811 IMF-B_v conditions. Panels (e) to (g) show the "L"-shaped TPAs with the nightside ends distorted 812 toward midnight or post-midnight during positive (duskward) IMF-B_v intervals. Panels (d) and 813 (h) show an "L"-shaped, and a "J"-shaped TPAs in the Southern Hemisphere during negative 814 and positive IMF-B_v intervals. These panels are orientated in the same way, with noon 815 (midnight) at the top (bottom), and dusk (dawn) on the left (right) of each plot. The yellow 816 concentric circles show the magnetic latitude (MLat) from 60 degrees to 80 degrees. The color 817 code is assigned according to Analogic-Digital Units (ADU), which is comparable to a detector 818 count rate, being proportional to the observed auroral brightness (accounting for the spectral 819 820 response of the instrument).

- 821
- 822
- 823

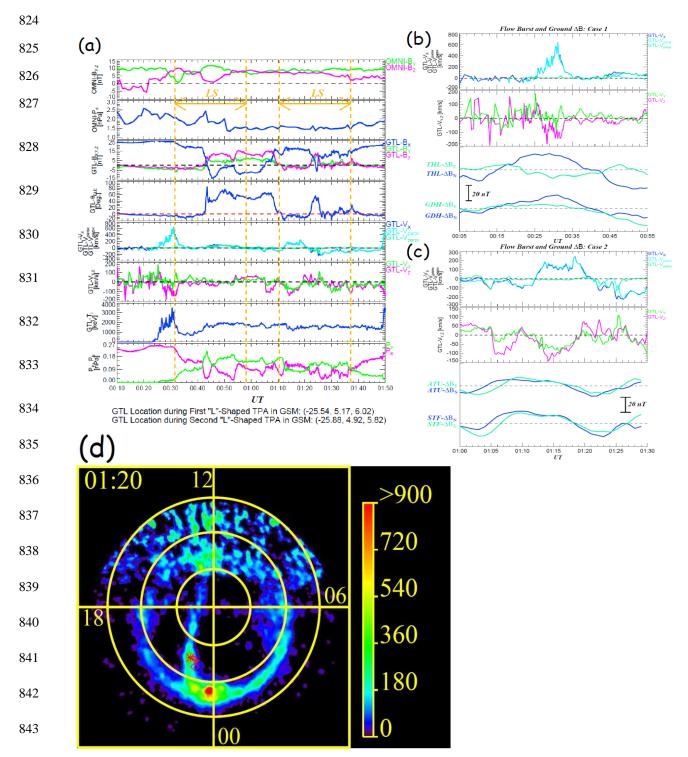


Figure 2: The summary plots of in-situ solar wind, duskside magnetotail and corresponding ground-based magnetic field observations, and the footpoints of the Geotail orbit projected onto the IMAGE FUV-WIC data on 12th March 2002 are displayed. Panel (a) shows a summary plot of OMNI-solar wind and Geotail magnetic field and plasma data in the duskside magnetotail

| 848 | during a 1 h 40 minute interval from 0:10 UT to 1:50 UT. The panels from top to bottom show |
|-----|--|
| 849 | the IMF-B $_{\rm y}$ and -B $_{\rm z}$ components in GSM coordinates, solar wind dynamic pressure, the three |
| 850 | components of the duskside magnetotail magnetic field in GSM, associated magnetic field |
| 851 | elevation angle, GSM-X and x-directional components of the plasma flow velocity parallel and |
| 852 | perpendicular to the local magnetic field, the GSM-Y and -Z components of the magnetotail |
| 853 | plasma velocity, the plasma (ion) temperature, and magnetic and plasma pressures, respectively. |
| 854 | Two clear intensified intervals of the "L"-shaped TPA are each bracketed with two gold broken |
| 855 | lines. Zoomed-in plasma flow velocity in GSM-X and x-directional components of parallel and |
| 856 | perpendicular to the local magnetic field, including significant V_x enhancements which suggest |
| 857 | an earthward plasma flow burst, and corresponding geomagnetic field variations observed at two |
| 858 | representative ground observatories close to the "L"-shaped TPAs are shown in panels (b) and |
| 859 | (c). The geomagnetic field fluctuations are calculated by a subtraction of the magnetic field |
| 860 | average during the presented interval from the observed magnetic field data. Panel (d) shows the |
| 861 | footpoints of the Geotail trajectory during 1 hour 40 minutes from 0:10 UT (asterisk) to 1:50 UT |
| 862 | (diamond), projected onto the IMAGE FUV-WIC data observed on 1:20 UT, using the |
| 863 | Tsyganenko 96 magnetic field empirical model (Tsyganenko and Stern, 1996). |
| 864 | |
| 865 | |
| 866 | |
| 867 | |
| 868 | |
| 869 | |
| 870 | |
| 871 | |
| 872 | |
| 873 | |
| 874 | |
| 875 | |
| 876 | |
| 877 | |
| 878 | |

