

Conditions for defocusing around more general metrics in Infinite Derivative Gravity

James Edholm¹

¹*Lancaster University*

Infinite Derivative Gravity is able to resolve the Big Bang curvature singularity present in general relativity by using a simplifying ansatz. We show that it can also avoid the Hawking-Penrose singularity, by allowing defocusing of null rays through the Raychaudhuri equation. This occurs not only in the minimal case where we ignore the matter contribution, but also in the case where matter plays a key role.

We investigate the conditions for defocusing for the general case where this ansatz applies and also for more specific metrics, including a general Friedmann-Robertson-Walker (FRW) metric and three specific choices of the scale factor which produce a bouncing FRW universe.

INTRODUCTION

The theory of General Relativity (GR) has been shown to describe gravity very accurately through a huge range of experimental tests over the past century [1]. However, it suffers from problems, generating both black hole and cosmological singularities [2].

Previous attempts to solve this problem include $f(R)$ gravity and higher derivative gravity. Higher derivative gravity suffers from the Ostrogradsky instability which produces ghosts [3], which are physical excitations with negative kinetic energy [4]. Infinite Derivative Gravity (IDG) solves this problem by adding an infinite sum of the d'Alembertian operator $\square = g^{\mu\nu}\nabla_\mu\nabla_\nu$, acting on the curvature. There is no highest derivative operator, so the Ostrogradsky instability does not apply. We must show that there are no ghosts in other ways. It has been shown that if the modification to the propagator contains at most a single pole then the problem of ghosts is avoided [5, 6].

Infinite derivative actions, which are used in string theory [7], were first applied to gravity by Biswas, Gerwick, Kovisto and Mazumdar [8]. IDG has been investigated around flat Minkowski backgrounds [9], (Anti) de Sitter backgrounds [10–13], a rotating metric [14] and the Schwarzschild black hole solution [15].

It is possible to find the gravitational entropy of a black hole within IDG [12, 16]. The propagator can also be found for this theory [5, 8, 10, 17–20], and it has been shown that one can curtail the divergences of 1 and 2 loop diagrams [21, 22], while other work has investigated the UV finiteness [23].

Within the Arnowitt-Deser Misner (ADM) decomposition, one can find the boundary terms of the theory [24]. IDG can be thought of as an extension to Starobinsky inflation [25–29], allowing us to put a constraint on the mass scale M [30]. It is also possible to put constraints on M using the deflection of light by the Sun [31], or by comparing the Newtonian potential to experimental evidence on the strength of gravity at small distances [32, 33].

GR suffers from the singularity problem. Due to the Einstein equations, the Raychaudhuri equations [34]

make it impossible for null rays to defocus as long as the Null Energy Condition holds, which by the Hawking-Penrose singularity theorem means there must be a singularity. IDG has different equations of motion to GR which generate the possibility of avoiding this singularity, through adding extra terms to the Einstein-Hilbert action. The conditions necessary to allow defocusing were investigated around a Minkowski background [33, 36, 37], around an (Anti) de Sitter background [38] and for an FRW metric near the bounce [39].

THE RAYCHAUDHURI CONDITION

The Raychaudhuri equation is a model-independent geometrical equation that tells us the expansion of a congruence of null rays emerging from the centre of our coordinate system, which have the tangent vectors k_μ where $k^\mu k_\mu = 0$ [34, 35]. If we imagine that the outgoing null rays form the surface of a sphere, then the expansion parameter $\theta = \nabla_\mu k^\mu$ describes the change in volume of that sphere. We would expect that for these outgoing null rays, the sphere would be expanding, but for a *trapped surface*, the volume actually decreases and $\theta < 0$. By the Raychaudhuri equation, θ fulfils the condition

$$\frac{d\theta}{d\tau} + \frac{1}{2}\theta^2 \leq -R_{\mu\nu}k^\mu k^\nu, \quad (1)$$

where $R_{\mu\nu}$ is the Ricci curvature tensor and τ is the affine parameter. Here we have taken the twist to be zero and we have used that the shear term is strictly positive to turn the Raychaudhuri equation into an inequality¹. If the null rays are to defocus, then the expansion parameter must be both positive and expanding, which implies that

$$R_{\mu\nu}k^\mu k^\nu < 0, \quad (2)$$

¹ The rotation term vanishes if we take the congruence of null rays to be orthogonal to a hypersurface and the shear term is strictly positive because the shear tensor is purely transverse [34, 35].

which we call the defocusing condition. Unless this condition is satisfied, the Hawking-Penrose singularity theorem says that a singularity will be generated [40].

