

Distribution and properties of magnetic flux ropes in Titan's ionosphere

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Key Points:

- 85 flux ropes are detected in Titan's ionosphere using all of Cassini's flybys of Titan
- Flux ropes are present when Titan is in dynamic areas of Saturn's magnetosphere
- Fitting a force free flux rope model leads to the conclusion that on average they are unlikely to be force free

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Abstract

Titan's magnetic environment is a dynamic and unique place. We detect 85 flux ropes during all the Cassini flybys of Titan from 2005-2017. Analysis describing the location of flux ropes in Titan's ionosphere as well as where Titan is in Saturn's magnetosphere shows that the flux ropes are more often found when Titan is in the noon sector of Saturn's magnetosphere. A secondary peak of occurrence is found in the post midnight area, where it is expected that Saturn's magnetosphere is highly dynamic. We also find that the flux rope occurrence is correlated with the average electron density profile of Titan's ionosphere. A force-free model is utilised to estimate the radii and core magnetic field strength of the flux ropes. We find a large range of radii from 150-500 km with a small number of very large flux ropes (500-1000 km) and a range of core field strengths of 1-15 nT, again with some much larger values (20-40 nT). The model also shows that more flux ropes are right-handed than left-handed in twist, however we are unable to determine if there is a physical reason or if this is due to an observer bias. Additionally, we evaluate the goodness of fit for the model in each instance and conclude that, on average, the flux ropes may be better represented by a non-force-free model.

1 Introduction

Flux ropes are bundles of magnetic flux that are twisted around a central axial core magnetic field with an increasingly tangential field with radial distance from the centre. Flux ropes are ubiquitous in plasma regimes across the solar system and have been found in the ionospheres of the terrestrial planets: Venus and Mars (e.g., Russell & Elphic, 1979; Vignes et al., 2004), in the magnetosphere of Mercury (e.g., Slavin et al., 2010), on the solar surface (e.g., Mouschovias & Poland, 1978) and in the solar wind where they are named magnetic clouds (e.g., Burlaga et al., 1982). Flux ropes are the results of a highly dynamic and interactive magnetic and plasma environment. Flux ropes on the solar surface are thought to be a precursor to coronal mass ejections and solar flares (Chen & Shibata, 2000), and as such flux ropes could be a useful diagnostic tool to uncover the various dynamics of the ionospheres of the terrestrial planets and giant planet moons.

The location of flux ropes around a planetary body can infer their source mechanism, i.e. production at Earth's magnetopause would elicit a conclusion that the solar wind and reconnection are a source mechanism for flux transfer events at Earth (Russell et al., 1985). Similarly, flux ropes in the magnetotail current sheet show that a possible source can be found in tail reconnection (Slavin et al., 2003). Flux ropes at Venus have been theorised to have a number of generation mechanisms, most of which would show a locational bias in detection. For example, if flux ropes were generated by the Kelvin-Helmholtz instability at the ionopause (Wolff et al., 1980) then we might expect an occurrence bias close to the ionopause with properties of flux ropes changing with distance from the ionopause. If they were generated by gravity waves and dynamo action in the ionosphere (Luhmann & Elphic, 1985) then we might expect an occurrence bias close to the "nose" of the induced magnetosphere. In contrast, flux ropes produced by random magnetohydrodynamic flows in the ionosphere (Kleerorin et al., 1994) might not show any bias.

Flux ropes in Titan's ionosphere were first examined by Wei, Russell, Zhang, and Dougherty (2010), where the authors found two unusually large spikes in the magnetic field magnitude in Cassini magnetometer data. The authors compared the magnetic signatures to a force-free flux rope model (e.g., Burlaga, 1988) and concluded that they were flux ropes. Additionally, Wei, Russell, Zhang, and Dougherty (2010) compared Titan's flux ropes to matured flux ropes in Venus' ionosphere. During the T42 flyby, a very pronounced peak in magnetic field magnitude was discovered by Wei et al. (2011) which exceeded all known magnetic field values previously seen at Titan. The authors show that the structure is likely a flux rope, and discuss the possibility of the rope originating from

63 the interaction of Titan with solar wind plasma when Saturn’s magnetopause is pushed
64 back inside Titan’s orbit.

65 Titan is in mildly eccentric orbit around Saturn with a semi-major axis of $1.22 \times$
66 10^6 km ($20.27 R_S$, where $1 R_S = 60268$ km) and an eccentricity of ($\varepsilon = 0.0288$) which
67 gives Titan a $1.17 R_S$ difference in distance from Saturn between perikrone and apokrone.
68 It’s orbit is almost in Saturn’s equatorial plane with an inclination of 0.3485 degrees. The
69 sub-solar position of Saturn’s magnetopause is often found at a bimodal distribution of
70 positions (21 and $27 R_S$ (Pilkington et al., 2015)). During times of high solar wind dyn-
71 amic pressure, it is common to find the magnetopause inwards of $20 R_S$ and hence Ti-
72 tan is often found within Saturn’s magnetosheath and even, on rarer occasions, in the
73 solar wind itself (Bertucci et al., 2015). However, in general, Titan resides inside the mag-
74 netosphere of Saturn where the equatorial region is a varied and dynamic place compared
75 to the solar wind.

76 Saturn’s equatorial current sheet (e.g., C. Arridge, Russell, et al., 2008) is found
77 in all local time sectors unless high solar wind dynamic pressure contracts the magne-
78 tosphere and the current sheet is no longer present in the day side (C. Arridge, Russell,
79 et al., 2008). Additionally, the current sheet has a number of movements that cause the
80 current sheet to be offset from the equator near which Titan lies. The seasonal bowl shape
81 (C. Arridge, Khurana, et al., 2008) causes the whole current sheet to move off the equa-
82 tor increasingly with radial distance due to the solar wind dynamic pressure and the cur-
83 rent sheet mass. Superposed are periodic flapping movements (C. Arridge et al., 2011;
84 Provan et al., 2012), which cause the current sheet to flap at the near the rotation rate
85 of the planet, and aperiodic random waves that cause large disturbances to the current
86 sheet (Martin & Arridge, 2017, 2019). Together these processes act to form a highly dy-
87 namic environment in which Titan resides (Simon et al., 2010, 2013; Achilleos et al., 2014).