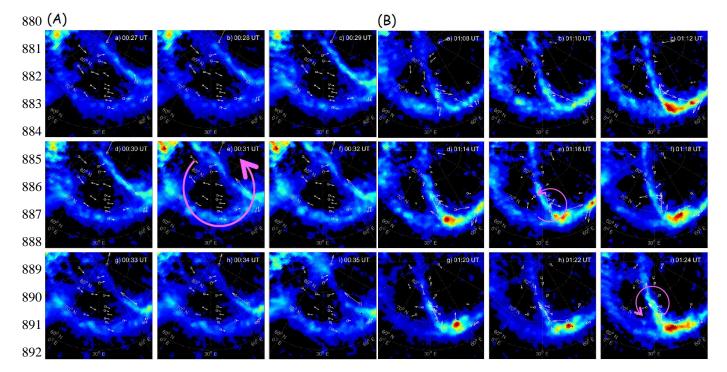


Figure 3: The vortex-like electric current structures detected by ground magnetic observatories beneath and in close proximity to the growth region of the "L"-shaped TPA from 00:27 UT to 00:35 UT with one minute time-step (panels A), and from 01:08 UT to 01:24 UT with two minutes time-step (panels B) on 12th March 2002 are shown. The electric current vectors are derived based on the ground magnetic field fluctuations during the time intervals including the first (a) and second (b) plasma earthward flow bursts, projected onto IMAGE FUV-WIC data in geomagnetic coordinates. Squares and circles with different sizes denote the polarity (positive and negative) and scale of the vertical directional magnetic field fluctuation component (ΔB_{z}). Magenta circle arrows denote large- and small-scale -clockwise current vortices as seen in panels A and B, respectively.

- / 0 0

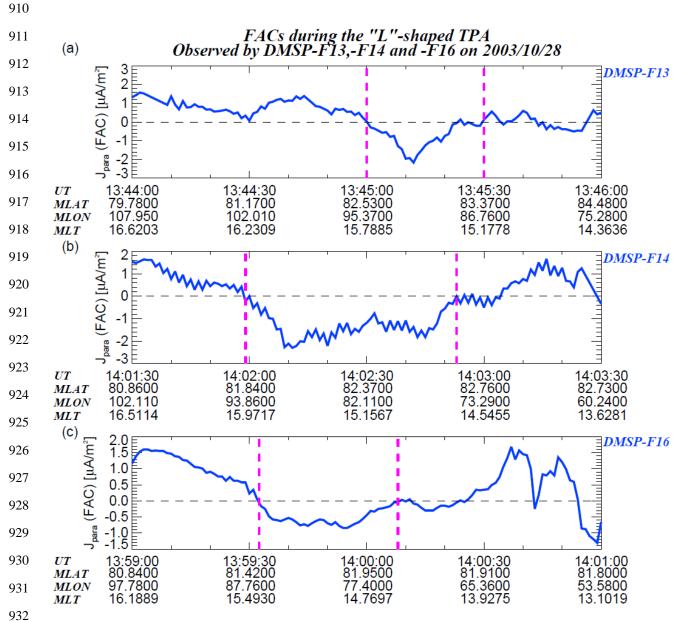


Figure 4: The temporal variations of the current density along the magnetic field lines (J_{para} , i.e., field-aligned current: FAC) against the universal time (UT), magnetic latitude (MLAT), magnetic longitude (MLON), and magnetic local time (MLT) are shown. All J_{para} values are derived from the magnetic field fluctuations observed during the DMSP-F13 (panel a) -F14 (panel b), and -F16 (panel c) crossings of the dayside straightforward bar-shaped part of the "L"shaped TPA, observed on 28th October 2003. The detailed theory and methodology to deduce the J_{para} values from the magnetic field data are given in Wang et al. (2005) and Lühr et al. (2016).

- 940 The TPA crossing time intervals of the three DMSP satellites are bracketed by two magenta
- 941 broken lines.