The Null Energy Condition (NEC) says that for non-exotic matter, the stress-energy tensor $T_{\mu\nu}$ contracted with the tangent vectors is non-negative, i.e. $T_{\mu\nu}k^\mu k^\nu \geq 0$. If one inserts the Einstein equation into (2) it can be seen that the defocusing condition is not fulfilled and therefore by the Hawking-Penrose singularity theorem, GR generates singularities. The defocusing condition has also been investigated for f(R) gravity [41]. It was shown for perturbations around Minkowski and for a bouncing FRW solution near the bounce that IDG could allow de-

focusing of null rays [33, 36, 37, 39].

INFINITE DERIVATIVE GRAVITY

We will look at the IDG action

$$S = \frac{1}{2} \int d^4x [M_P^2 R + RF(\square)R - 2\Lambda], \quad (3)$$

where R is the Ricci curvature scalar, Λ is the cosmological constant with mass dimension 4 and M_P is the Planck mass. $F(\square) \equiv \sum_{n=0}^{\infty} f_n \square^n / M^{2n}$ is the infinite sum of the d'Alembertian operator $\square \equiv g^{\mu\nu} \nabla_\mu \nabla_\nu$, regulated by the mass scale M and with the dimensionless coefficients f_n . This action produces the equations of motion [44]

$$T_{\alpha\beta} = M_P^2 \left(R_{\alpha\beta} - \frac{1}{2} g_{\alpha\beta} R \right) + g_{\alpha\beta} \Lambda + 4 \left(R_{\alpha\beta} - \frac{1}{2} g_{\alpha\beta} R \right) F(\square)R + g_{\alpha\beta} RF(\square)R - 4 (\nabla_\alpha \nabla_\beta - g_{\alpha\beta} \square) F(\square)R + \sum_{n=1}^{\infty} f_n \sum_{l=0}^{\infty} [g_{\alpha\beta} ((\partial^\sigma \square^l R) \partial_\sigma \square^{n-l-1} R + \square^l R \square^{n-l} R) - 2 (\partial_\alpha \square^l R) \partial_\beta \square^{n-l-1} R]. \quad (4)$$

where $T_{\alpha\beta}$ is the stress-energy tensor and $R_{\alpha\beta}$ is the Ricci curvature tensor. If we contract (4) with tangent vectors $k^\alpha k^\beta$, where $k^\alpha k_\alpha = 0$, we can find the full condition for null rays to defocus (2), which allows us to avoid the Hawking-Penrose singularity²³

$$k^\alpha k^\beta R_{\alpha\beta} = \frac{1}{M_P^2 + 4F(\square)R} \left[k^\beta k^\alpha T_{\alpha\beta} + 4k^\beta k^\alpha \nabla_\alpha \nabla_\beta F(\square)R + 2k^\beta k^\alpha \sum_{n=1}^{\infty} \frac{f_n}{M^{2n}} \sum_{l=0}^{n-1} (\partial_\alpha \square^l R) \partial_\beta \square^{n-l-1} R \right] < 0. \quad (5)$$

The simplifying ansatz

In the next section we will show IDG also allows null rays to defocus for solutions where the d'Alembertian operator \square acting on the Ricci scalar R fulfills the ansatz $\square R = r_1 R + r_2$, where r_1 and r_2 are constants, as studied in previous work [5, 6, 29, 42, 43]. In fact, [29] showed that a metric of this form was the most general spatially flat FRW solution within IDG.

Using this ansatz, [5, 6, 42, 43] were able to find vacuum solutions to the equations of motion (4). These fo-

cused on generating bouncing FRW cosmologies, in particular solutions where $a(t) = \cosh(\sigma t)$ and $a(t) = e^{\frac{\lambda}{2} t^2}$. A simple toy model $a(t) = 1 + a_2 t^2$ was studied as a perturbation to flat space in [39]. These bouncing solutions exhibit time symmetry about $t = 0$ and avoid the curvature singularity at $t = 0$ present in GR, known as the Big Bang singularity problem. We will examine the most general metric where this ansatz is satisfied, and then look at these more specific metrics, showing that we can avoid the singularity generated by General Relativity through the Hawking-Penrose singularity theorem [40].