88 Titan’s ionosphere itself is mainly formed through magnetospheric electron impacts
89 and solar radiation (e.g., Cravens et al., 2005; Agren et al., 2007), from the thick and
90 extended atmosphere. Titan’s low gravity and dense atmosphere means that the atmo-
91 sphere is extended to over 1000 km in altitude (Yelle et al., 2006), and as such the peak
92 ionospheric electron density is found at altitudes between 1100 - 1200 km (Keller et al.,
93 1992). As Titan has no measurable intrinsic magnetic field (Wei, Russell, Dougherty, et
94 al., 2010), the interaction between the moon and Saturn’s magnetic field is that of an
95 induced draped magnetosphere where Titan’s extended ionosphere ‘captures’ the rotat-
96 ing magnetic field of Saturn to form a draped magnetic field with defined tail and cur-
97 rent sheet (Ness et al., 1982). Similar in form to the draping of the interplanetary mag-
98 netic field around Venus.

99 The same draping process occurs when Titan is in the magnetosheath or solar wind,
100 however the draped field is now sourced from the IMF. Titan is also expected to form
101 a well-defined bow shock when present in the solar wind (Bertucci et al., 2015; Feyer-
102 abend et al., 2016; Omid et al., 2017). Titan’s interaction with Saturn’s magnetopause
103 can also give rise to another magnetic phenomena at Titan - fossil fields (Neubauer et
104 al., 2006; Bertucci et al., 2008). Fossil fields are remnants of magnetic field deposited in
105 Titan’s ionosphere when the magnetic field in a different magnetic environment was draped
106 around Titan (i.e., solar wind plasma and magnetic field can be found in Titan’s iono-
107 sphere when Titan is inside of Saturn’s magnetopause as it was deposited there during
108 a recent excursion into the magnetosheath.) Bertucci et al. (2008) estimates that the life-
109 time of fossil fields are on the order of 3 hours. Wei et al. (2011) discussed the possibil-
110 ity of flux ropes originating, like fossil fields, from a different magnetic environment from
111 the current one, where possible solar wind plasma was found inside the flux rope sur-
112 rounded by magnetospheric plasma.

113 In this paper, we use the entire Cassini magnetometer data set to survey the Ti-
114 tan flybys for flux ropes to contribute to our understanding of Titan’s unique induced

115 magnetosphere. This study presents a location analysis of the 85 flux ropes at Titan, with
 116 emphasis on where on Titan the flux ropes are found along with where Titan is located
 117 in Saturn’s magnetosphere. A force-free model is adapted to fit the magnetic field sig-
 118 natures where radii, axial magnetic field, handedness, and flux content of the flux ropes
 119 is explored along with the restrictions and assumptions of the model used. Finally, a dis-
 120 cussion of the flux ropes will conclude with a comparison with flux ropes at Mars and
 121 Venus.

122 2 Methods

123 Cassini magnetometer data (Dougherty et al., 2004), with a resolution of 1 second,
 124 is utilised to detect peaks in magnetic field magnitude along each of the 127 Titan fly-
 125 bys. Elphic and Russell (1983b, 1983a) stated that a flux rope is a discrete individual
 126 excursion of magnetic field where a peak in magnetic field magnitude is larger than the
 127 surrounding magnetic field. Within this study we define a ‘larger magnitude’ as > 1 nT
 128 due to the average background field being on the order of 1-5 nT and fluctuations within
 129 the background unrelated to flux ropes or environment are < 0.5 nT. Hence, we plot
 130 magnetic field magnitude with altitude to determine if a flux rope is present. An exam-
 131 ple of an altitude plots with a single flux rope on the inbound flybys of T30 is shown in
 132 Figure 1 along with an example of a flyby with two flux ropes on the inbound of T84.
 133 Cassini flybys of Titan are named T(*number of letter corresponding to flyby*) where TA
 134 is the first flyby and T3 is the first flyby after the release of the Huygen probe.

135 We then verify that the magnetic field fits a flux rope signature using minimum vari-
 136 ance analysis (Sonnerup & Cahill Jr, 1967) (MVA) on the smaller time series between
 137 the initial increase in magnetic field magnitude to when it reaches the background level
 138 again. In the ideal case, MVA allows us to determine the directions of maximum, inter-
 139 mediate and minimum variance. At Mercury it was found that maximum variance di-
 140 rection was tangential to the flux rope, the intermediate direction was along the flux rope
 141 axis, and the minimum variance direction was in the radial direction (in cylindrical co-
 142 ordinates). However, in this study we found that at Titan the maximum variance direc-
 143 tion was along the axial direction, thus the intermediate and minimum variance direc-
 144 tions could be combined to give the tangential direction, assuming the radial field is zero.

145 The variance directions are identified as eigenvectors with associated eigenvalues
 146 λ_1, λ_2 and λ_3 for maximum, intermediate and minimum directions respectively. Ratios
 147 of these eigenvalues allow us to infer the ability of the variance analysis to find a distinct
 148 co-ordinate system, where the ratio, $\frac{\lambda_1}{\lambda_2}$, gives the degeneracy of the eigenvectors. As we
 149 are constructing a co-ordinate system, we require each eigenvector to be distinct and non-
 150 degenerate and as such the eigenvalues are required to be distinct, and the ratio gives
 151 a measure of this. We accept eigenvalue ratios of two or more for both ratios, however
 152 most magnetic field time series in this study give eigenvalue ratios of 10 or more, i.e. the
 153 variance directions are distinct and non-degenerate. For comparison, studies at other plan-
 154 etary bodies use an eigenvalue limits of five (Smith et al., 2017) and eight (Briggs et al.,
 155 2011) however, some believe that changes in the acceptance limit does not strongly ef-
 156 fect the results (e.g., DiBraccio et al., 2015).

157 The handedness of the flux rope corresponds to the direction of which the tangen-
 158 tial field curves around the axial field. If a flux rope is right handed and the axial field
 159 is along z then the tangential field is anti-clockwise around z . Conversely, a left handed
 160 flux rope has a clockwise tangential field.

161 3 Spatial Distribution of Flux Ropes

162 At Titan, we find that the flux ropes are detected in the ionosphere. Figure 2 shows
 163 that the highest occurrence rate of flux ropes is around 1050-1100 km which is near the

164 estimated peak of the electron density (e.g., Keller et al., 1992; Edberg et al., 2010; Ågren
 165 et al., 2009) and decreases either side. We also note that due to Cassini’s closest approach
 166 to Titan being above 900 km, we cannot show a lower limit for their detection as they
 167 may occur lower than Cassini has probed, and as such the decrease in occurrence near
 168 900 km is assumed to be due to an observer bias.

