

Euclid Quick Data Release (Q1)

The role of cosmic connectivity in shaping galaxy clusters

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ABSTRACT

The matter around galaxy clusters is distributed over several filaments, reflecting their positions as nodes in the large-scale cosmic web. The number of filaments connected to a cluster, i.e. its connectivity, is expected to affect the physical properties of clusters. Using the first *Euclid* galaxy catalogue from the *Euclid* Quick Release 1 (Q1), we investigated the connectivity of galaxy clusters and how it correlates with their physical and galaxy member properties. Around 220 clusters located within the three fields of Q1 (covering $\sim 63 \text{ deg}^2$) were analysed in the redshift range $0.2 < z < 0.7$. Due to the photometric redshift uncertainty, we reconstructed the cosmic web skeleton, and measured the cluster connectivity, in 2D projected slices with a thickness of 170 comoving $h^{-1} \text{ Mpc}$ and centred on each cluster redshift, by using two different filament finder algorithms on the most massive galaxies ($M_* > 10^{10.3} M_\odot$). In agreement with previous measurements, we recovered the mass-connectivity relation independently of the filament detection algorithm, showing that the most massive clusters are, on average, connected to a larger number of cosmic filaments, consistent with hierarchical structure formation models. Furthermore, we explored the possible correlations between connectivities and two cluster properties: the fraction of early-type galaxies and the Sérsic index of galaxy members. Our result suggests that the clusters populated by early-type galaxies exhibit higher connectivity compared to clusters dominated by late-type galaxies. These preliminary investigations highlight our ability to quantify the impact of the cosmic web's connectivity on cluster properties with *Euclid*.

Key words. Cosmology: observations – Galaxies: cluster: general – large-scale structure of Universe – Methods: statistical

1. Introduction

The hierarchical model of matter assembly dictates not only the formation and evolution of galaxy clusters, but also their positioning within the large-scale cosmic web, comprising filaments, sheets, and voids. Observations (Jêveer et al. 1978; Barrow et al. 1985; de Lapparent et al. 1986; Sousbie et al. 2008), theory (Zel'Dovich 1970; White & Silk 1979; Bond et al. 1996), and cosmological simulations (Klypin & Shandarin 1983; Davis et al. 1985; Thomas & Couchman 1992; Springel et al. 2005; Angulo & Hahn 2022) have demonstrated that galaxy clusters reside at the nodes of this vast network. They grow through mergers (Bardeen et al. 1986; Kauffmann et al. 1993) and anisotropic accretion of matter from their connected filaments (Bond et al. 1996; Aragón-Calvo et al. 2010; Pichon et al. 2011; Gouin et al. 2017; Vurm et al. 2023). Thus, galaxy clusters serve as unique tracers and laboratories for the study of the structure of the cosmic web, and its environmentally driven influence on galaxy evolution in the densest regions of the Universe.

In the context of cosmic web topology, connectivity—defined as the number of filaments linked to a cosmic node—is one aspect of the broader spectrum of morphological elements (see e.g. Codis et al. 2018, for formal definitions). As a key statistical parameter, it probes the geometrical structure of the cosmic web surrounding a cluster. Currently, increasing numbers of investigations on cluster connectivity are being undertaken to explore the influence of filamentary accretion on the physical and dynamical properties of galaxy clusters, hence yielding better constraints on structure formation. Theoretically, it was shown that mass primarily controls the halo connectivity (Pichon et al. 2010; Aragón-Calvo et al. 2010; Codis et al. 2018), such that massive clusters have a larger number of connected filaments than low-mass ones. Observationally, Einasto et al. (2018a) have shown that groups in superclusters also have higher connectivity than groups of the same richness in voids (see also Einasto et al. 2018b; Sarron et al. 2019). This was explained from first principle by Cadiou et al. (2020). The mass–connectivity relation can