Minimum defocusing condition using the ansatz

If we take the ansatz $\square R = r_1 R + r_2$, then for $n > 0$, $\square^n R = r_1^n \left(R + \frac{r_2}{r_1} \right)$ and

$$F(\square)R = \sum_{n=0}^{\infty} f_n \frac{\square^n R}{M^{2n}} = F(r_1) \left(R + \frac{r_2}{r_1} \right) - f_0 r_2 / r_1, \quad (6)$$

where f_0 is the zero-order coefficient of $F(\square)$, i.e. the coefficient of R^2 in the action. The minimum defocusing condition (5) becomes⁴

be clearly seen in the denominator of (7). It is not explicitly seen in the numerator of (7) because the derivatives of r_2 vanish, so $r_1 R = -r_2$ implies that the derivatives of $r_1 R$ also vanish, so the higher order terms do not contribute.

² (5) reduces to the defocusing condition found for a bouncing FRW cosmology [39] if we take the Ricci scalar R to be a function of t only and consider only even powers of t (as happens at the bounce), so that $H \rightarrow 0$ and take $\alpha = 1/2$.

³ The cosmological constant does not feature in (5) due to the condition $k^\alpha k^\beta g_{\alpha\beta} = 0$.

⁴ We can return to a local theory by taking $r_1 R = -r_2$, which removes the dependence on higher orders of \square in (6). This can

$$\frac{1}{M_p^2 + 4F(r_1) \left(R + \frac{r_2}{r_1} \right) - 4f_0 r_2 / r_1} \left[k^\beta k^\alpha T_{\alpha\beta} + 4k^\beta k^\alpha \sum_{n=0}^{\infty} \frac{f_{1n}}{M^{2n}} \nabla_\alpha \nabla_\beta (r_1^n R) + 2k^\beta k^\alpha \sum_{n=1}^{\infty} \frac{f_n}{M^{2n}} \sum_{l=0}^{n-1} \partial_\alpha (r_1^l R) \partial_\beta (r_1^{n-l-1} R) \right] < 0, \quad (7)$$

where we are looking at vacuum solutions, so neglect the stress-energy tensor term and have also removed the derivatives of the constants r_1 and r_2 . We can simplify (7). r_1^{n-2} and the partial derivatives of R have no l de-

pendence and can therefore be pulled out of the sum. By observing that $\sum_{l=0}^{n-1} 1 = n$ and $k^\beta k^\alpha (\partial_\alpha R) \partial_\beta R = (k^\alpha \partial_\alpha R)^2$, our final condition for the conditions for defocusing for any metric which fulfils $\square R = r_1 R + r_2$ can be written as⁵

$$\frac{1}{M_p^2 + 4F(r_1) \left(R + \frac{r_2}{r_1} \right) - 4f_0 r_2 / r_1} \left[k^\beta k^\alpha T_{\alpha\beta} + 4 \sum_{n=0}^{\infty} \frac{f_{1n}}{M^{2n}} r_1^n k^\beta k^\alpha \nabla_\alpha \nabla_\beta R + 2 \sum_{n=1}^{\infty} \frac{n f_n}{M^{2n}} r_1^{n-1} (k^\alpha \partial_\alpha R)^2 \right] < 0, \quad (8)$$

or more succinctly as⁶

$$\frac{1}{M_p^2 + 4F(r_1) \left(R + \frac{r_2}{r_1} \right) - 4f_0 \frac{r_2}{r_1}} \left[k^\beta k^\alpha T_{\alpha\beta} + 4F(r_1) k^\beta k^\alpha \nabla_\alpha \nabla_\beta R + 2F'(r_1) (k^\alpha \partial_\alpha R)^2 \right] < 0, \quad (9)$$

where $F'(\square)$ is defined as

$$F'(\square) \equiv \sum_{n=1}^{\infty} \frac{n f_n}{M^{2n}} \square^{n-1}. \quad (10)$$

THE FRW METRIC

FRW metrics describe universes that are time-dependent, homogenous and isotropic. These metrics can suffer from curvature singularities at $t = 0$. We will investigate some bouncing universes that can avoid these curvature singularities and see whether they can avoid the Hawking-Penrose singularity.