169 As described earlier, we know that Titan is exposed to a number of different mag-
 170 netic and plasma regimes as it travels along its orbit. In the noon sector, Titan is likely
 171 to interact with the magnetopause, in the tail and flanks Titan is exposed to the cur-
 172 rent sheet dynamics which can see Titan in either lobe of Saturn’s magnetic field or within
 173 the current sheet (Simon et al., 2010). Figure 3 shows the occurrence rate of flux ropes
 174 normalised to the number of flybys in each local time sector which sample below 2000km
 175 (taken from figure 2) of 1 hour SLT (Saturn Local Time) width. The vast majority of
 176 flux ropes occur in the dayside of Saturn with a smaller number in the 23-3 SLT sector.
 177 The events that make up the 15-20 SLT sector total five events which occur when Ti-
 178 tan is within Saturn’s magnetospheric current sheet.

179 The axial direction for each flux rope, as determined from MVA, allows us to in-
 180 fer the orientation of the flux rope compared to the surface of Titan, where 90° is par-
 181 allel to the surface of Titan and $0/180^\circ$ is perpendicular. We are only interested in de-
 182 termining if the flux rope is parallel or perpendicular to the surface of Titan, hence ori-
 183 entations of $180^\circ - 360^\circ$ are folded into the $0^\circ - 180^\circ$ parameter space. We find a large
 184 spread of orientations with a maximum occurrence around 90° and minima at 0° and
 185 180° . This demonstrates that the flux ropes are found in the plane of the expected draped
 186 field around Titan, e.g. tangential to the surface. Figure 4 shows this orientation bias.
 187

188 The Titan-Ionosphere-Interaction-System (TIIS) is an orthogonal coordinate sys-
 189 tem centred at the centre of Titan where \hat{y} is positive along the Titan-Saturn line, \hat{x} is
 190 along the expected corotation direction, and \hat{z} completes the right handed system. In
 191 the TIIS coordinate system, we can find the position of the flybys with respect to the
 192 expected corotation direction. Figure 5 shows a depression in occurrence where the tail
 193 of Titan’s induced magnetosphere is expected if assuming Saturn’s magnetospheric plasma
 194 population and magnetic field are strictly corotating, i.e., the paucity of flux ropes in fig-
 195 ure 5a, between $0 < Y_{TIIS} < 1$, $-1 < X_{TIIS} < 0.5$). This assumption is as accurate
 196 as possible for the majority of Titan flybys as there are no in-situ measurements of the
 197 velocity and direction of Saturn’s magnetospheric plasma whilst Cassini is in Titan’s iono-
 198 sphere, and the rate of change from when Cassini is sampling Saturn’s magnetosphere
 199 to the sampling of the flux ropes is large and the environment is dynamic in that time
 200 frame.

201 4 Force-free Modeling

202 Force-free assumes that the $\mathbf{j} \times \mathbf{B}$ force density in the flux rope is equal to zero.
 203 Hence, the assumption is that the currents present in the flux rope are negligible or field
 204 aligned, and the magnetic pressure force $\left(\frac{B^2}{2\mu_0}\right)$ is balanced with the magnetic tension
 205 force $\left(\frac{B^2}{2\mu_0 R_c}\right)$ where B is the magnetic field magnitude, μ_0 is the permeability of free
 206 space and R_c is the radius of curvature of the magnetic field. A number of authors have
 207 discussed how a force-free flux rope might represent a flux rope in equilibrium, in its low-
 208 est energy state (Osherovich et al., 1995). As such, a developing or evolving flux rope
 209 would not appear to be force free and the fit of a force-free model to a flux rope would
 210 depend on the "maturity" of the flux rope Wei, Russell, Zhang, and Dougherty (2010).

211 Minimum variance analysis, described earlier, is used to rotate the magnetic field
 212 into a flux rope coordinate system where at Titan we find that the maximum variance

213 direction in Cartesian coordinates gives the axial field direction. The final magnetic field
 214 component (radial) is assumed to be a constant zero in the force-free assumptions. The
 215 three magnetic field components are modeled from Ampere's Law in cylindrical coordi-
 216 nates to give:

$$217 \quad B_A = B_0 J_0(\alpha R) + b_0, \quad (1)$$

$$218 \quad B_T = H B_0 J_1(\alpha R), \quad (2)$$

$$219 \quad B_R = 0, \quad (3)$$

220 where B_0 is the core axial magnetic field, J_0 and J_1 are the zeroth and first order Bessel
 221 functions, b_0 is a magnetic offset in the axial direction. α is a constant of 2.4048 repre-
 222 senting the first root of the zeroth order Bessel function. H is the handedness of the flux
 rope and can take values 1 or -1.

223 R is radial distance from the flux rope centre during the fly-through, we do not di-
 224 rectly know R yet and as such we use a proxy, u . u ranges from -1 at the inbound edge
 225 of the flux rope, to +1 at the outbound edge. Equation 4 can be used to estimate αR :

$$226 \quad \alpha R = \alpha \left(\sqrt{\left(\frac{Y_0}{R_0}\right)^2 + u^2 \left(1 - \left(\frac{Y_0}{R_0}\right)^2\right)} \right), \quad (4)$$

227 where Y_0 is the closest approach distance to the centre of the flux rope, R_0 is the radius
 228 of the flux rope. Hence $\left(\frac{Y_0}{R_0}\right)$ is the impact factor, or ratio of closest approach to ra-
 229 dius which will equal unity at the flux rope edge and zero directly at the centre (Lepping
 230 et al., 2017). The actual radius can be calculated knowing the trajectory, velocity, and
 231 position of the spacecraft relative to the flux rope. We assume at Titan that the flux ropes
 232 are stationary. With this approximation, the radius is found using the following equa-
 tion derived from the geometry of the flux rope encounter:

$$233 \quad R_0 = \frac{V t \sin(\phi)}{2 \sin\left(\cos^{-1}\left(\frac{Y_0}{R_0}\right)\right)}, \quad (5)$$

234 where V is the spacecraft velocity, t is the time in seconds since entry to the flux rope,
 ϕ is the angle between the flux rope axis and the trajectory.

