

Article

# Critical Review of Scintillating Crystals for Neutron Detection

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**Abstract:** There exists an ongoing need to develop and improve methods of detecting radioactive materials. Since each radioactive isotope leaves a unique mark in a form of the particles it emits, new materials capable of detecting and measuring these particles are constantly sought. Neutrons and their detectors play a significant role in areas such as nuclear power generation, nuclear decommissioning and decontamination, border security, nuclear proliferation and nuclear medicine. Owing to the complexity of their detection, as well as scarcity of <sup>3</sup>He, which has historically been the preferred choice for neutron detection in many application fields, new sensitive materials are sought. Organic and inorganic scintillating crystals have been recognised as particularly good alternatives and as such systems that utilise them are increasingly common. Since they allow investigation of the neutron energy spectra, greater information about the radioactive source can be inferred. Therefore, in this article an extensive review of scintillating crystals used for neutron detection is presented. By describing the history of scintillating crystals and discussing changes that occurred in their use and development of methods for radiation detection, the authors present a comprehensive overview of the current situation. Supported by a practical example, possible future directions of the research area are also presented.

**Keywords:** Scintillators, Scintillating crystals, Neutron detectors, Gamma detectors, <sup>3</sup>He deficit

## 1. Introduction

Radiation detection plays an important role in many application fields such as nuclear medicine, power generation, border control and nuclear decommissioning. Regardless of the application field, radiation detectors are primarily deployed to ensure safety of the personnel either working with, or in the close proximity, of the radioactive substances [1]. Further, they are essential to border and security control, where they are used to prevent illegal transportation of dangerous items [2]. Irrespective of the way they are used, a sensitive material is required that interacts directly with the targeted or expected radiation field. A large number of these devices use scintillating materials as radiation sensitive medium.

The history of the scintillating materials used for radiation detection goes back to the work by Röntgen and his famous discovery of X-rays [3]. In his experiment, Röntgen was placing barium platinocyanide plates in the close vicinity of the vacuum tubes with CaWO<sub>4</sub> powder that were previously discovered by Crookes [4]. He discovered that materials such as lead are opaque to the X-rays, whereas other materials such as aluminium are transparent. Most famously, he discovered that X-rays can be used to image bones of a human body, because calcium absorbs the X-rays owing to its relatively high atomic number, while tissues in other body parts are built of elements characterised with lower density. As such, they are more transparent to this type of radiation.

34 This discovery was embraced by a large scientific community as it allowed them to investigate  
35 previously unknown properties of materials. One of the materials investigated was crystal, as described  
36 by Friedrich et al., where they discovered X-ray diffraction within the crystal [5]. Around the same  
37 time, the structure of crystals was described based on the X-ray diffraction [6]. What became apparent  
38 as a result of these experiments was that crystals are capable of scintillating when exposed to X-rays,  
39 and thus their interactions in crystals could be observed.

40 Initially, the fluorescence produced by scintillators was observed by the naked eye, which made it  
41 difficult to conduct a suitable investigation. The requirement for a suitable photodetector resulted in  
42 the discovery of a photomultiplier tube (PMT). There exists some controversy related to the discovery  
43 of PMT, but the first electrostatic PMT (similar to the devices still produced and used today) was  
44 presented in 1936 by Zworykin et al. [7]. Nonetheless, discovery of PMTs opened up a new chapter in  
45 the history of scintillating crystals, as it made the investigation of the new materials easier and enabled  
46 new properties to be found.

47 In this article, a review of the available crystal scintillators for radiation detection, with particular  
48 focus placed on neutron detection, is presented. In the following sections, an overview of types of  
49 crystals used for radiation detection with regard to their chemical structure and particle sensitivity is  
50 presented. Further, both organic and inorganic crystals used for neutron detection are discussed in  
51 detail, as well as their growing importance given the scarcity of  $^3\text{He}$  and limitations of other detection  
52 methods. The discussion is supported through numerous examples from the literature, as well as  
53 practical example of a response of an organic crystal to mixed neutron/gamma (n/g) field provided by  
54  $^{252}\text{Cf}$ . The article is concluded with a discussion about possible future directions and expectations of  
55 where crystals may be used to further support neutron detection capabilities.

## 56 2. Scintillating Crystals used in Radiation Detection Applications

57 Regardless of the chemical type of a scintillating material, the process of extracting information  
58 from an interaction occurring within a scintillator is largely the same. When energetic particles enter  
59 the scintillator, they cause ionisation, either directly or indirectly. In the case of charged particles, e.g.  
60 protons, electrons and alpha particles, they ionise the scintillator directly. Quanta and particles without  
61 charge, such as photons and neutrons, must first transfer their energy to ionising particles within  
62 the medium. For instance, photons can liberate electrons and neutrons undergo nuclear interactions  
63 resulting in a release of charged particles (e.g.  $\alpha$ , proton). All the charged particles produced can then  
64 ionise the material raising atoms and molecules to excited states.

65 These then emit photons of visible light as they de-excite, which can be later transformed into  
66 photoelectrons through a photocathode of a photodetector such as PMT. PMTs multiply the weak  
67 signal of photoelectrons and form an electrical pulse which carries important information about the  
68 incident radiation [8]. These can be easily detected through a combination of analogue and digital  
69 electronics.

70 Characteristics of pulses observed on the outputs of a photodetector, such as their length, height,  
71 rise time, decay time, are measured and used to infer the origin of the interaction within the scintillator.  
72 These characteristics differ between scintillators and incident particles, owing to distinctive interactions  
73 that govern the scintillation process. The differences can be observed and analysed, enabling the  
74 information about the incident particles to be inferred. The most basic distinction related to crystals is  
75 between organic and inorganic crystals.

### 76 2.1. Operation Principle of Inorganic Crystals

77 One of the most frequently used crystals in radiation detection is NaI. This single crystal of  
78 alkali halide is characterised by very good spectrometric response to gamma-rays. Pure NaI crystal  
79 is an example of an insulating material. As such, its energy band structure consists of a *valence band*,  
80 which is normally full, and a *conduction band*, which is normally empty. The two are separated by  
81 *gap band*, which is also known as *forbidden gap* or *energy gap* [8,9]. When exposed to ionizing radiation,

the electrons from the valence band can be excited and move onto the conduction band. A hole in the valence band is filled when an electron returns from the conduction band. This process is accompanied by the release of a photon. However, the width of the energy gap means that the energy of the photon released is too high to be in the visible region, resulting in low light yield in the pure NaI crystal [9].