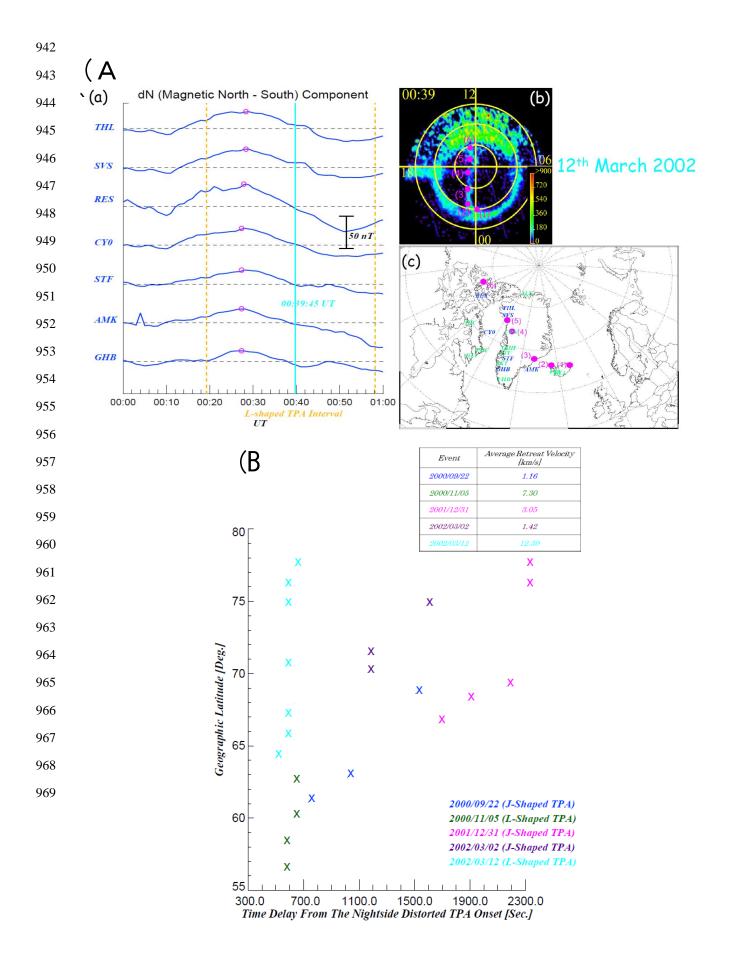


Figure 5: (A) The plots of in-situ geomagnetic magnetic field variations beneath and in close 970 proximity to the growth regions of a nightside distorted TPA ("L"-shaped TPA observed on 12th 971 March 2001) are displayed. Panels (a) to (c) show the magnetic field at several ground magnetic 972 973 observatories corresponding to the locations beneath or in close proximity to the regions of growth of the "L"-shaped TPA. Several magenta points labelled with numbers in the IMAGE 974 FUV-WIC plots (panel b) correspond to similarly labelled locations in geographical map (panel 975 c). Panel (a) shows the plots of fluctuations in the local magnetic north-south magnetic field 976 977 component (ΔB_N) at these observatories highlighted by blue. The fluctuation component, which was obtained by the subtraction of average magnetic field over the time interval of interest from 978 979 the observed magnetic field values at each station, is plotted upon their averages (horizontal grey broken lines), and its peak during the "L"-shaped TPA intensification intervals (vertical gold 980 981 broken lines), is marked by magenta open circle. The plots are sorted in decreasing order of latitude. The magnetic field fluctuation component at the time of panel (b) is indicated by a 982 horizontal solid line in the panel. The color code of the IMAGE FUV-WIC data is assigned 983 according to ADU. (B) The relationship between the magnetic peaks observed at several ground 984 985 observatories beneath and in close proximity to the growth regions of the 5 nightside distorted TPAs from geographical low- to high-latitudes, and the time delays from the 5 TPA onset times 986 to the magnetic peak times is shown. The magnetic field peaks seen in the local magnetic north-987 south magnetic field component (ΔB_N) are used. A rough estimation of the reconnection point 988 989 retreat speed, which was calculated based on the slope of a line of geographical latitude versus 990 the time delay between the magnetic peaks and the TPA onsets, is summarized in the table in the top-right of the panel. We adopted a value of 1 degree = 110.95 km to convert a unit of 991 geographic latitude (degree) to equatorial distance (km). 992