235 Furthermore, one can calculate the flux content, Φ , of the flux rope using the fol-
 236 lowing equation:

$$237 \quad \Phi = \frac{2\pi}{\alpha} B_0 R_0^2 J_1(\alpha) \quad (6)$$

238 The model, consisting of four free parameters B_0 , Y_0/R_0 (impact factor), b_0 and
 239 H , is fitted to the magnetometer data using non-linear least squares, minimising the re-
 240 duced χ^2 using the Levenberg-Marquardt algorithm. With the restriction that H can only
 241 take the values 1 and -1. This restriction is imposed by fitting the model twice, once for
 242 either handedness, and selecting the lowest χ^2 value. Uncertainties are estimated using
 243 the covariance matrix formed during the final iteration of fitting which is calculated us-
 244 ing the Jacobian matrix. The square root of the diagonal elements of the covariance ma-
 245 trix give the standard deviation of each fitted value, and hence a measure of uncertainty
 246 in each parameter. The off-diagonal elements give a measure of the covariance between
 247 each of the fitted parameters, and non-zero values of off-diagonal elements show the vari-
 ables to be fitted are dependant on each other.

248 An example of a force-free model fitted to magnetometer data is found in Figure
 249 6 for a flux rope on T5 in 2005. 49 of the total 85 flux ropes could be fitted using the
 250 force-free model to a satisfactory degree (χ^2_p equal to or below 5 and χ^2 P-value of 0.05
 251 or more). This restriction was to allow the fitting to attempt to fit flux ropes that may

not completely fit the force-free assumptions. The P-value is a probability that, under a null hypothesis (in this case the null hypothesis is that the fitting is no better than a straight line) the χ^2_ν will be greater than the value of χ^2_ν calculated from the data, a value of 0.05 or 5% is a generally accepted significance level. Values of χ^2_ν and χ^2 P-value for each flux rope can be found in supporting information S1.

Figure 7 shows the results of the 49 flux ropes that were fitted to a satisfactory degree, showing the maximum core field strength in a), the flux rope radii in b) and the flux content calculated from these values in c). The binning of the data is unequal to emphasis the spread without a number of null bins, field strength is binned into 1 nT bins from 0-10 nT, then 2 nT bins from 2-40 nT. Radii of the flux ropes is binned as 50 km for 0-500 km and 100 km for 500-1500 km. Flux content is binned as 5 Wb from 0-20 Wb, 20 Wb from 20-200 Wb and 100 Wb from 200-400 Wb.

The main results show that in each parameter of interest there are a large spread where the maximum field strength ranges from 1-15 nT, the radii from 50-500 km and the flux content from 5 Wb with a long tail up to 400 Wb. The spread of data is large and as such an average value will not be quoted as it is not instructive, however, it is noted that the distributions appear to maximise at lower values and have a large tail. This shows that smaller and weaker flux ropes are more common but large and strong flux ropes are also found at a lower occurrence rate. Additionally, the handedness of the ropes is investigated and it is found that 37 are right-handed and 12 are left-handed. Russell (1990) showed that the handedness of flux ropes at Venus corresponded to the location of the flux rope source. If the flux rope formed on the dusk flank of Venus (note that Venus exhibits retrograde rotation) then these flux ropes were more likely to be left-handed, and a right-handed preference is found on the dawn flank. However, there appears to be no statistical significance to the position or orientation of either handedness at Titan (not shown).

5 Discussion

Titan's environment is highly changeable in that the moon can find itself moving between different magnetic and plasma regimes in short time frames. Titan being present in the solar wind and magnetosheath is controlled by both the magnetopause position and hence solar wind dynamic pressure, and where Titan is along its orbit. Figure 3 shows that when normalised to the number of flybys present in each one hr local time bin, the majority of flux ropes are found between 9-14 SLT with a large peak at 12-13 SLT. We thus conclude that the highly dynamic magnetic environment of the dayside magnetosphere, near the magnetopause, may be a source of initiation of flux ropes, this is displayed in Figure 8 which shows the number of flux ropes found in each SLT sector for each magnetospheric environment.

Additionally, Figure 9 shows the occurrence of flux ropes found in each magnetospheric environment. The red bars show the relative occurrence, relative to the occurrence rate of all flux ropes in each environment. We find that a much higher relative occurrence rate is found when Titan is in the magnetosheath or solar wind. The environment classification is found using Simon et al. (2010, 2013); Kabanovic et al. (2017).

A secondary peak in occurrence is found when Titan is in the post-midnight sector (1-3 SLT). This is presumably driven by a general disturbed environment, related to activity in the magnetotail driven by magnetic reconnection (Vasyliunas, 1983) and where we note that the observed (Arridge et al., 2016) and statistical reconnection X-line from plasmoids/dipolarisations (e.g., Thomsen et al., 2013; Smith et al., 2016, 2018) is close to Titan's orbit.

The final peak in occurrence, between 16-19 SLT, consists of five flux ropes that occur only when the moon is inside the equatorial current sheet, shown in figure 8. Ad-

302 Additionally, all local time sectors are sampled except 7-8 SLT, however most notable is that
 303 flybys with no flux ropes detected are all found within Saturn’s magnetosphere and have
 304 very few environment changes as defined by Simon et al. (2010, 2013); Kabanovic et al.
 305 (2017) compared to encounters with one or more flux ropes. These non-detection flybys
 306 are all during times when Titan mainly resides inside either magnetospheric lobe of Sat-
 307 urn (a table of number of flux ropes and environmental changes can be found in the sup-
 308 porting information).

309 We also discuss the altitude distribution of flux ropes when normalised by the num-
 310 ber of hours Cassini spent in each altitude bin. We find that the peak in occurrence is
 311 around 1100-1150 km above Titan’s surface, which is colocated with the expected elec-
 312 tron density peak. There is a large spread in occurrence where the sharp fall at 800-900
 313 km may be an artefact of the Cassini closest approach distances, and that Cassini never
 314 goes below 800 km and as such we cannot give a lower estimate on flux rate occurrence
 315 height. Hence, we can conclude that the flux ropes are occurring where electron density
 316 peaks and their occurrence reduces in correlation with electron density in the ionosphere
 317 above 1100 km.

318 The orientation of the flux ropes, with respect to the surface normal of Titan, is
 319 explored in figure 4, where a broad but pronounced peak in orientation is found around
 320 90° . This means that the flux ropes are lying parallel to the surface of Titan, in the plane
 321 of the expected draped magnetic field. Implying that the mechanism for production is
 322 possibly the interaction of the draped magnetic field and the external environment.