In order to alleviate this problem, impurities are introduced to inorganic crystals. These are called *activators* and are introduced to increase the likelihood of emitting photons that can be detected through conventional photodetectors. When an electron is returning to the valence band, in an insulating material such as pure NaI, a photon may be emitted. However, due to the width of the energy band, it may be self-absorbed. Therefore, the energy band structure of the crystal matrix is changed when an activator is added. The activator introduces states within the energy gap of the pure crystal matrix. Thus photons, which can be easily detected through conventional methods, can be emitted.

One of the most common activators is Tl. As an example, this activator alters the maximum emission wavelength from 303 nm in pure NaI to 450 nm in thallium doped NaI crystal, and notation NaI(Tl) is used [8]. Generally, activators create new regions within the crystalline structure of a scintillator, which are sometimes referred to as *luminescence centres* or *emission centres*. These enable the scintillators emitted wavelengths to be more closely matched with the sensitivity regions of the PMTs.

Depending on the application different properties of the inorganic crystals may be sought. However, there exists a basic set of requirements that is desirable across many application fields which includes a fast response, high light yield, high density and high atomic number [10]. Excellent gamma-ray sensitivity and energy resolution should naturally lie above the mentioned characteristics. A material meeting all of these criteria does not exist. For instance, NaI(Tl) and CsI(Tl) are characterised by the high light yield, but relatively slow response time. In contrast, pure CsI crystal exhibits very fast response but low light yield in the room temperature range. One of the inorganic crystals that was utilised in varied application areas due to its unique combination of the specified characteristics is Lu<sub>2</sub>SiO<sub>5</sub>(Ce) (LSO) [11]. As such, it was successfully exploited, together with its modified version containing yttrium - e.g. Lu<sub>1.8</sub>Y<sub>0.2</sub>SiO<sub>5</sub>(Ce) (LYSO) in e.g. nuclear medicine for Positron Emission Tomography (PET) applications.

Inorganic crystals were primarily developed for application in gamma-ray detection and characterisation applications, due to their suitability in areas requiring excellent energy resolution. However, there have been numerous inorganic crystals developed, which are directly aimed at low-energy neutron detection. This is possible because the crystals contain a high neutron cross-section material such as Li [9].

## 2.2. Inorganic Crystals Capable of Neutron Detection

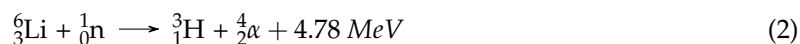
Owing to their high cross-section for low-energy neutron capture, the most commonly used isotopes are <sup>10</sup>B, <sup>6</sup>Li and <sup>3</sup>He. The most common nuclear reaction with <sup>3</sup>He used for neutron detection is defined in Eq. 1. It is accompanied by the release of 0.764 MeV of kinetic energy, and cross-section for this particular reaction is 5330 barns, for thermal neutrons [9]. Fast neutron detectors based on <sup>3</sup>He have also been implemented, where appropriate moderating material is added to thermalize the fast neutrons [12]. However, scarcity of <sup>3</sup>He, caused by the decline in tritium production for nuclear weapons maintenance, requires that other alternatives be sought [13].



One of the proposed alternatives are organic scintillation detectors utilising *elastic scattering* of neutrons with light atoms, such as hydrogen [14,15]. When considered as an alternative for <sup>3</sup>He detectors, organic scintillation detectors exhibit gamma-ray sensitivity which requires particles to be separated. However, detection systems exploiting both scattering and particle separation techniques (will be discussed in the following section) have shown a promising performance with regard to source localisation, as well as particle identification [16].

### 128 2.3. Detectors Utilising ${}^6\text{Li}$ Neutron Reaction

129 Out of the remaining two isotopes,  ${}^6\text{Li}$  has been most widely adapted in inorganic crystals. One  
130 of the examples of a scintillating crystal capable of neutron detection, which contains Li, is another  
131 alkali halide - LiI(Eu). Detectors containing Li represent a group of potential candidates for detection  
132 of low-energy neutrons owing to the  ${}^6\text{Li}(n,\alpha)$  reaction, as defined in Eq. 2.



133 When a scintillator is sensitive to both neutrons and gamma-ray photons, it is necessary to separate  
134 the two particle types. This phenomenon is often referred to as pulse shape discrimination (PSD)  
135 and is very common in the domain of organic scintillators. The  $\alpha$  particle resulting from neutron's  
136 interaction with  ${}^6\text{Li}$  can be detected and easily classified through the PSD methods [17].

137 Relatively recent study investigating the doping of the pure LiI crystal with  $\text{Eu}^{2+}$  show that  
138 appropriate doping level, as well as heat treatment may hold an answer to the light yield problem,  
139 when used for neutron detection. It should be noted that the heat treated LiI: $\text{Eu}^{2+}$  scintillator examined  
140 by Boatner et al. [18] also shows excellent spectral response to gamma-rays from  ${}^{137}\text{Cs}$  calibration  
141 source.

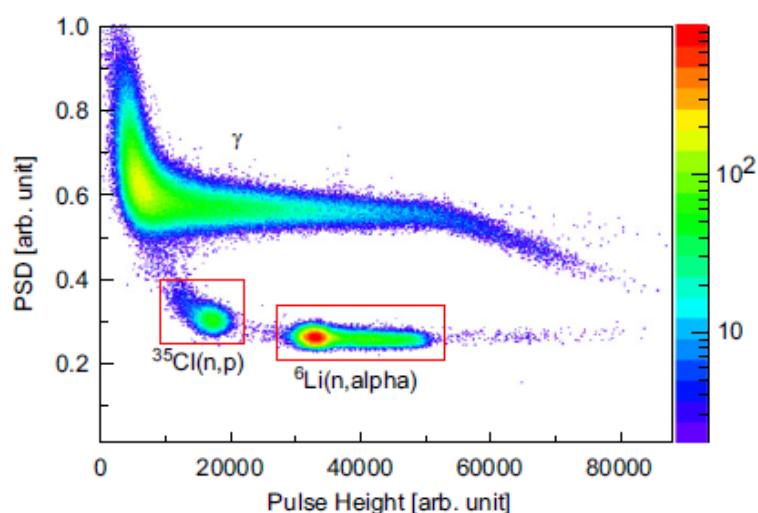
142 Another detector utilising the high thermal neutron cross-section of  ${}^6\text{Li}$  isotope is  $\text{Ce}^{3+}$  doped  
143 LiCaAlF<sub>6</sub> inorganic crystal. When experimentally tested, this detector's performance was compared to  
144 that of a commercially available Li-glass scintillator [19]. Samples of two different sizes of LiCaAlF<sub>6</sub>  
145 were manufactured, and tested in regard to the light yield, n/g separation capabilities and neutron  
146 detection efficiency. Regardless of the sample size the light yield was considerably lower than measured  
147 for Li-glass detector. However, n/g separation capabilities were deemed as high, and the intrinsic  
148 neutron detection efficiency (for the large size sample - 50.8 mm × 2 mm) was estimated to 80% of the  
149 Li-glass counterpart.