- 993
- 994
- 995
- 996
- 997
- 998
- 999
- 1000

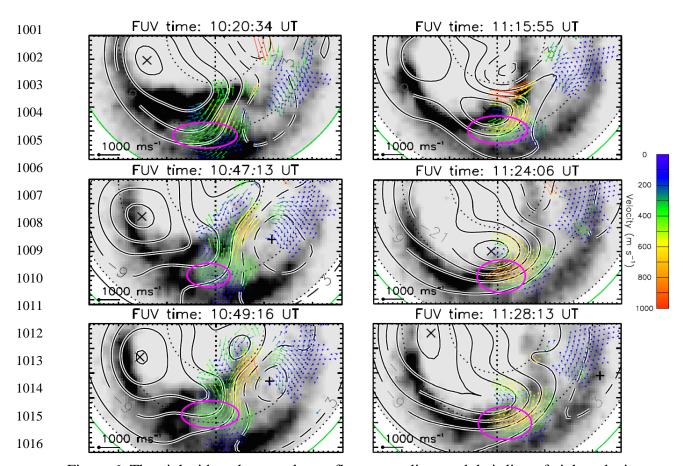


Figure 6: The nightside polar cap plasma flow streamlines and their line-of-sight velocity vectors 1017 1018 measured by SuperDARN in the Northern Hemisphere, overlaid by the IMAGE FUV-WIC auroral image data, are shown. The dotted circles indicate the magnetic latitude (MLat) from 60 1019 1020 degrees to 80 degrees. The left, bottom and right sides in each panel show 18h, 24h and 6h in magnetic local time (MLT), respectively. The time resolutions of the SuperDARN and IMAGE 1021 1022 FUV-WIC data are 2 minutes. These streamlines and velocity vectors are projected onto the geomagnetic grids, and positive (maximum denoted by a plus) and negative (minimum shown 1023 1024 with a cross) electrostatic potential models, which are controlled by the IMF conditions, as 1025 shown with black solid and broken contours on dawn and dusk. The equipotential values are also 1026 overlaid. The green curves show the lower latitude limit of the plasma convection pattern in the polar cap (Heppner and Maynard, 1987), determined from the line-of-sight plasma velocities 1027 1028 measured by the radars. Each dot shows a SuperDARN radar measurement. The length of the 1029 vectors and color code are assigned according to the flow orientation and speed in units of m/s. Westward "Tail Reconnection during IMF Northward and Non-substorm Interval" (TRINNI) 1030 1031 flows are marked with magenta ovals.

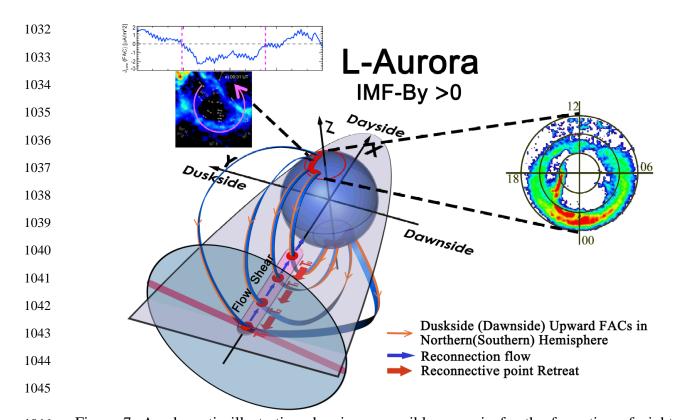
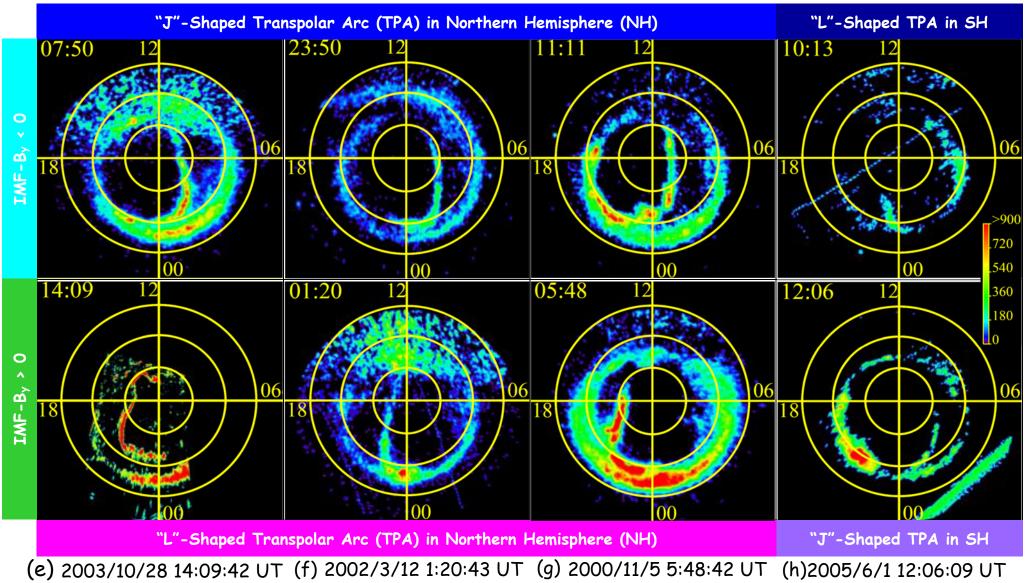