An FRW metric in spherical coordinates takes the form

$$ds^2 = -dt^2 + a^2(t) \left(\frac{dr^2}{1 - \kappa r^2} + r^2 d\Omega^2 \right), \quad (11)$$

where $a(t)$ is the scale factor of the universe and κ is the

spatial curvature.

Using the condition on the tangent vectors $k^\mu k_\mu = 0$ gives (noting that $R = R(t)$ even if the spatial curvature κ is non-zero)

$$k^\beta k^\alpha \nabla_\alpha \nabla_\beta R(t) = (k^0)^2 (\partial_0^2 - H \partial_0) R(t). \quad (12)$$

We can also simplify the defocusing condition (9) by noting that for an FRW metric, the background d'Alembertian acting on a time-dependent scalar quantity $S(t)$ is given by $\square S(t) = -\partial_0^2 S(t) - 3H \partial_0 S(t)$, so $\partial_0^2 S(t) - H \partial_0 S(t) = -\square S(t) - 4H \partial_0 S(t)$, where H is the Hubble parameter. Therefore by comparing this to the ansatz, we find $\partial_0^2 R(t) - H \partial_0 R(t) = -r_1 R(t) - 4H \partial_0 R(t)$. We can divide by $(k^0)^2$ as it is strictly positive and use that $H \approx 0$ near the bounce point. Finally, note that for FRW, the null tangent vectors contracted with the stress-energy tensor gives $\kappa^\mu \kappa^\nu T_{\mu\nu} = (k^0)^2 (\rho + p)$ where ρ and p are density and pressure respectively.

Thus the defocusing condition for an FRW metric near the bounce fulfilling $\square R = r_1 R + r_2$ is

$$\frac{1}{M_p^2 + 4F(r_1) \left(R + \frac{r_2}{r_1} \right) - 4f_0 \frac{r_2}{r_1}} \left[2F(r_1) (r_1 R + r_2) - F'(r_1) (\partial_0 R)^2 - (\rho + p) \right] > 0. \quad (13)$$

⁵ This matches what we find if we apply the ansatz to perturbations around a Minkowski or de Sitter background when we

where f_0 is the zeroth order coefficient of $F(r_1) = F(\square = r_1)$. Note that (13) does not depend explicitly on the curvature κ , which is encoded into the Ricci scalar $R(t)$.

SIMPLIFYING OUR CONDITION

We can simplify the defocusing condition by placing conditions on $F(\square)$ and $F'(\square)$ by inserting a simple solution to the equations of motion.

Inserting the equations of motion

The vacuum trace equation for metrics satisfying the ansatz is [6, 42, 43]

$$A_1 R + A_2 (2r_1 R^2 + \partial_\mu R + \partial^\mu R) + A_3 = 0, \quad (14)$$

where the coefficients A_i are given by

$$\begin{aligned} A_1 &= -M_P^2 + 4F'(r_1)r_2 - 2\frac{r_2}{r_1}(F(r_1) - f_0) + 6F(r_1)r_1, \\ A_2 &= F'(r_1), \\ A_3 &= 4\Lambda + \frac{r_2}{r_1}(M_P^2 + A_1) - 2\frac{r_2^2}{r_1}F'(r_1). \end{aligned} \quad (15)$$

The simplest solution to (14) is to assume $A_i = 0$ which gives the identities

$$\begin{aligned} F'(r_1) &= 0, \quad r_2 = -\frac{r_1(M_P^2 - 6F(r_1)r_1)}{2(F(r_1) - f_0)}, \\ \Lambda &= -\frac{r_2 M_P^2}{4r_1} = M_P^2 \frac{(M_P^2 - 6F(r_1)r_1)}{8(F(r_1) - f_0)}, \end{aligned} \quad (16)$$

which allows us to remove f_0 from our defocusing condition (13), which becomes

$$\frac{2F(r_1)(r_1 R + r_2) - (\rho + p)}{F(r_1)(4R + 6r_1) - M_P^2} > 0, \quad (17)$$

where again the spatial curvature κ does not feature explicitly, but is encoded within R .