150 PSD methods have also been applied to successfully separate neutrons from gamma-ray photons  
151 in crystals such as LiAlO<sub>2</sub> and LiGaO<sub>2</sub> [20]. In this case, Cherenkov radiation can be used to distinguish  
152 between neutrons and gammas, as it provides a cut-off point between the fast and slow component in  
153 the pulse decay. As tested with  ${}^{252}\text{Cf}$ , the researchers show that scintillators are capable of detecting  
154 fast neutrons. It is believed that detector's sensitivity could potentially be extended to thermal energy  
155 region.

156 A very good potential for neutron detection via PSD methods is presented by detectors utilising  
157 LiBaF<sub>3</sub> crystal doped with Ce. The discrimination between various particles, across broad energy  
158 spectrum, is possible due to the occurrence of core-valence luminescence (CVL). It is a very short pulse  
159 (sub-nanoseconds) resulting from a hole in the conduction band of an ionic crystal that is being filled  
160 by an electron travelling from the valence band [10]. It appears alongside the self-trapped-exciton  
161 (STE) luminescence, when the crystal is exposed to gamma-ray field. When it is exposed to the neutron  
162 field, only the STE luminescence is observed. It is reported to have a very decent energy resolution, as  
163 well as being able to discriminate between gammas, thermal and fast neutrons [21].

164 Another group of crystals capable of neutron detection are elpasolites, which include scintillators  
165 such as Cs<sub>2</sub>LiYCl<sub>6</sub> (CLYC) and Cs<sub>2</sub>LiLa(Br,Cl)<sub>6</sub> (CLLBC). When doped with Ce, these crystals present  
166 excellent n/g separation characteristics, as well as very high energy resolution [22]. An example of  
167 PSD capabilities of CLYC scintillator is presented in Fig. 1. Fast neutron detection can also be facilitated  
168 by growing the crystals using  ${}^7\text{Li}$ , rather than the traditionally used  ${}^6\text{Li}$  to maximise thermal neutron  
169 sensitivity. Moreover, a number of composite detectors has been developed, consisting of CLYC crystal  
170 incorporated into an organic plastic, to further extend the sensitive spectrum to fast neutrons [23,24].

171 Further example of an inorganic scintillator for neutron detection that is popularly used is  
172  ${}^6\text{LiF}/\text{ZnS:Ag}$  [26]. At the heart of this scintillator lies ZnS crystalline powder, which was famously  
173 used by Rutherford in his work on the stability of atoms [27]. ZnS:Ag powder is characterised by a  
174 very good light yield of 75000 photons/MeV and relatively slow decay time of 1.4  $\mu\text{s}$  [28]. In the same



**Figure 1.** An example of PSD capabilities of CLYC scintillator when exposed to a specific neutron field of 1.3 MeV produced by a generator. Reproduced from [25].

175 study, the author attempts to characterise pure ZnS single crystal. The analysis presented suggests  
 176 that due to the absence of the Ag dopant, the light yield is reduced significantly. It is therefore clear  
 177 that a scintillator in this form would not be capable of detecting neutrons. However, when  ${}^6\text{LiF}$  is  
 178 added to the mix it becomes an efficient thermal neutron detector with low gamma-ray sensitivity. It is  
 179 commercially available from Eljen Technology as EJ-426 [29].

#### 180 2.4. Detectors Utilising Other Properties of Inorganic Crystals

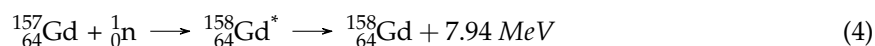
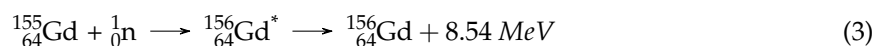
181 The ongoing research into finding an appropriate alternative for  ${}^3\text{He}$  detector has resulted in  
 182 new ways of using well established inorganic crystals. One of such examples is  $\text{YAlO}_3:\text{Ce}^{3+}$  which  
 183 was successfully used for gamma radiation detection. Neutron sensitivity was in this case facilitated  
 184 by adding *converter* in a form of a powder to the surface of the scintillator. Depending on the energy  
 185 group of neutrons targeted possible candidates are lithium, boron, gadolinium (thermal neutrons) and  
 186 thorium, hydrogen (fast neutrons).

187 The discrimination between gamma-ray and neutron interactions is performed via pulse height  
 188 discrimination (PHD) and has been successfully presented with PuBe source [30]. A detector utilising  
 189  $\text{YAlO}_3:\text{Ce}^{3+}$  with neutron *converter* would benefit from the intrinsic properties of the perovskite detector  
 190 such as fast decay time, high light yield and good stopping power. Simultaneously, the size of the  
 191 detector could be kept small which is often desired in applications such as nuclear medicine. However,  
 192 as with all inorganic scintillation crystals it is characterised by very high gamma-ray sensitivity which  
 193 makes the analysis and discrimination process difficult.

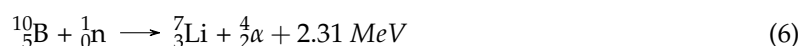
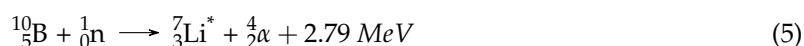
194 One of the materials mentioned in the preceding paragraph (gadolinium) is characterised by the  
 195 highest thermal neutron cross-section known. Apart from being used as a *converter*, gadolinium based  
 196 detectors form another group of good fast neutron detecting crystals.  $\text{Gd}_3\text{Al}_2\text{Ga}_3\text{O}_{12}:\text{Ce}$  (GAGG:Ce)  
 197 crystal is characterised by excellent light yield and good stopping power. Neutron interactions with  
 198 gadolinium are primarily driven by  ${}^{155}\text{Gd}(n,\gamma)$  and  ${}^{157}\text{Gd}(n,\gamma)$  reactions, for which the cross-sections  
 199 are 60900 and 255000 barns, respectively. The reactions are defined in Eq. 3 and Eq. 4, where the  
 200 unstable products return to the ground state with a release of gamma-rays.