323 A total of 49 flux ropes fitted the force-free flux rope model to a satisfactory de-
 324 gree (MSE equal to 0.5 or below, χ^2 probability of 5% or less). The main results show
 325 that the flux ropes have a large range of core magnetic field values and radii leading to
 326 a large range of flux content. As the spread of data is very large, quoting average val-
 327 ues for the flux ropes at Titan would be misleading. However, they have a range of mag-
 328 netic field strength from 1-15 nT and radii of 50-500 km. Figure 10 shows the radius and
 329 magnetic field information on a positional plot. The figure shows that there is little de-
 330 pendence on position for radius and magnetic field strength for the flux ropes that were
 331 fitted. Non-fitted examples are shown in grey, indicating there is no positional depen-
 332 dence of the effectiveness of fitting a force-free model. Therefore, we conclude that the
 333 position of the flux ropes at Titan is not an effective indicator of which mechanisms may
 334 be inducing their formation.

335 The closest comparison of ionospheric flux ropes may be found in Venus’ ionosphere,
 336 where flux ropes are frequently observed. The flux ropes here are generally on the or-
 337 der of 10s km in size and varied field strengths inside the flux ropes that range from 1-
 338 100 nT (e.g., Elphic et al., 1981; Elphic & Russell, 1983b; Luhmann & Elphic, 1985; Klee-
 339 orin et al., 1994; Ledvina et al., 2002; Wei, Russell, Zhang, & Dougherty, 2010). Zhang
 340 et al. (2012) shows a secondary population of giant flux ropes on the order of 100s km
 341 in size and a flux content of 1000s Wb. In comparison, the flux ropes at Titan have sim-
 342 ilar spatial scales to the larger population described by Zhang et al. (2012), however have
 343 much lower magnetic field strength values on average, attributed to the lower surround-
 344 ing magnetic field strength.

345 At all points of this study, we have attempted to reduced the bias produced by Cassini’s
 346 trajectory, however as described earlier when discussing the altitude occurrence, it is some-
 347 times difficult to remove all biases due to position of Cassini in TIIS coordinates or the
 348 changing value of closest approach to the moon. As described, Cassini encounters Ti-
 349 tan in every local time hour except 7-8 SLT, however each encounter has a different alti-
 350 tude profile and only shows one *slice* through Titan’s ionosphere. Hence it is cautioned
 351 that averaging values from all flybys is not an accurate representation of the whole sys-
 352 tem and as such we have shown a large range of variables for flux rope at Titan.

Improvements to modelling flux ropes, and the different models that can be fitted are explored in a companion paper (Martin et al., Submitted) where elliptical cross sections and bent flux ropes are also explored to aid the fitting of the force-free model and ultimately understand the spacial and temporal evolution of flux ropes. The force-free model describes a flux rope at equilibrium, we hence conclude that the force-free model shows a mature and stable flux rope. Other models, such as the non-force free Hidalgo et al. (2002) model, may be able to more accurately describe a developing and changing flux rope with time and hence why we see a reduced number of flux ropes at Titan being adequately fitted by a force-free model.

6 Summary

85 flux rope signatures are detected in Cassini magnetometer data during all 126 flybys of Titan, where numerous flybys have multiple flux ropes along Cassini's trajectory. We discuss the position of Titan in Saturn's magnetosphere and find that, when normalised to the number of flybys in each local time sector, a higher occurrence is present in the noon (9-15 SLT) sector where according to the classification criteria of Simon et al. (2010), Titan experienced both current sheet and magnetosheath magnetic field. Fluctuations between magnetic field regimes can occur over timescales of less than a few minutes and can last indefinitely depending on the location of Titan and Cassini. Additionally, we find that a secondary occurrence peak is present in the post-midnight area where it is expected to be a turbulent and dynamic magnetic field regime. Additionally, the flybys found with no flux ropes are found only when the magnetic field is lobe-like or has very few regime changes - i.e. enters the current sheet.

The position of the flux ropes with respect to Titan's surface show a correlation with electron density, and large trajectory and sampling biases. Additionally, the orientation is explored using MVA and found that the flux ropes lie parallel to the Titan's surface, implying that the ropes are initiated by the interaction of draped field with surrounding magnetic field, or alternatively have a similar source to fossil fields.

A force-free flux rope model is examined and fitted to the magnetometer data where the model shows a large range of flux rope radii and central magnetic field strengths of the orders of 50-500 km and 1-15 nT respectively. There appears to be no spatial correlation with handedness, size or field strength however this may be a result of a lack of data. We showed that a large number of flux ropes were unable to be fitted with the force-free flux rope model and additional studies of non-force-free models and changes to the force-free model will be examined in the companion paper Martin et al. (Submitted).

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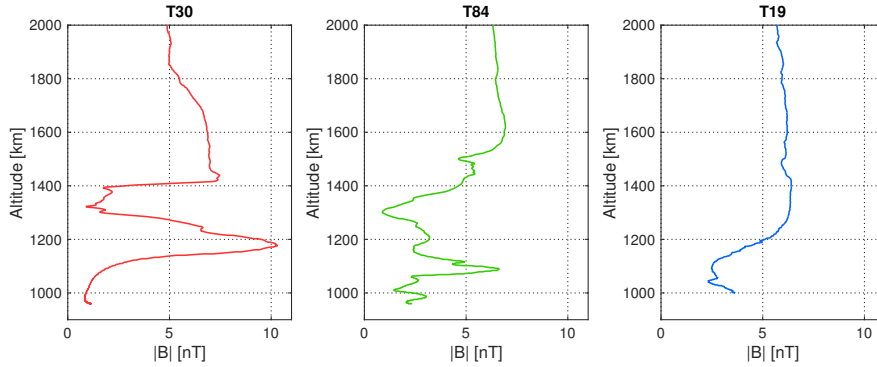


Figure 1. Figure showing three examples of altitude profiles of magnetic field magnitude. Red: T30 with a large flux rope at 1150 km when Titan is at 13.6 SLT. Green: T84 flyby with flux rope at 1100 km when Titan is at 13.7 SLT. Both of these profiles are taken from times when Titan is in the magnetosheath or current sheet according to the classification scheme from (Simon et al., 2010, 2013). Blue: T19, an example with no flux ropes when Titan is at 2.2 SLT and primarily in the southern lobe of Saturn according to (Simon et al., 2010).