be explained by the hierarchical structure formation scenario, in which clusters are the result of merging haloes (Codis et al. 2018). By probing mass assembly history in simulations, Darragh Ford et al. (2019) have shown that major merging events increase the connectivity of haloes (see also Lee et al. 2021; Galárraga-Espinosa et al. 2024, for connectivity evolution with redshift). In agreement with this cluster evolutionary picture, Gouin et al. (2021) showed that, regardless of halo mass, highly connected objects grow faster than low-connected ones, linking halo connectivity to its dynamical state captured by its relaxation level. However, this finding is debated by Santoni et al. (2024), owing to differences in filament-finding methods, cosmic web tracers (galaxies versus gas), and the physics of the underlying cosmological simulations. Regarding halo shapes, cosmological simulations predict that cluster shapes tend to align preferentially with their main connected filaments (Chisari et al. 2015; Gouin et al. 2017; Okabe et al. 2020a; Kuchner et al. 2020; Morinaga & Ishiyama 2020), a trend increasingly supported by observational evidence (Einasto et al. 2020; Okabe et al. 2020b; Gouin et al. 2020; Smith et al. 2023).

In addition to its impact on the clusters, the connectivity of haloes is a relevant ingredient for galaxy evolution. At large scales, Poudel et al. (2017) suggested that cosmic filaments play a role in shaping the properties of groups and their central galaxies. Darragh Ford et al. (2019) suggested that high connectivity might be the result of past mergers, which in turn boost the growth of the supermassive black hole and the active galactic nucleus (AGN) feedback (Dubois et al. 2013), and quench the central galaxy. Interestingly, investigations at smaller scales tend to relate galaxy properties with connectivity. By using both hydrodynamical simulations and SDSS-DR10 observations, Kraljic et al. (2020) showed that more massive, less star-forming, and less rotation-supported galaxies tend to have higher connectivity (see also Tillson et al. 2015; Galárraga-Espinosa et al. 2023).

To study the role of connectivity in shaping the properties of clusters and galaxies, it is crucial to reconstruct the skeleton of the cosmic web. However, this requires addressing two key challenges: the definition of filaments and their detection

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(see e.g. Libeskind et al. 2017, for a review of filament finders). Over the past decade, various filament-finding techniques have been developed and applied to both simulations and observations. Among these techniques, DISPERSE (DIScrete PERsistent Structure Extractor, Sousbie 2011) and T-REX (Tree-based Ridge Extractor, Bonnaire et al. 2020) are widely used on discrete data. In DISPERSE, filaments are defined as links between maxima and saddles in the density field, using a topological segmentation of the galaxy catalogue that consistently identifies all geometric features of the cosmic web (walls, voids, filaments, peaks). Only the most topologically robust structures are retained. In contrast, T-REX defines filaments as a graph-based tree structure modelling the cosmic web skeleton. These algorithms allow for filament detection in observations. As an alternative, the MMF/Nexus formalism is an explicitly multi-scale method used to analyse the cosmic web in simulations (Aragón-Calvo et al. 2007; Cautun et al. 2014). Dynamically motivated approaches have also been proposed for some time (Arnold et al. 1982; Bond & Myers 1996; Feldbrugge & van de Weygaert 2024) and offer valuable insights into the formation of the cosmic web. However, these methods are less straightforward to implement in observational datasets.

Different studies have tested the capability of DISPERSE to detect filaments around clusters. New generations of surveys such as the WEAVE Wide Field Cluster Survey (Kuchner et al. 2021; Cornwell et al. 2022; Kraljic et al. 2022; Cornwell et al. 2023), and *Euclid* (Euclid Collaboration: Mellier et al. 2025) have motivated a renewed interest. In contrast to detections based on the galaxy distribution, other recent studies have attempted to detect filaments from the gas density field in simulations (Schimd et al. 2024; Santoni et al. 2024), and in X-ray observations (Sousbie 2011; Gallo et al. 2024).