201 The resulting neutron and gamma-ray induced pulses must be separated via appropriate method.  
 202 However, there is no need for material enrichment due to exceptional neutron sensitivity of gadolinium.  
 203 Moreover, it is possible to retrieve incident kinetic energy of a neutron interacting within the crystal  
 204 which opens up the possibility of performing neutron spectroscopy. Recent study performed with

205 AmBe source showed a superior performance of this crystal, when compared with an established  
 206  $^6\text{Li}$ -glass detector [31]. Given the fast response of the crystal to gamma-ray photons, it is also feasible  
 207 to explore time-of-flight based discrimination. Therefore, it comes at no surprise that a lot of research  
 208 effort is currently going into the improvement of this detector. However, as with most of inorganic  
 209 crystals high cost, and long growing time may be unacceptable in many applications.



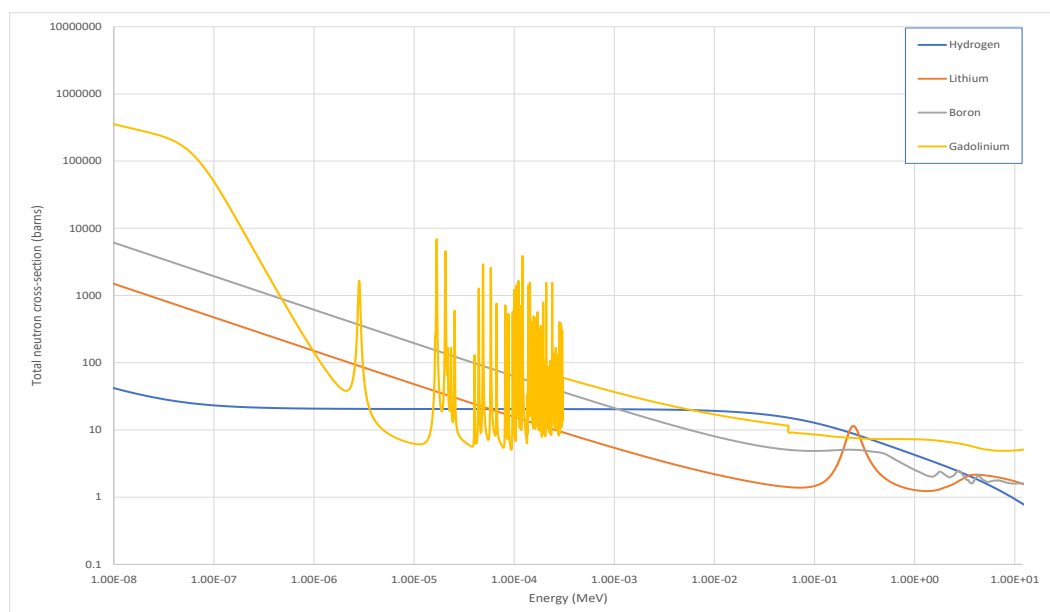
210 Detection of thermal neutrons using  $^{10}\text{B}$  reactions is well established in the domain of organic  
 211 scintillators [32]. Doping with  $^{10}\text{B}$  enables the sensitivity spectrum of organic scintillators, which is a  
 212 very good fast neutron detector, to be extended to the thermal region.  $^{10}\text{B}(n,\alpha)$  reactions, as defined in  
 213 Eq. 5 and Eq. 6, are probably most widely used mechanism for detection of thermal neutrons, owing to  
 214 high thermal neutron cross-section (3840 barns) [9]. The reaction can lead to a stable or an unstable  $^7\text{Li}$   
 215 isotope and is accompanied by the release of  $\alpha$  particle that can be easily detected using conventional  
 216 methods.



217 Although popular in the domain of organic scintillators, there are not many examples of  
 218 inorganic crystals utilising  $^{10}\text{B}$  based reactions. However,  $\text{Li}_6\text{Y}(\text{BO}_3)_3:\text{Ce}$  has been computationally  
 219 and experimentally tested showing good potential for thermal neutron detection. It is reported to be a  
 220 relatively fast scintillator with a decay time for thermal neutrons of  $38 \pm 18 \text{ ns}$ , and to show a greater  
 221 thermal neutron detection efficiency than Li-glass scintillator. However, its light yield is estimated to  
 222 be six times lower than NaI:Tl, and  $\alpha/\gamma$  ratio is ten times lower than that of Li-glass. The  $\alpha/\gamma$  ratio  
 223 is a measure used to assess scintillator's ability to separate  $\alpha$  and  $\gamma$  interactions. The assessment is  
 224 based on the pulse height information. Generally, the light yield produced as a result of  $\alpha$  interactions  
 225 is lower than that resulting from  $\gamma$  interactions for the same amount of energy deposited [33]. Another  
 226 potential area of application for boron doped crystals capable of neutron detection is considered to  
 227 be space instrumentation, with initial experiments showing reasonable results in regard to thermal  
 228 neutron detection efficiency [34].

229 Total neutron cross-section for the discussed elements is presented in Fig. 2. It can be observed  
 230 that gadolinium (shown in yellow) has the highest overall cross-section for the low energy regions. In  
 231 agreement with the quoted barn values lithium (shown in orange) has the lowest cross-section out of  
 232 the three considered candidates. However, there is a noticeable spike between 100 keV and 1 MeV  
 233 that could be exploited in a specific application targeting this energy region. Boron (shown in grey)  
 234 appears to be the most stable, out of the three thermal detector options, across the energy spectrum.  
 235 For comparison, hydrogen's cross-section (shown in blue) is considerably lower than the other three  
 236 elements in the thermal energy region. Therefore, organic scintillators are primarily used to detect fast  
 237 neutrons, due to their high hydrogen content.

238 It is also worth noting that as early as 1968, it was attempted to perform neutron detection  
 239 using NaI(Tl) crystal [35]. The experiment was performed with  $^{127}\text{I}$  to observe crystal's response to  
 240 low energy neutrons (via radiative capture) and fast neutrons (inelastic scattering). When tested in  
 241 monoenergetic field of 1 MeV neutrons, overall efficiency was measured as 0.5 %, considerably lower  
 242 than that obtained for organic scintillators. As a result, research into suitable fast neutron detection  
 243 was pursued within the organic scintillators' domain.



**Figure 2.** Total neutron cross-sections for the discussed elements: hydrogen, lithium, boron and gadolinium. The cross-sections of the selected isotopes were generated using ENDF/B-VIII.0 libraries.

244 Heavy oxide scintillator crystals represent another group of detectors showing potential of neutron  
245 detection. Most commonly used examples of this group are  $\text{CdWO}_4$  and  $\text{PbWO}_4$  crystals.  $\text{CdWO}_4$  is  
246 capable of providing a very good spectral response to fast neutrons, but there exist handling issues in  
247 some places (e.g. UK) related to this crystal due to toxicity of Cd [36]. Similarly, has been tested for  
248 its fast neutron sensitivity [37]. Despite relatively good response in comparison to other counterparts  
249 tested, its low light yield makes it unsuitable for many applications [22].