In the early universe, we expect the curvature to be very large and we can assume $r_1, r_2 \ll R$, so

$$\frac{2r_1 R F(r_1) - (\rho + p)}{4R F(r_1) - M_P^2} > 0. \quad (18)$$

There are two cases where the defocusing condition is fulfilled

1. $2r_1 R F(r_1) - (\rho + p) > 0$ and $4R F(r_1) - M_P^2 > 0$
2. $2r_1 R F(r_1) - (\rho + p) < 0$ and $4R F(r_1) - M_P^2 < 0$

Note that if $\rho + p$ is very large, as it would be at the bounce point, then it is unsatisfactory (but possible) for this to be cancelled out by an even larger curvature term. However, we do note that it would be natural for the curvature to be large at this time. It is indeed pleasing that in the second case, having $r_1 > 0$ and $R F(r_1) < 0$ means that the condition is fulfilled and the singularity is avoided.

We have shown that it is possible to not only avoid the singularity in the minimal case where we ignore matter, but actually for the matter to aid in the defocusing! We now look at examples where this can apply.

SPECIFIC BOUNCING MODELS

Toy model

If we take a simplistic flat bouncing model with the scale factor

$$a(t) = 1 + a_2 t^2, \quad (19)$$

which fulfills the ansatz with $r_1 = -6a_2$ and $r_2 = 48a_2^2$ and was investigated in [39].

The defocusing condition for the simple solution to the equations of motion (17) at small times and neglecting matter becomes

$$\frac{M_P^2 + 16a_2 f_0 - \frac{5}{12}(\rho + p)}{M_P^2 + 96f_0 a_2} > 0 \quad (20)$$

which is fulfilled for $-\frac{M_P^2}{16} < a_2 f_0 < -\frac{M_P^2}{96}$.

Assuming a cosh solution

In this section we assume the solution $a(t) = \cosh(\sigma t)$, where σ is a constant. This vacuum solution to the equations of motion was investigated in [5, 6, 43] and return to the de Sitter metric at late times. This solution requires a cosmological constant and radiation. This metric fulfills the ansatz ⁷ $\square R = r_1 R + r_2$ with $r_1 = 2\sigma^2$ and $r_2 = -24\sigma^4$.

Therefore the defocusing condition (13) at $t \approx 0$ near the bounce point for an FRW metric with the scale factor

discard the appropriate terms [36, 37].

⁶ This correctly reduces to $R + R^2$ gravity [45] in the limit $r_1 = r_2 = 0$.

⁷ Note that r_1 and r_2 have opposite signs for both the simple solution $1 + t^2$ and the cosh solution. Therefore by (16), this is equivalent to a positive Λ , whereas for the exponential bouncing solution, a negative Λ is produced.

$a(t) = \cosh(\sigma t)$ and non-zero spatial curvature κ is⁸

$$\frac{24\sigma^2 (\kappa - \sigma^2) F(r_1) - (\rho + p)}{M_P^2 + 24F(2\sigma^2)(\kappa - \sigma^2) + 48\sigma^2 f_0} < 0. \quad (21)$$

Taking the simple solution (16) to the equations of motion and neglecting matter, (21) is fulfilled for zero spatial curvature κ if $-\frac{1}{24} \frac{M_P^2}{\sigma^2} < f_0 < -\frac{1}{96} \frac{M_P^2}{\sigma^2}$. If matter is included, and gives a very large contribution, then defocusing occurs for $f_0 < -\frac{1}{96} \frac{M_P^2}{\sigma^2}$.

We have shown that it is possible for a bouncing cosh solution, which was already shown to avoid the Big Bang curvature singularity, to avoid the Hawking-Penrose singularity for certain curvature.

Exponential bouncing solution

Finally we look at a flat FRW metric with the exponential bouncing scale factor which was studied in [42, 43] where

$$a(t) = e^{\frac{\lambda}{2} t^2}. \quad (22)$$

where λ is a positive constant with dimensions of mass squared. This scale factor gives the Ricci scalar $R = 3\lambda(1 + \lambda t^2)$ which fulfills the ansatz with $\square R(t) = -6\lambda R(t) - 12\lambda^2$. The defocusing condition (13) for small times becomes

$$\frac{5\lambda^2 F(r_1) + (\rho + p)}{M_P^2 + 20\lambda F(r_1) - 8f_0\lambda} < 0. \quad (23)$$

and with the simple solution to the equations of motion (16), we find that there is defocusing for $-\frac{1}{4} M_P^2 < \lambda f_0 < \frac{1}{28} M_P^2$ if we neglect matter, or $\lambda f_0 < \frac{1}{28} M_P^2$ if the contribution of matter is large.