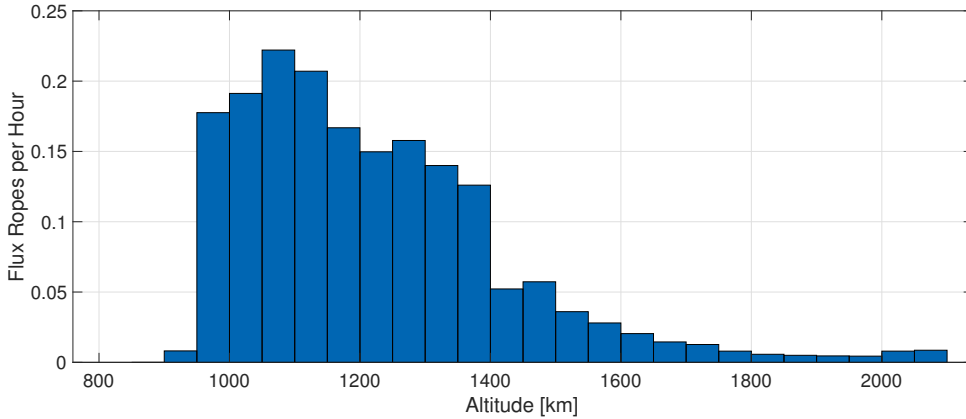


Figure 2. Figure showing the altitude distribution of flux ropes in Titan's ionosphere normalised by the number of hours that Cassini spent within the altitude bin for all flybys.

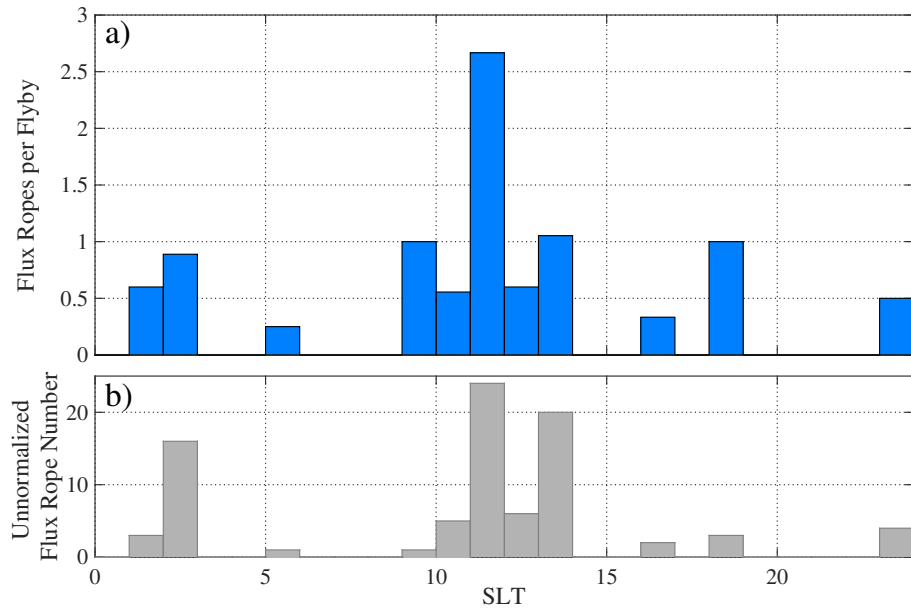


Figure 3. Figure showing the Saturn local time distribution of flux ropes in Titan's ionosphere, a) normalized by number of flybys in each SLT sector and b) unnormalized.

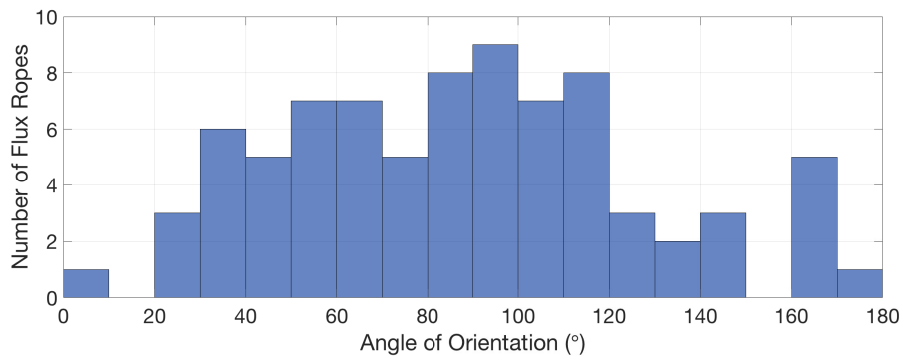


Figure 4. Figure showing the orientation of flux ropes at Titan with respect to the surface of Titan.

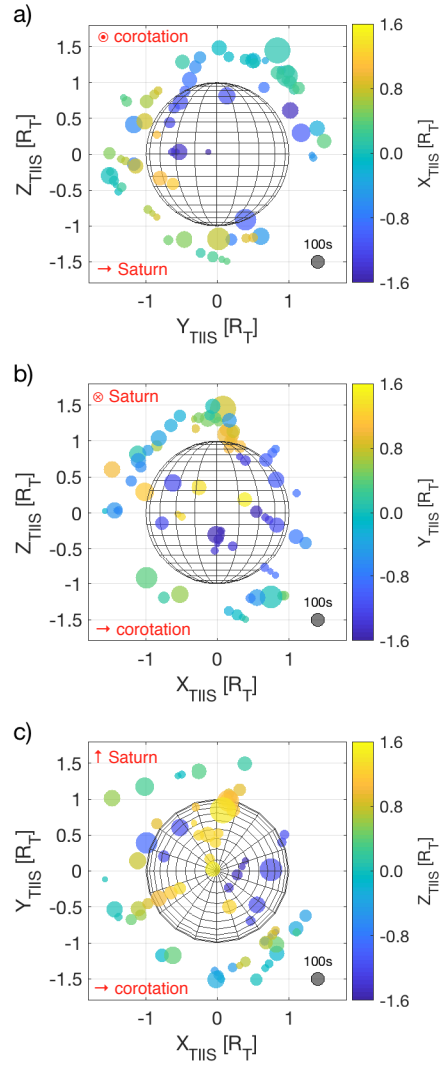


Figure 5. Figure showing the position of flux ropes at Titan in the TIIS coordinate system. The size of each point is determined by the length of time that Cassini was detecting the flux rope, where 100 seconds is shown by the key. Titan's outline is shown in black and each flux rope is a coloured circle, coloured by position in the third dimension.