### 250 2.5. Organic Crystals Operation

251 Regardless of their state (solid or liquid), organic scintillators are generally sensitive to both fast  
252 neutrons and gamma-ray photons. Therefore, many PSD methods have been investigated to facilitate  
253 low misclassification probability. The difference between the two particles can be inferred from the  
254 varying rate of energy loss of the particle, when scattered in the scintillation medium. Fast neutrons  
255 primarily undergo *elastic scattering* with a proton, while gamma-ray photons interact with the atoms of  
256 the scintillant via *Compton scattering*. These result in *fluorescence*, whose decay time is proportional to  
257 the rate of energy loss of the incident particle. Appropriate photodetector is then capable of detecting  
258 the *fluorescence*, and gives rise to a proportional electronic pulse. The rate of energy loss is greater for  
259 Compton electrons (resulting from gamma-ray interactions), when compared to protons (resulting  
260 from neutron interactions). This difference is reflected in the tail of the electronic pulse produced by  
261 the detector [9].

262 There are only two pure organic crystals that have been widely exploited in radiation detection  
263 applications: anthracene and stilbene. Anthracene was popularly used due to its scintillation efficiency,  
264 which is the greatest of all organic scintillators [9]. Scintillation efficiency of organic scintillators is often  
265 quoted as a percentage of anthracene's light output. Stilbene on the other hand, was characterised by  
266 an excellent n/g separation capabilities and was originally used by Brooks [38] when investigating PSD  
267 methods in the analogue domain. However, due to the issues related to growing of these crystals in

268 greater dimensions, they have been left aside for many years. In the first decade of the 21st century, an  
269 interest has grown back due to new growing methods developed by the team at Lawrence Livermore  
270 National Laboratory (LLNL) in the US led by Natalia Zaitseva [39].

271 Given its excellent light yield anthracene still remains as the material that is characterised by the  
272 best scintillation efficiency available and is often used as a reference when developing new crystals.  
273 It was also tested for its PSD capabilities and even though inferior to stilbene decent separation  
274 was observed [39]. One of the disadvantages of using organic crystals is their anisotropic response  
275 to incident radiation, which affects the performance when the orientation of the detector changes.  
276 However, this property can also be exploited to infer the location of the interaction via the angle of the  
277 scattered proton. It was successfully used by Brubaker and Steele [40] to perform neutron imaging.

278 Traditionally, *trans-stilbene* crystals were grown using the melt growth method. Growth process  
279 was associated with both high complexity of the growth process and high cost. Hence, they were  
280 only grown in sizes not exceeding 10 cm. However, when new solution growth method was applied,  
281 the growth time was reduced, and samples of greater sizes were grown. It also partially addresses  
282 the well-recognised issue of high misclassification between neutrons and gamma-ray photons in  
283 the low energy region. Furthermore, when tested in regard to its light yield and PSD capabilities,  
284 solution grown stilbene crystal performed considerably better than equivalent melt grown stilbene  
285 and organic liquid scintillator - EJ-309 [41]. It also shows better PSD characteristics than other PSD  
286 plastic scintillators [42].

287 As a solid, non-hygroscopic, not hazardous material, light-weight stilbene crystal is suitable for  
288 many applications such as nuclear decommissioning and portable security devices [43]. Although it is  
289 now possible to grow these crystals in larger sizes, the cost of manufacturing is still relatively high  
290 suggesting that organic liquids may still be more cost effective for large scale detectors. Nevertheless,  
291 the continuous interest in the field of organic crystals has led to the development of a new stilbene  
292 crystal, where hydrogen is replaced with deuterium. This deuterated stilbene is reported to have  
293 even better PSD capabilities than the standard stilbene [44]. Another organic crystal that should be  
294 mentioned at this stage is rubrene crystal, that is also grown from solution and is reported to show  
295 clear response to  $\alpha$  particles, and a moderate response to fast neutrons [45].

296 Based on the presentation of the scintillating crystals currently utilised in neutron detection  
297 applications, it can be noticed that there is no single choice that would account for all the requirements  
298 of a neutron detector. Therefore, it is essential to carefully analyse the requirements of a detector  
299 and choose the sensitive material accordingly. In the following section, a practical example of an  
300 organic stilbene crystal tested in the mixed field of  $^{252}\text{Cf}$ , in regard to its pulse shape discrimination  
301 capabilities is presented. This particular scintillator was chosen, as it illustrates the feature of lower  
302 misclassification probability at lower neutron energies. Results obtained are then analysed, and the  
303 article is concluded with the future outlook for scintillating crystals in neutron detection field.

## 304 2.6. Summary

305 There exist a vast number of scientific resources available, where the most important properties  
306 of scintillating materials have been documented. However, these are generally focusing on specific  
307 particles (e.g. gamma-ray detectors) or subset of the particle group (e.g. thermal neutrons). In this  
308 work, an attempt was made to present the properties of the most promising candidates that have  
309 been examined in respect to neutron detection potential. In Table 1, a comparison of the selected  
310 inorganic and organic scintillating crystals is shown. A broad range of materials is covered, including  
311 both inorganic and organic crystals, capable of gamma-ray detection as well as n/g detection. For  
312 comparison, typical liquid and plastic scintillators are also included.

313 Data in Table 1 presents a list of potential candidates for the specific applications with regard to  
314 the target particle types. There are two particular materials that bring the distinct advantages to n/g  
315 detection and are aimed at different areas of the neutron energy spectrum. In the region of thermal



**Table 1.** Comparison of the most prominent properties of scintillating crystals capable of gamma-ray and n/g detection. Data presented below was compiled based on the following references [9,20,46–58].