The exponential bouncing solution (22) generated using ghost-free Infinite Derivative Gravity was already known to avoid the Big Bang curvature singularity, but we have shown it can avoid the Hawking-Penrose singularity, and that the addition of matter can actually aid this process.

CONCLUSION

We have shown that IDG can allow defocusing of null rays, and thus avoid Hawking-Penrose singularities, for more general spacetimes as long as our simplifying ansatz holds. We first looked at completely general spacetimes where this ansatz applies, before restricting ourselves to time-dependent Ricci scalars and then time-dependent

metrics. We looked at an FRW metric and showed that this could defocus given certain conditions.

We investigated three specific scale factors within an FRW metric. We first looked at a simple polynomial bouncing universe, then a cosh scale factor and finally an example of a exponentially bouncing time-dependent metric. We showed that all of these metrics allowed defocusing in the minimal case, but they also allowed defocusing in the case where matter gave a large contribution, as in the early universe.

It has previously been shown that IDG would permit a bouncing universe which avoided a curvature singularity. We have now shown that it is also possible to avoid the Hawking-Penrose singularity in these universes.

ACKNOWLEDGEMENTS

We would particularly like to thank Anupam Mazumdar for his guidance and suggestions on this project. We would also like to thank Alexey Koshelev and Aindriu Conroy for their help on this work.

-
- [1] C. M. Will, *Living Rev. Rel.* **17** (2014) 4 doi:10.12942/lrr-2014-4 [arXiv:1403.7377 [gr-qc]].
 - [2] S. W. Hawking and G. F. R. Ellis, “The Large Scale Structure of Space-Time,” doi:10.1017/CBO9780511524646
 - [3] N. Barnaby and N. Kamran, “Dynamics with infinitely many derivatives: The Initial value problem,” *JHEP* **0802** (2008) 008 doi:10.1088/1126-6708/2008/02/008 [arXiv:0709.3968 [hep-th]].
 - [4] P. Van Nieuwenhuizen, “On ghost-free tensor lagrangians and linearized gravitation,” *Nucl. Phys. B* **60** (1973) 478. doi:10.1016/0550-3213(73)90194-6
 - [5] T. Biswas, A. Mazumdar and W. Siegel, “Bouncing universes in string-inspired gravity,” *JCAP* **0603** (2006) 009 doi:10.1088/1475-7516/2006/03/009 [hep-th/0508194].
 - [6] T. Biswas, T. Koivisto and A. Mazumdar, “Towards a resolution of the cosmological singularity in non-local higher derivative theories of gravity,” *JCAP* **1011** (2010) 008 doi:10.1088/1475-7516/2010/11/008 [arXiv:1005.0590 [hep-th]].
 - [7] A. A. Tseytlin, “On singularities of spherically symmetric backgrounds in string theory,” *Phys. Lett. B* **363** (1995) 223 doi:10.1016/0370-2693(95)01228-7 [hep-th/9509050].
 - [8] T. Biswas, E. Gerwick, T. Koivisto and A. Mazumdar, “Towards singularity and ghost free theories of gravity,” *Phys. Rev. Lett.* **108**, 031101 (2012) doi:10.1103/PhysRevLett.108.031101 [arXiv:1110.5249 [gr-qc]].
 - [9] A. S. Koshelev and A. Mazumdar, “Absence of event horizon in massive compact objects in infinite derivative gravity,” arXiv:1707.00273 [gr-qc].
 - [10] T. Biswas, A. S. Koshelev and A. Mazumdar, “Gravitational theories with stable (anti-)de Sitter backgrounds,”

⁸ Note that (21) matches (13) when we take $r_1 = 2\sigma^2$ and $r_2 = -24\sigma^4$.