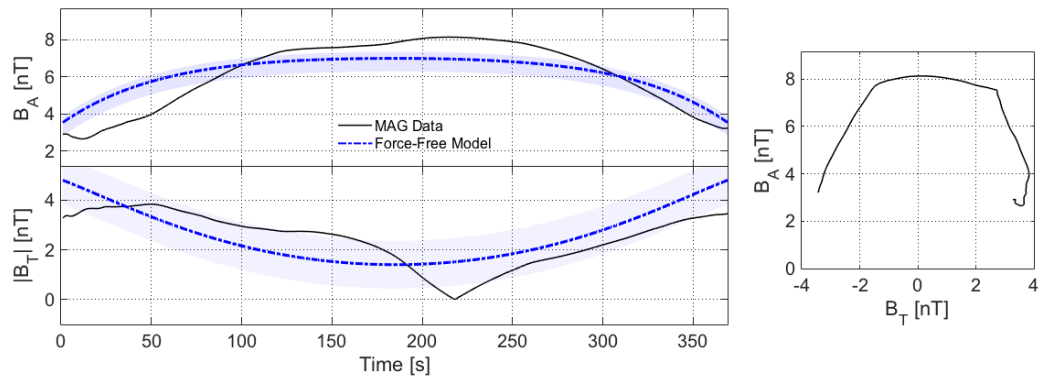


Figure 6. Figure showing an example of a flux rope signature in magnetometer data (black) fitted with a force-free model (blue) with uncertainty range (pale blue range). The figure shows axial magnetic field and tangential magnetic field. On the right is a hodogram showing the relationship between the axial and tangential field components. This flux rope occurred at 5.3 SLT on 16th April 2005 at 19:05 during flyby T5.

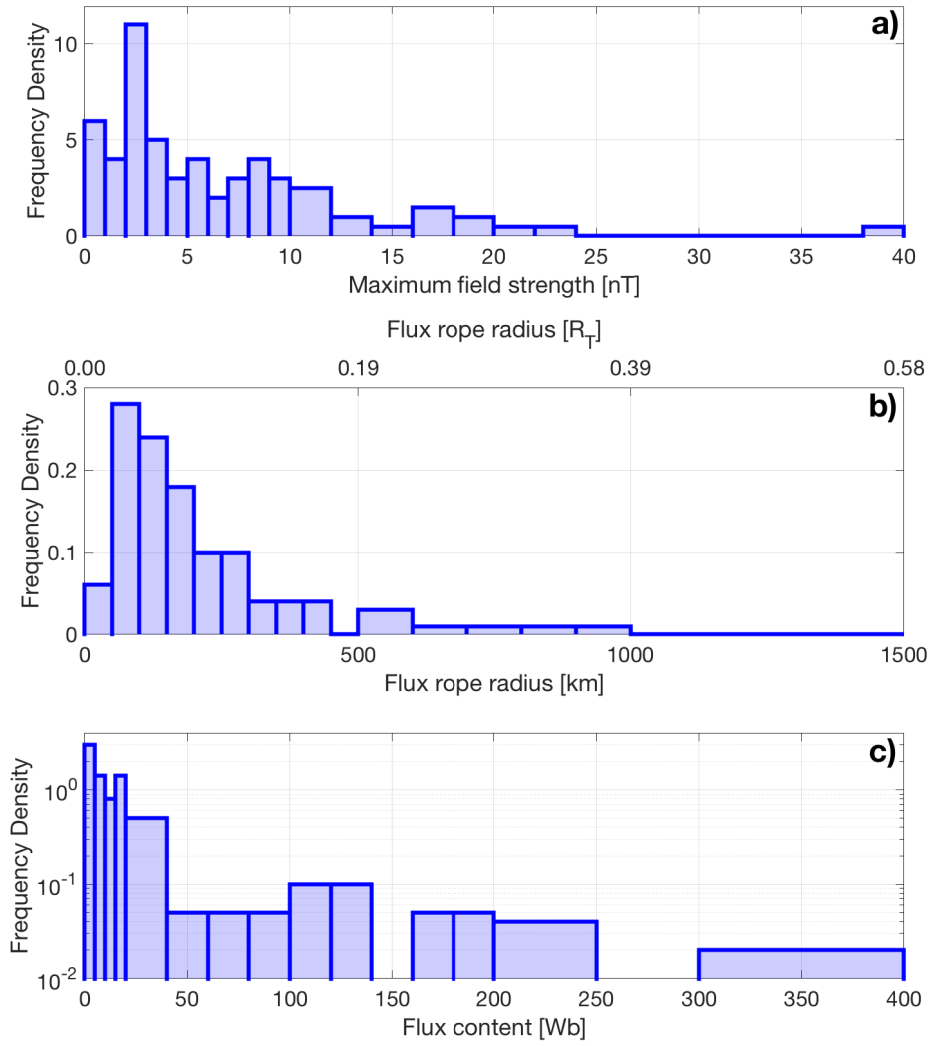


Figure 7. Figure showing a histogram of a) maximum core magnetic field, b) flux rope radius and c) flux content for 49 flux rope in Titan's ionosphere that fit to the force-free model with an MSE of 0.5 or below and $P(\chi^2)$ of 5% or less.

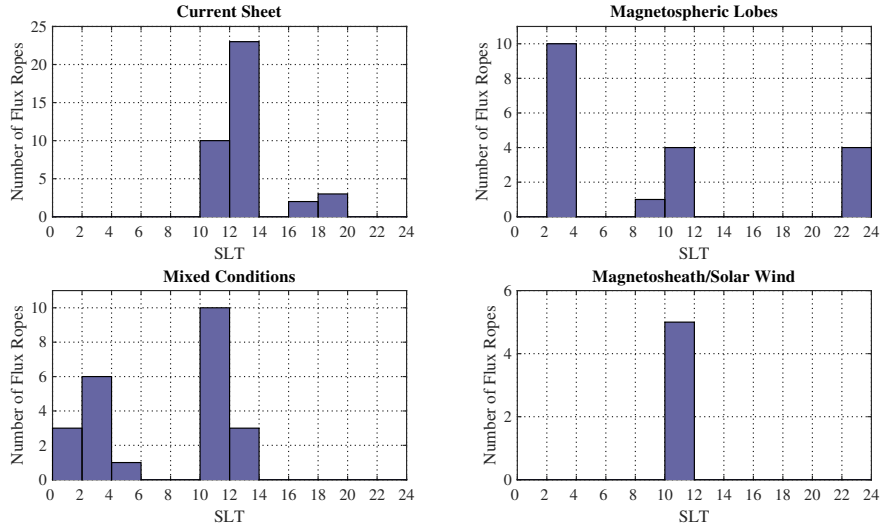


Figure 8. Histogram of the number of flux ropes with SLT ordered by magnetic environment which is determined by (Simon et al., 2010, 2013; Kabanovic et al., 2017).