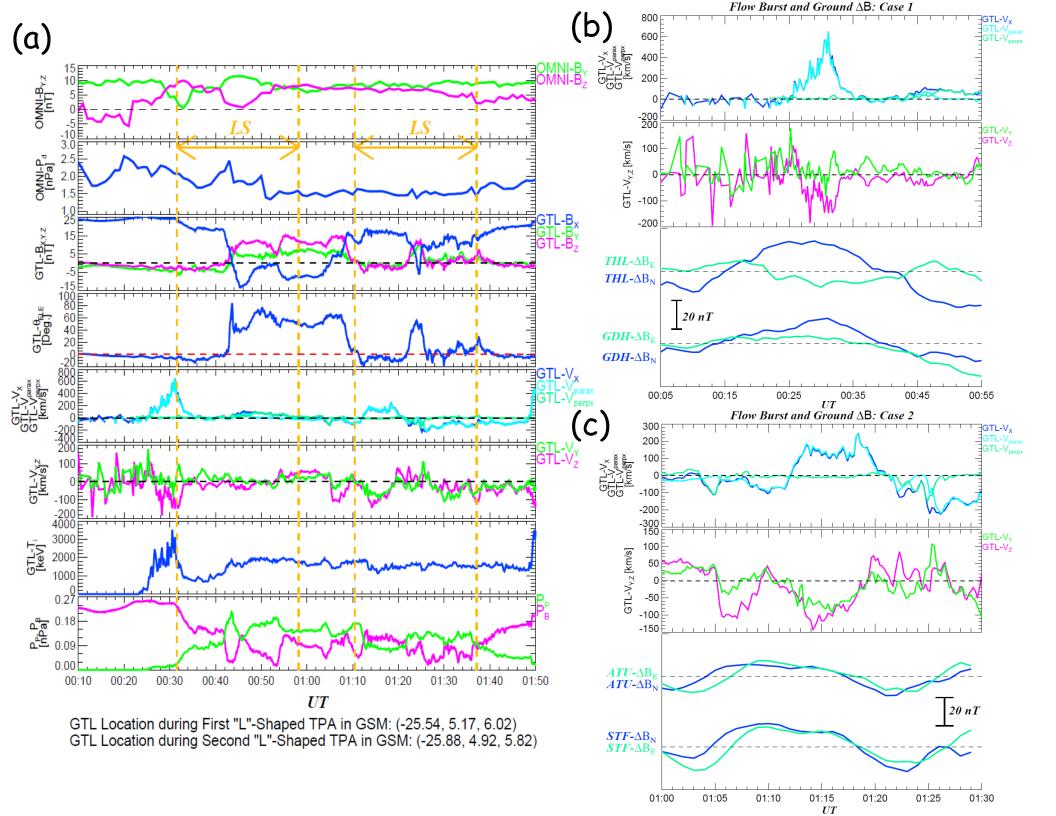
Figure 7: A schematic illustration showing a possible scenario for the formation of nightside 1046 distorted TPAs in terms of the magnetic field configuration changes, field-aligned currents 1047 (FACs), the magnetic reconnection plasma flows and the reconnection point retreats is shown. 1048 This illustration includes the observational examples of the "L"-shaped TPA, obtained by the 1049 IMAGE FUV-WIC on November 5th 2000, the counter-clockwise current vortex in close 1050 proximity to the "L"-shaped TPA on March 12th 2002, induced by FACs flowing out of the 1051 ionosphere, and direct measurement of upward FACs across the "L"-shaped TPA on October 1052 28th 2003, detected by DMSP-F14. The magnetotail cross section and twisted plasma sheet are 1053 1054 shown with a gray-shaded circle and red bar, respectively. FACs flowing toward magnetotail are indicated by orange curved arrows. Thin blue arrows show the fast plasma flows generated by 1055 magnetotail magnetic reconnection. The progressive retreat profile of the reconnection points 1056 1057 (red dots) from T_0 to T_3 is shown with thick red arrows.

Figure 1.



(a) 2000/9/22 7:50:26 UT (b) 2000/12/27 23:50:07 UT (c) 2001/12/31 11:11:49 UT (d) 2004/6/5 10:13:40 UT

Figure 2.