- Fundam. Theor. Phys. **183** (2016) 97 doi:10.1007/978-3-319-31299-6 [arXiv:1602.08475 [hep-th]].
- [11] T. Biswas, A. S. Koshelev and A. Mazumdar, “Consistent higher derivative gravitational theories with stable de Sitter and anti de-Sitter backgrounds,” *Phys. Rev. D* **95** (2017) no.4, 043533 doi:10.1103/PhysRevD.95.043533 [arXiv:1606.01250 [gr-qc]].
- [12] A. Conroy, A. Mazumdar, S. Talaganis and A. Teimouri, “Nonlocal gravity in D dimensions: Propagators, entropy, and a bouncing cosmology,” *Phys. Rev. D* **92** (2015) no.12, 124051 doi:10.1103/PhysRevD.92.124051 [arXiv:1509.01247 [hep-th]].
- [13] J. Edholm, “Revealing Infinite Derivative Gravity’s true potential: The weak-field limit around de Sitter backgrounds,” arXiv:1801.00834 [gr-qc].
- [14] A. S. Cornell, G. Harmsen, G. Lambiasi and A. Mazumdar, “Rotating metric in Non-Singular Infinite Derivative Theories of Gravity,” arXiv:1710.02162 [gr-qc].
- [15] G. Calcagni and L. Modesto, “Stability of Schwarzschild singularity in non-local gravity,” *Phys. Lett. B* **773** (2017) 596 doi:10.1016/j.physletb.2017.09.018 [arXiv:1707.01119 [gr-qc]].
- [16] A. Conroy, A. Mazumdar and A. Teimouri, “Wald Entropy for Ghost-Free, Infinite Derivative Theories of Gravity,” *Phys. Rev. Lett.* **114**, no. 20, 201101 (2015) doi:10.1103/PhysRevLett.114.201101 [arXiv:1503.05568 [hep-th]].
- [17] E. Tomboulis, “Renormalizability and Asymptotic Freedom in Quantum Gravity,” *Phys. Lett. B* **97**, 77 (1980).
- [18] L. Buoninfante, “Ghost and singularity free theories of gravity,” arXiv:1610.08744 [gr-qc].
- [19] W. Siegel, “Stringy gravity at short distances,” hep-th/0309093.
- [20] T. Biswas, T. Koivisto and A. Mazumdar, “Nonlocal theories of gravity: the flat space propagator,” arXiv:1302.0532 [gr-qc].
- [21] S. Talaganis, T. Biswas and A. Mazumdar, “Towards understanding the ultraviolet behavior of quantum loops in infinite-derivative theories of gravity,” *Class. Quant. Grav.* **32** (2015) no.21, 215017 doi:10.1088/0264-9381/32/21/215017 [arXiv:1412.3467 [hep-th]].
- [22] S. Talaganis, “Quantum Loops in Non-Local Gravity,” *PoS CORFU 2014* (2015) 162 [arXiv:1508.07410 [hep-th]].
- [23] S. Talaganis, “Towards UV Finiteness of Infinite Derivative Theories of Gravity and Field Theories,” arXiv:1704.08674 [hep-th].
- [24] A. Teimouri, S. Talaganis, J. Edholm and A. Mazumdar, “Generalised Boundary for Higher Derivative Theories of Gravity,” arXiv:1606.01911 [gr-qc]. *JHEP* **1608** (2016) 144 doi:10.1007/JHEP08(2016)144
- [25] B. Craps, T. De Jonckheere and A. S. Koshelev, *JCAP* **1411** (2014) no.11, 022 doi:10.1088/1475-7516/2014/11/022 [arXiv:1407.4982 [hep-th]].
- [26] F. Briscese, A. Marcian, L. Modesto and E. N. Saridakis, “Inflation in (Super-)renormalizable Gravity,” *Phys. Rev. D* **87** (2013) no.8, 083507 doi:10.1103/PhysRevD.87.083507 [arXiv:1212.3611 [hep-th]].
- [27] A. S. Koshelev, L. Modesto, L. Rachwal and A. A. Starobinsky, “Occurrence of exact R^2 inflation in non-local UV-complete gravity,” *JHEP* **1611** (2016) 067 doi:10.1007/JHEP11(2016)067 [arXiv:1604.03127 [hep-th]].