Scintillation material	Density (gm/cm <sup>3</sup> )	Wavelength (nm)	Refractive index	Decay time (ns)	Light yield (Photons/MeV)		Energy resolution (% at 662 keV)
					Neutron	Gamma	
NaI(Tl)	3.67	415	1.85	230	-	41,000	5.6
CsI(Tl)	4.51	550	1.8	800	-	66,000	6.6
CsI(Na)	4.51	420	1.84	630	-	40,000	7.4
LSO(Ce)	7.4	420	1.82	40	-	26,000	7.9
LYSO(Ce)	7.2	400	1.81	30-35	-	32,000	8.5
LiI(Eu)	4.1	470	1.96	1400	50,000	12,000	8
LiCaAlF <sub>6</sub> (Eu)	2.94	370	1.4	40	30,000	29,000	-
LiCaAlF <sub>6</sub> (Ce)	2.94	300	1.4	40	4,000	1,600	-
LiAlO <sub>2</sub>	2.61	330	-	790(5400 not enriched) <sup>6</sup> Li)	5,900	7,000	-
LiGaO <sub>2</sub>	4.18	330	-	12(680 not enriched) <sup>6</sup> Li)	5,500	5,000	-
CLYC	3.3	380	1.81	50; 1,000	70,000	20,000	4
CLLBC	4.1	410	1.9	55; <270	180,000	60,000	3.5
<sup>6</sup> LiF/ZnS:Ag	2.6	450	-	80,000(neutron),100(gamma)	160,000	75,000	-
YAlO <sub>3</sub> :Ce <sup>3+</sup>	5.37	370	1.95	30	-	21,000	4.3
GAGG:Ce	6.63	520	1.9	100	-	56,000	-
Li <sub>6</sub> (BO <sub>3</sub> ) <sub>3</sub> :Ce	2.8	420	-	27	-	1,200	-
CdWO <sub>4</sub>	7.9	495	-	5000	-	20,000	6.8
PbWO <sub>4</sub>	8.28	420	2.16	6,30	-	205	-
Stilbene	1.25	390	1.626	3.5 - 4.5	10,700	14,000	-
Anthracene	1.16	447	1.62	30	20,000	20,000	-
EJ-309	0.96	424	1.57	3.5(short component)	12,300	12,300	-
EJ-276	1.096	425	-	g(13, 35, 270);n(13, 59, 460)	8,600	8,600	-

neutrons, CLYC appears to be a very promising candidate, as it presents a very decent results across the considered properties, and its PSD capabilities are exceptional, as presented in Fig. 1.

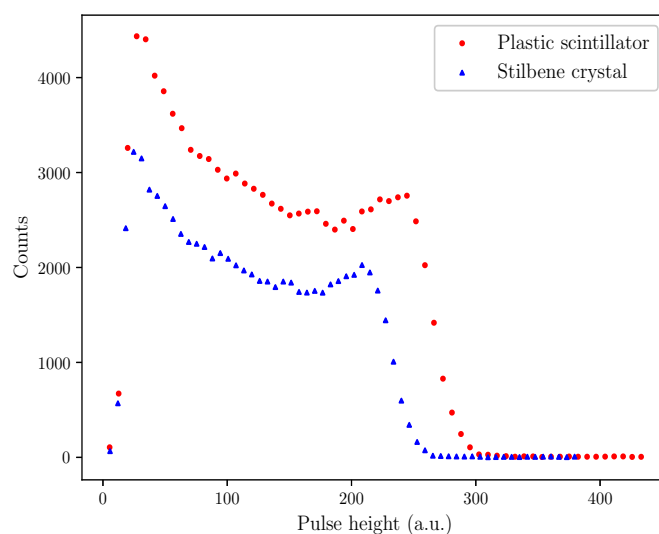
Fast neutron detection is primarily targeted by organic scintillation materials. These are presented in the last four rows of Table 1. It can be noticed that continuous development of the new crystal growing methods results in improved light yield of stilbene crystal which used to only achieve approx. 50% of anthracene's yield [9]. Similarly to CLYC for thermal neutrons, stilbene's PSD performance is superior to other organic scintillators in the region of fast neutrons. A comparison of stilbene's PSD performance to that of plastic scintillator is shown in further section of this article.

### 3. Methodology

This section describes the methodology of the work performed in order to present the PSD potential of single stilbene crystal through a comparison of its performance with that of an organic plastic scintillator. Firstly, the energy calibration process is described. It is followed by the description of a PSD technique used in this experiment and concluded with the explanation of the PSD quality assessment method used in this study.

#### 3.1. Energy Calibration

Prior to the experiments performed within the mixed-field environment of <sup>252</sup>Cf both scintillators were calibrated using <sup>137</sup>Cs gamma-ray source of 319 kBq current activity. Each detector assembly was in turn exposed to the gamma-ray field of <sup>137</sup>Cs by placing the detector assembly 15 cm away from the point source. Each detected pulse was processed through a bespoke pile-up rejection algorithm where a pulse was rejected if two peaks within one trigger window were detected. Also baseline subtraction was performed by calculating the average over the periods before and after the pulse within the trigger window. Given that the pulse was detected between sample no. 50 and 100, baseline was calculated over samples 1-45 and 105 - 128. There were 104,069 pulses accepted for the plastic scintillator sample and 74,684 pulses for the single stilbene crystal. These were subsequently used to plot the pulse height spectra, as presented in Fig.3 and adjust the equivalent energy scale for PSD considerations.



**Figure 3.** Pulse height spectra of each scintillator obtained with  $^{137}\text{Cs}$  used to perform energy calibration of the detectors.

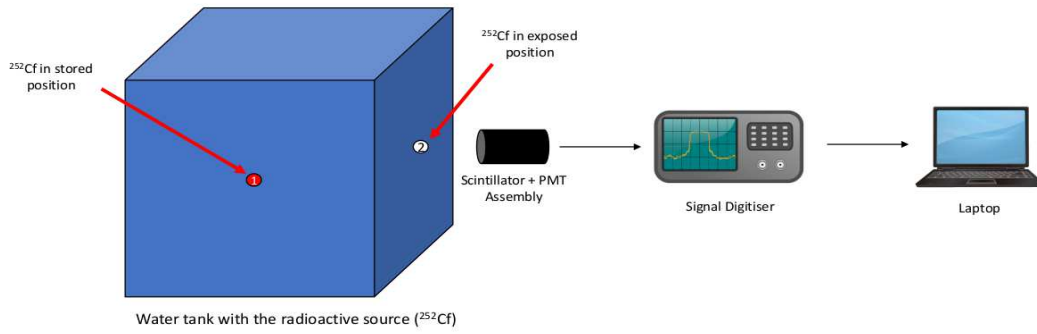
### 341 3.2. Pulse Shape Discrimination

342 In order to illustrate the capabilities of organic scintillators in regard to fast neutron detection,  
343 two solid state organic detectors have been tested in the mixed-field ( $n/\gamma$ ) environment provided  
344 by  $^{252}\text{Cf}$  at Lancaster University, UK. A single stilbene organic crystal scintillator was obtained from  
345 Inrad Optics in 2016. PSD performance of this cylindrical crystal (20 cm  $\times$  20 cm) was compared with  
346 that of an organic plastic cylindrical sample (25.4 cm  $\times$  25.4 cm) obtained from Lawrence Livermore  
347 National Laboratory (LLNL) in the US, with LLNL sample number 5706. Samples have been covered  
348 with reflective coating on the side and back to minimise the chance of photons escaping the scintillator  
349 without being detected. Each scintillator was then in turn attached to a single channel ET Enterprises  
350 9107B PMT using EJ-550 silicon grease. The PMT anode signals were collected via FPGA based signal  
351 digitiser operating at the sampling frequency of 500 MS/s with 12-bit resolution.