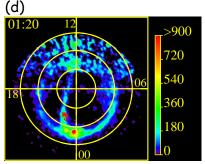
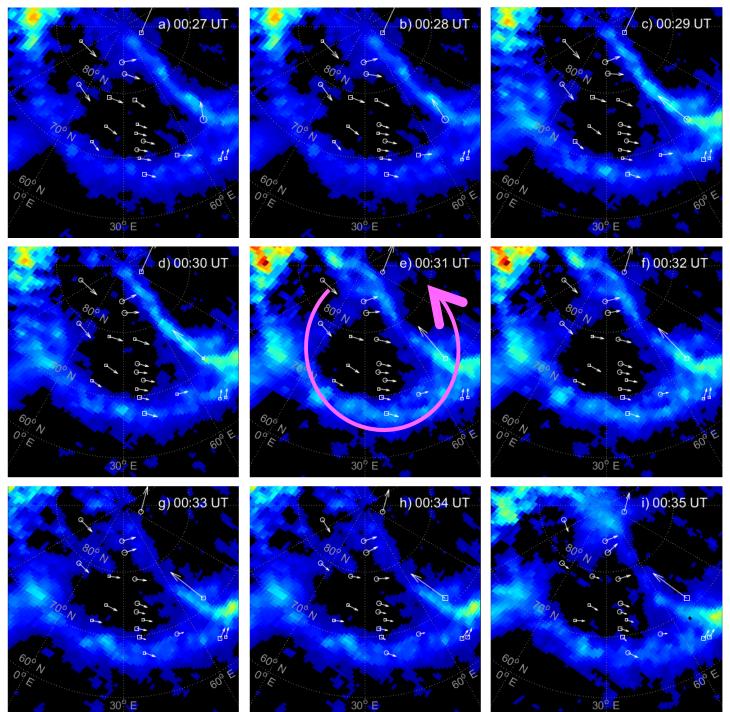
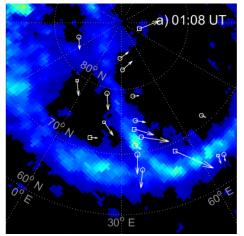


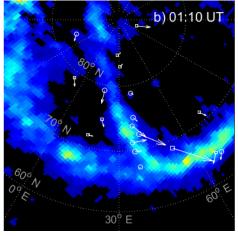
Figure 3.

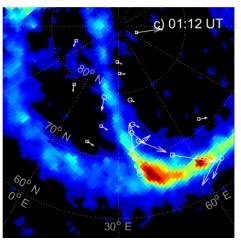
(A)

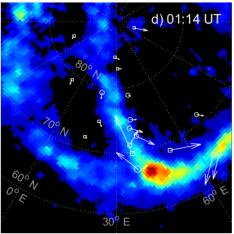


(B)





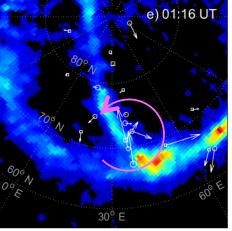


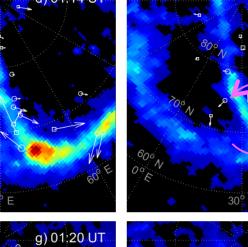


ę. ę

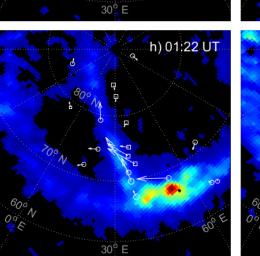
30° E

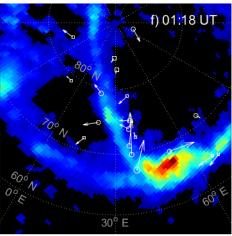
60.°N





60° E





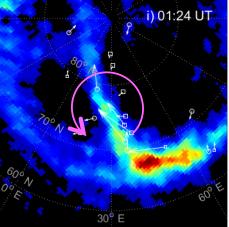


Figure 4.