- [28] T. Biswas, A. S. Koshelev, A. Mazumdar and S. Y. Vernov, “Stable bounce and inflation in non-local higher derivative cosmology,” *JCAP* **1208** (2012) 024 doi:10.1088/1475-7516/2012/08/024 [arXiv:1206.6374 [astro-ph.CO]].
- [29] A. S. Koshelev, K. Sravan Kumar and A. A. Starobinsky, “ R^2 inflation to probe non-perturbative quantum gravity,” arXiv:1711.08864 [hep-th].
- [30] J. Edholm, “UV completion of the Starobinsky model, tensor-to-scalar ratio, and constraints on non-locality,” *Phys. Rev. D* **95** (2017) no.4, 044004 doi:10.1103/PhysRevD.95.044004 [arXiv:1611.05062 [gr-qc]].
- [31] L. Feng, “Light Bending in the Infinite Derivative Theories of Gravity,” arXiv:1703.06535 [gr-qc].
- [32] J. Edholm, A. S. Koshelev and A. Mazumdar, “Universality of testing ghost-free gravity,” arXiv:1604.01989 [gr-qc]. *Phys. Rev. D* **94** (2016) no.10, 104033 doi:10.1103/PhysRevD.94.104033
- [33] A. Conroy and J. Edholm, “Newtonian Potential and Geodesic Completeness in Infinite Derivative Gravity,” arXiv:1705.02382 [gr-qc].
- [34] Robert M. Wald, “General Relativity,” Chicago, Usa: Univ. Pr. (1984)
- [35] Eric Poisson, “A Relativist’s Toolkit: The Mathematics of Black-Hole Mechanics,” Cambridge University Press, (2004)
- [36] A. Conroy, A. S. Koshelev and A. Mazumdar, “Criteria for resolving the cosmological singularity in Infinite Derivative Gravity,” arXiv:1605.02080 [gr-qc].
- [37] A. Conroy, “Infinite Derivative Gravity: A Ghost and Singularity-free Theory,” arXiv:1704.07211 [gr-qc].
- [38] J. Edholm and A. Conroy, “Criteria for resolving the cosmological singularity in Infinite Derivative Gravity around expanding backgrounds,” *Phys. Rev. D* **96** (2017) no.12, 124040 doi:10.1103/PhysRevD.96.124040 arXiv:1710.01366 [gr-qc].
- [39] A. Conroy, A. S. Koshelev and A. Mazumdar, “Geodesic completeness and homogeneity condition for cosmic inflation,” *Phys. Rev. D* **90** (2014) no.12, 123525 doi:10.1103/PhysRevD.90.123525 [arXiv:1408.6205 [gr-qc]].
- [40] S. Hawking and G. Ellis, “The Large Scale Structure of Space-Time,” Cambridge University Press (1973)
- [41] S. Carloni, P. K. S. Dunsby and D. M. Solomons, “Bounce conditions in f(R) cosmologies,” *Class. Quant. Grav.* **23** (2006) 1913 doi:10.1088/0264-9381/23/6/006 [gr-qc/0510130].
- [42] A. S. Koshelev and S. Y. Vernov, “On bouncing solutions in non-local gravity,” *Phys. Part. Nucl.* **43** (2012) 666 doi:10.1134/S106377961205019X [arXiv:1202.1289 [hep-th]].
- [43] A. S. Koshelev, “Stable analytic bounce in non-local Einstein-Gauss-Bonnet cosmology,” *Class. Quant. Grav.* **30** (2013) 155001 doi:10.1088/0264-9381/30/15/155001 [arXiv:1302.2140 [astro-ph.CO]].
- [44] T. Biswas, A. Conroy, A. S. Koshelev and A. Mazumdar, “Generalized ghost-free quadratic curvature gravity,” *Class. Quant. Grav.* **31** (2014) 015022 Erratum: [*Class. Quant. Grav.* **31** (2014) 159501] doi:10.1088/0264-9381/31/1/015022, 10.1088/0264-9381/31/15/159501 [arXiv:1308.2319 [hep-th]].
- [45] F. D. Albareti, J. A. R. Cembranos, A. de la Cruz-Dombriz and A. Dobado, “On the non-

attractive character of gravity in $f(R)$ theories,” JCAP

1307 (2013) 009 doi:10.1088/1475-7516/2013/07/009
[arXiv:1212.4781 [gr-qc]].