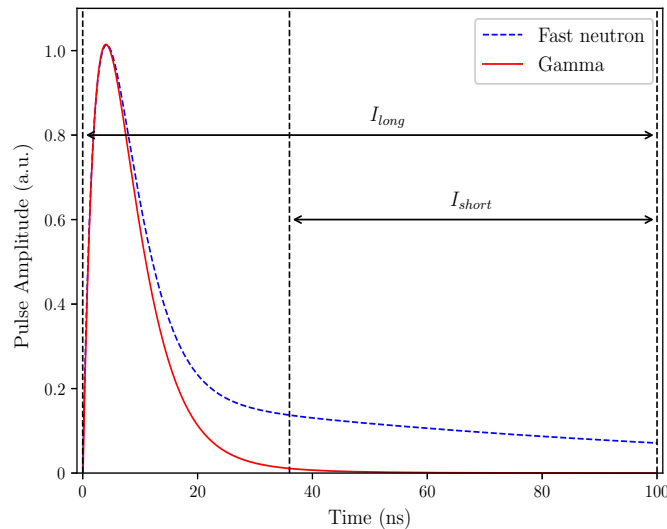
352 The complete assembly, comprising scintillator and the PMT, was placed in a cylindrical  
353 light-proof box and placed in front of the water tank, where the radioactive isotope is normally  
354 stored. The radioactive source is normally located in the centre of a water-filled tank, as shown in Fig.  
355 4. For experiments the source is pneumatically moved to the edge of the tank, which stops approx.  
356 20 cm away from the edge. The detector assembly was placed 15 cm away from the edge of the tank,  
357 resulting in the total distance of 35 cm between the source and the detector front. Each scintillator was  
358 exposed for the duration of 1 hour. The FPGA based digitiser collected raw data, with each sample  
359 collected every 2 ns. Detection window consisted of 128 samples, collected over 256 ns trigger period.

360 Before any further analysis was performed, quality of each pulse detected was assessed through  
361 the pile-up rejection algorithm in the same way as for the energy calibration. Similarly, the baseline  
362 removal was performed. Charge Comparison Method (CCM) was applied in the digital domain to  
363 assess  $n/g$  separation capabilities of the scintillator samples.

364 The CCM is the most popularly used method, where the pulse is analysed by calculating integrals  
365 over two different time intervals [38]. As the difference between the neutron and gamma-ray induced  
366 interactions is most prominent in the tail of the pulse, the *short integral* is calculated between a point  
367 some time after the peak of the pulse and the end of the pulse, as specified in Fig. 5. The *long integral*  
368 is calculated over the entire duration of the pulse. These can then be used to calculate the discrimination  
369 factor, as described below, and generate a plot exploiting the PSD capabilities of the detector.



**Figure 4.** Diagram presenting the experimental set-up, with the radioactive isotope in the centre of a water-filled steel tank (position 1), where it is normally stored. For experiments the source is pneumatically moved to the edge of the tank (position 2).



**Figure 5.** Illustration of the implementation of the pulse shape discrimination method used in this study. Long and short integrals used in CCM calculations are clearly marked on the plot. Theoretical fast neutron and gamma-ray pulses were obtained based on the data from Knoll [9] and Zaitseva et al [59].

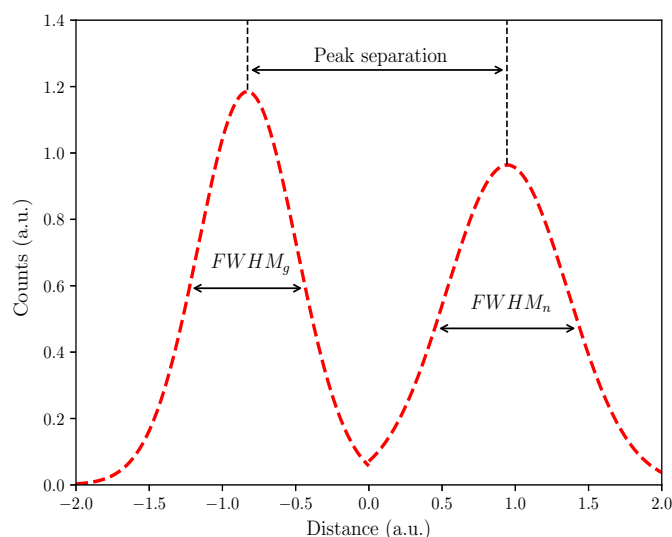
370 There are numerous ways of presenting the implementation results of CCM. One of the most  
 371 reliable methods is to calculate a discrimination factor and present it with respect to the electron  
 372 equivalent energy for each detected interaction. In this work, the discrimination factor  $D_f$  was  
 373 calculated using the equation presented in Eq. 7. The remaining terms in Eq. 7 ( $I_{short}$ ,  $I_{long}$ ) correspond  
 374 to the integrals introduced in 5. The discrimination factor was then plotted against the equivalent  
 375 energy of the pulse, following the calibration process described in the preceding subsection.

$$D_f = 1 - \frac{I_{short}}{I_{long}} \quad (7)$$

### 376 3.3. PSD Quality Assessment

377 The concept of FOM as a measure for particle separation quality was originally introduced by  
 378 Winyard et al. [60]. In order to estimate the FOM, the data needs to be presented in a form of a

379 plot, where the distribution of the particles is illustrated. For neutrons and gamma-ray photons it is  
 380 expected that they will show normal distribution spread. An example n/g distribution is presented in  
 381 Fig. 6. Terms identified in Fig. 6 are then used to calculate the FOM, as presented in Eq. 8.