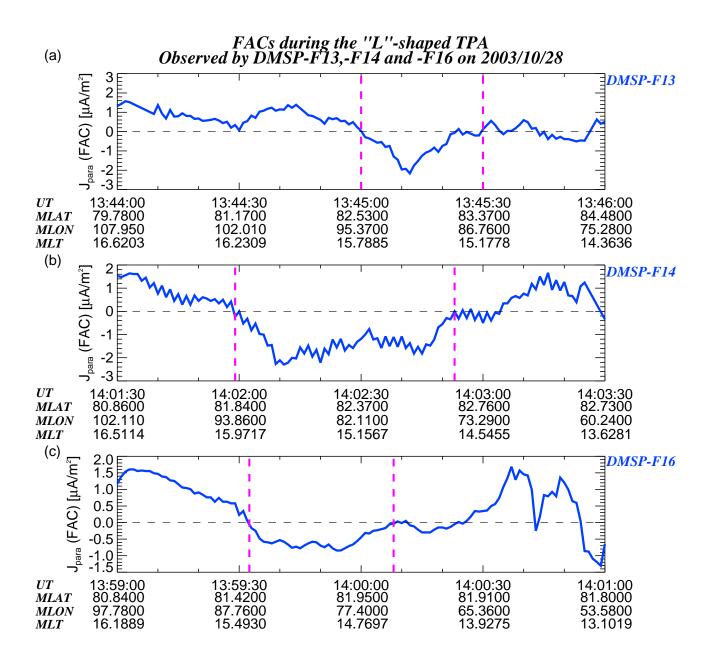
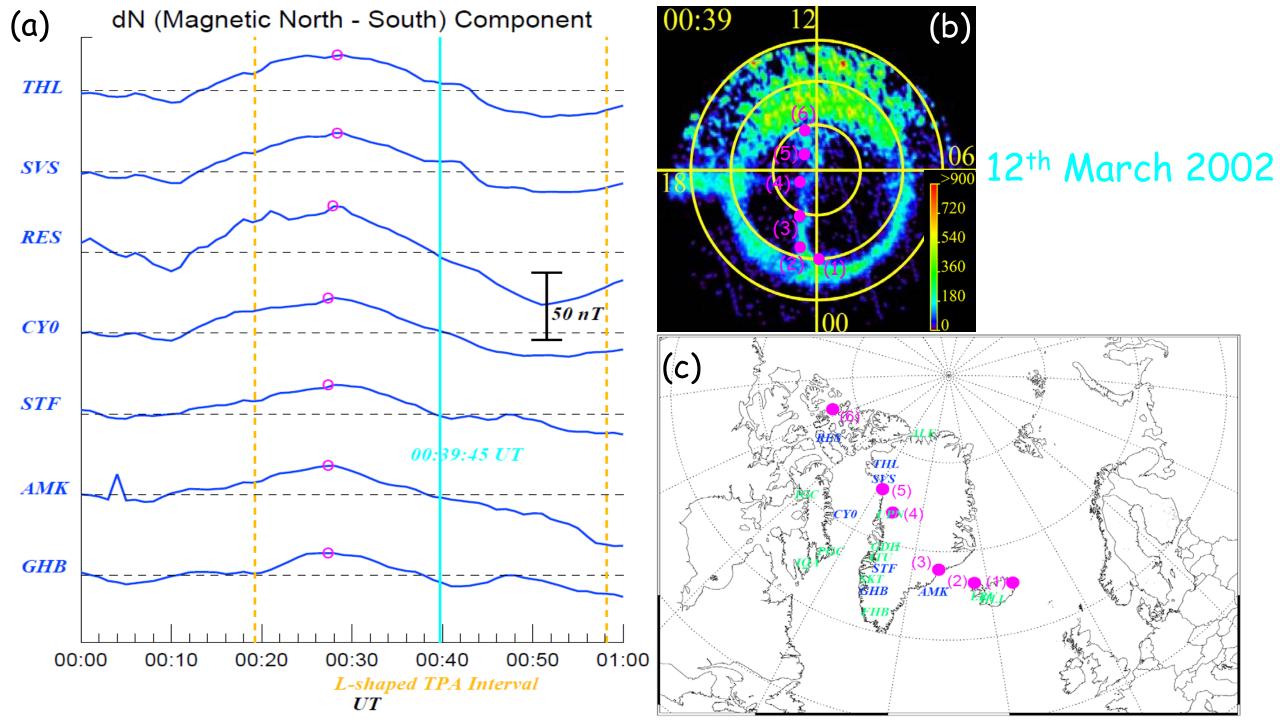


Figure 5.