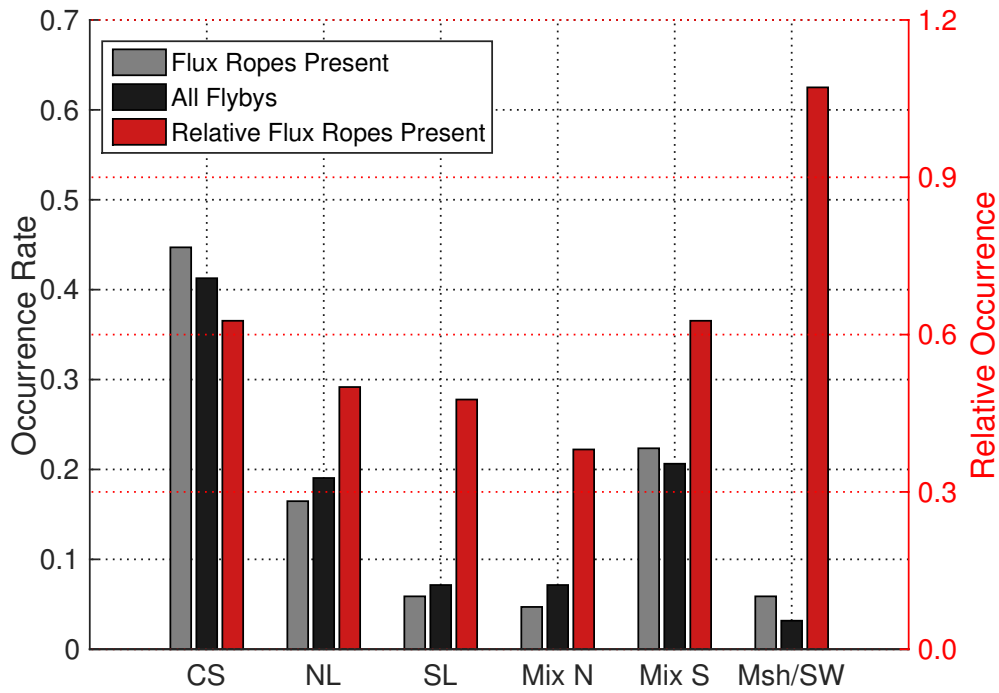


Figure 9. Histogram of the occurrence rate of different magnetospheric environments at ± 3 hours of all flybys of Titan (black, black axis), all flybys with flux ropes present (grey, black axis), and a relative occurrence of each environment where flux ropes are present relative to the occurrence rate of all flybys (red, red axis). Where CS is the current sheet, NL is the northern lobe, SL is the southern lobe, Mix N/S are a mixture with primarily northern (N) or southern (S) signatures and Msh/SW is magnetosheath or solar wind. Magnetic environment is determined by (Simon et al., 2010, 2013; Kabanovic et al., 2017).

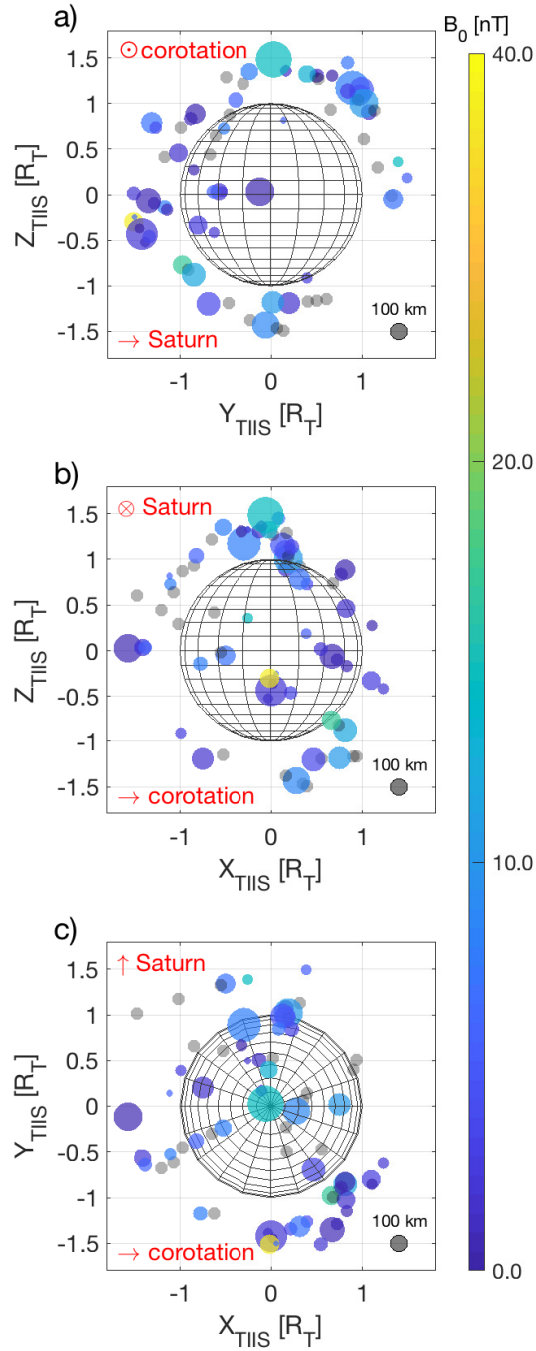


Figure 10. Figure showing the position of flux ropes at Titan in the TIIS coordinate system. The size of each point is determined by the radius found using the force-free models, where a 100 km example is shown by the key. Titan's outline is shown in black and each flux rope is a coloured circle, coloured by magnetic field strength in the centre of the flux rope. Flux ropes which were not fitted by the force-free model are shown in grey with a 100 km equivalent radius.