**Figure 6.** Example neutron and gamma-ray distributions based on the distance to the discrimination line.

$$\text{FOM} = \frac{\text{Peak separation}}{\text{FWHM}_g + \text{FWHM}_n} \quad (8)$$

## 382 4. Results

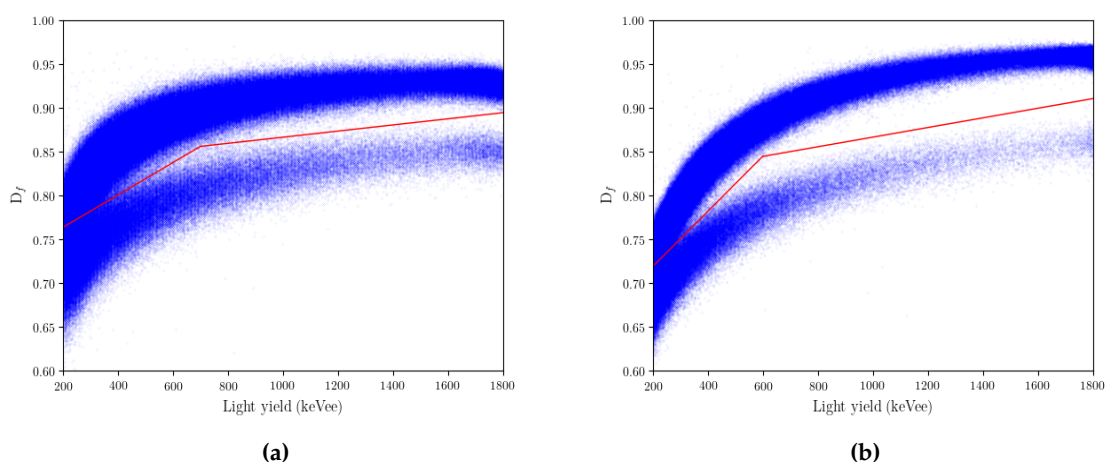
383 Each scintillator was in turn exposed to the mixed-field environment provided by  $^{252}\text{Cf}$  for the  
 384 duration of 60 min. There were 902,564 pulses accepted for the plastic scintillator sample, and 840,583  
 385 pulses for the organic crystal sample. PSD scatter plots for each sample are presented in Fig. 7a (plastic)  
 386 and Fig. 7b (crystal). Discrimination factor  $D_f$ , as defined in previous sections, has been plotted against  
 387 the electron equivalent energy. The resulting plumes represent the neutron and gamma-ray photon  
 388 interactions, with gamma-rays depicted by the upper plume and neutrons by the lower plume.

389 Following that, PSD separation quality was assessed for each scintillator using FOM. Given the  
 390 way data are presented in this study, a discrimination line was plotted to mark the visible separation  
 391 between the plumes. The distance from each point to the discrimination line was then plotted in form  
 392 of a histogram in order to show the distribution of the considered particles. This method was used to  
 393 estimate the FOM in the current study, with the resulting values of 0.637 for the plastic and 0.892 for  
 394 the crystal scintillator sample.

## 395 5. Discussion and Conclusions

396 Given the increasing need for reliable neutron detection alternatives for  $^3\text{He}$  detectors, the authors  
 397 attempted to present a review of the most viable options available among the crystal scintillators.  
 398 Given the complexity of neutron detection, various methods are required to target specific neutron  
 399 energy range. Both organic and inorganic options were considered. Each group presents advantages  
 400 for certain application areas.

401 It appears that inorganic crystals utilising isotopes with high thermal neutron cross-section  
 402 (lithium, boron, gadolinium) provide a very good alternative for low energy neutron detectors.



**Figure 7.** Comparison of CCM plots for the two organic scintillator samples when exposed to  $^{252}\text{Cf}$  and data were collected with 500 MS/s digitiser: a) Cylindrical PSD Plastic from LLNL, and b) Single Stilbene Crystal. The upper plume is associated with gamma-ray interactions, whereas the lower plume with neutron events.

403 However, the manufacturing cost is still high, and the growing process is long. Fast neutron region, on  
 404 the other hand, has been targeted by organic scintillators for a long time, due to  $^1\text{H}$  content, which  
 405 allows elastic scattering of neutrons with a proton. Stilbene crystal is arguably the best available  
 406 scintillator detector capable of n/g separation. Nonetheless, growing large size detectors using stilbene  
 407 crystals is expensive in comparison to organic plastics and liquids.

408 There have been attempts to develop a neutron detector targeting a larger energy spectrum.  
 409 However, due to different mechanisms governing neutron interactions with matter at various energy  
 410 levels, this is not possible with a single material detector. Up to date literature reports on multi-detector  
 411 systems, where different detectors are used independently to detect specific group of neutrons. Readout  
 412 electronics attached to such system can combine the results into one system. Another method,  
 413 stemming from the multi-detector approach described, is based on composite detectors, where a  
 414 detector such as CLYC is incorporated into plastic scintillator to detect gammas, and thermal and fast  
 415 neutrons. Regardless of the target energy range, it is clear that scintillating crystals will continue to  
 416 play a key role in neutron detectors.

#### 417 5.1. Example of Neutron Detection Capabilities Using Single Stilbene Crystal

418 An example of detecting neutrons originating from  $^{252}\text{Cf}$  using organic solid state scintillators is  
 419 presented in Fig. 7. Due to scintillators' sensitivity to both neutrons and gamma-ray photons, both  
 420 particle types are detected resulting in two corresponding plumes. These tend to overlap slightly  
 421 in the low energy level. A significant overlap in that region leads to higher probability of particle  
 422 misclassification. As evidenced by the plots in Fig. 7, the overlap is most prominent in the low energy  
 423 region. In order to illustrate the difference in PSD performance in the low energy region between the  
 424 two scintillator sample, the low energy limit was set to 200 keVee. The high energy limit was set to  
 425 1800 keVee for both scintillators.

426 Based solely on the observation of the two graphs presented in Fig. 7, it is clear that the single  
 427 stilbene crystal (Fig. 7b) provides superior PSD, when compared with the LLNL plastic sample (Fig.  
 428 7a). Given that a similar number of pulses was accepted by the system for each scintillator, the shape  
 429 and intensity of the plumes appear quite dissimilar. Most importantly, the low energy cut-off point can  
 430 be observed at approx. 300 keVee for the single stilbene crystal. The corresponding cut-off point for  
 431 the plastic scintillator is found at approx. 400 keVee. Moreover, the overlap in the low energy area is

visibly smaller for the single stilbene than it is for plastic. The density of each plume is also higher for the stilbene crystal which again allows PSD to be performed with the higher level of accuracy.

These general observations agree with the quantitative analysis performed. The FOM was estimated for each detector, where 0.637 was observed for the plastic, and 0.892 for the single stilbene crystal. Despite various unique considerations required in the process of FOM estimation, presented results strongly support the claim that stilbene crystal is characterised by significantly superior PSD for fast neutron detection. The FOM estimated for stilbene crystal is considerably higher than the FOM value calculated for the plastic.

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