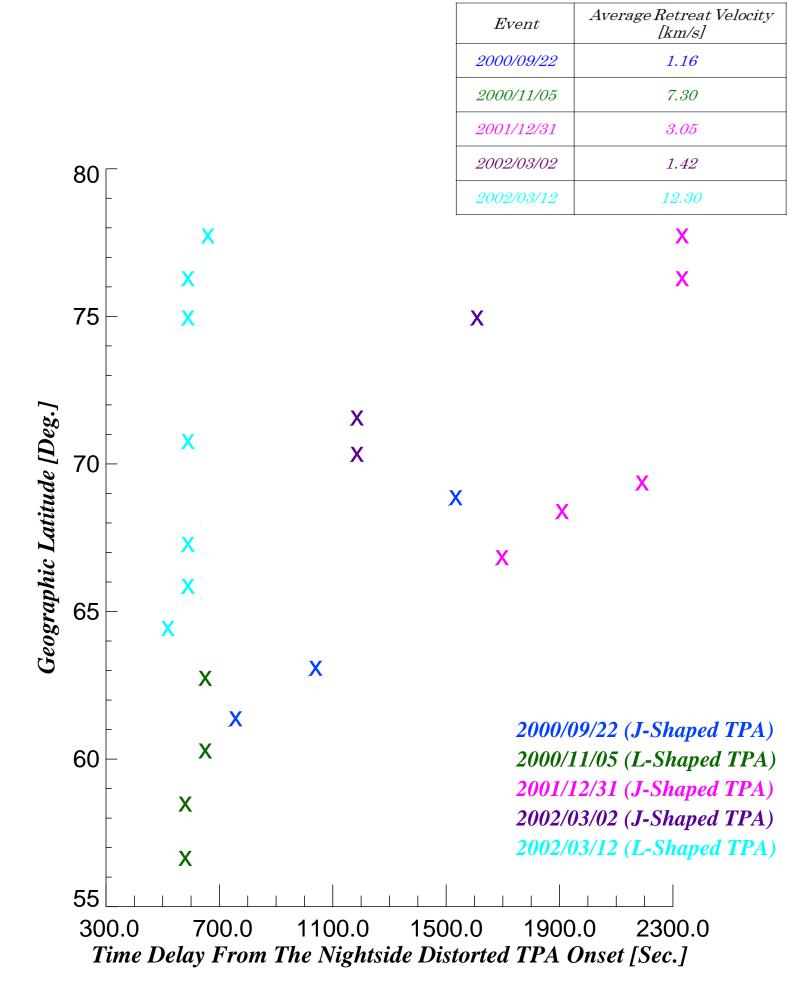


Figure 6.

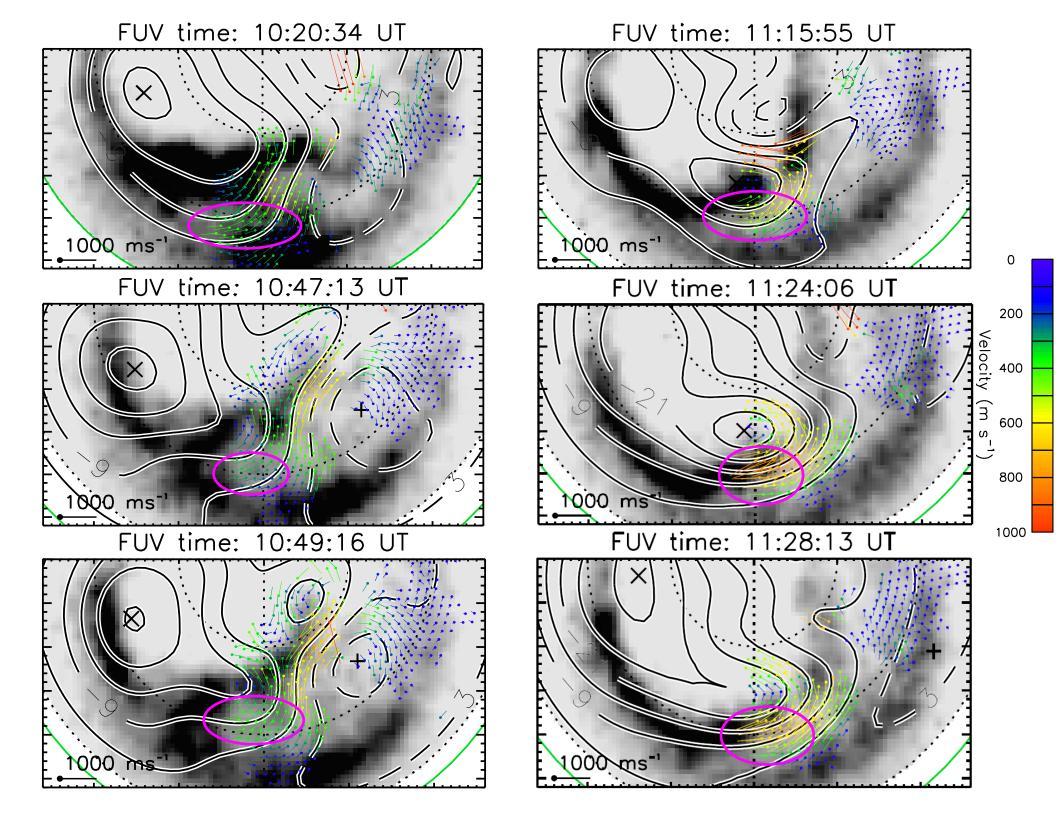


Figure 7.