In this context, while cosmic web connectivity is increasingly well understood in Lagrangian space (Codis et al. 2018), large cosmological dark matter simulations (Aragón-Calvo et al. 2010; Codis et al. 2012) and hydrodynamical cosmological simulations (Darragh Ford et al. 2019; Kraljic et al. 2020; Lee et al. 2021; Gouin et al. 2021) measuring it in observations remains a challenge. Beyond detecting individual filaments connected to clusters or superclusters (Einasto et al. 2020; Malavasi et al. 2020; Aghanim et al. 2024; Gallo et al. 2024), large and representative statistical samples of groups and clusters of galaxies with measured connectivities are still rare (Darragh Ford et al. 2019; Sarron et al. 2019; Kraljic et al. 2020); to date, this has hindered a full exploration of the impact of cosmic web environments on cluster evolution. The situation will change with the ongoing sky survey by *Euclid*, which will eventually provide us with the largest sample of clusters of galaxies containing hundreds of thousands sources (Sartoris et al. 2016) and with about ten billion galaxies used to reconstruct the cosmic web skeleton.

For this work, we used the very first data from the *Euclid* Quick Release Q1 (2025) to explore the connectivity of galaxy clusters as a function of cluster mass and galaxy member properties, and we showcase the capabilities of future analyses of *Euclid* data when the survey is completed. To do so, we combined a sample of about 220 already-known clusters, detected in optical and X-ray surveys (Voges et al. 1999; York et al. 2000; Flaughner et al. 2015; Merloni et al. 2012), with the skeleton of the cosmic web surrounding them and derived from the galaxy distribution in *Euclid* (Euclid Collaboration: Laigle et al. 2025). In this study, filament finders are applied on a selected sample of galaxies in 2D projected slices centred on each cluster redshift. The paper is organised as follows. In Sect. 2 the selected sample of clusters and the method used to reconstruct the cosmic web

skeleton around them are described. In Sect. 3 we explore the connectivity of clusters as a function of cluster properties such as cluster mass and galaxy member properties. Finally, in Sect. 4 we summarise our main results. We consider here a flat Λ CDM cosmology with cosmological parameters from the *Planck* mission (Planck Collaboration et al. 2014), namely $\Omega_\Lambda = 0.693$, $\Omega_m = 0.307$, $\Omega_b = 0.04825$, $\sigma_8 = 0.8288$, and $h = 0.6777$.

2. Cosmic web extraction around clusters

2.1. Q1 photometric data catalogue

We present in this section our cluster selection, and the method for extracting the cosmic web skeleton around them in Q1 data, which are divided in three patches: the Euclid Deep Field North of 20 deg^2 (EDF-N), the Euclid Deep Field South of 23 deg^2 (EDF-S), and the Euclid Deep Field Fornax of 10 deg^2 (EDF-F) (Euclid Collaboration: Aussel et al. 2025). In the Q1 dataset, those fields are representative of the Euclid Wide Survey (EWS) in terms of their 5σ depths: 26.0, 23.8, 24.0, and 24.0 in the I_e , Y_e , J_e , and H_e filters, respectively (Euclid Collaboration: Tucci et al. 2025). Galaxy number counts in the band for each of the fields are presented in Euclid Collaboration: McCracken et al. (2025).

2.2. Cluster selection

We focused on the four publicly available catalogues that contain clusters in the *Euclid* Q1 fields. These are eROSITA (Bulbul et al. 2024), MCXC (Piffaretti et al. 2011), DES-Y1 (Rykoff et al. 2016), and WHL-SDSS (Wen et al. 2012). Starting from an initial sample of 322 clusters corresponding to the union of these four catalogues, we finally retained 258 clusters after removing clusters that are excluded by Q1 masks or that appear multiple times in different catalogues. We summarise in Table 1 the resulting cluster sample and show its properties in Fig. 1. On the one hand, we estimated the cluster masses from their richness for the DES-Y1 and WHL-SDSS cluster catalogues, following the relation of McClintock et al. (2019) for DES-Y1, and that of Wen et al. (2012) for WHL-SDSS. On the other hand, for the clusters from eROSITA and MCXC, we used the mass values in their catalogues. For each cluster, we estimated the cluster radius R_{500c} , which is defined as the radius of a sphere that encloses a mass M_{500c} with an average density equal to 500 times the critical density of the Universe at the cluster redshift.

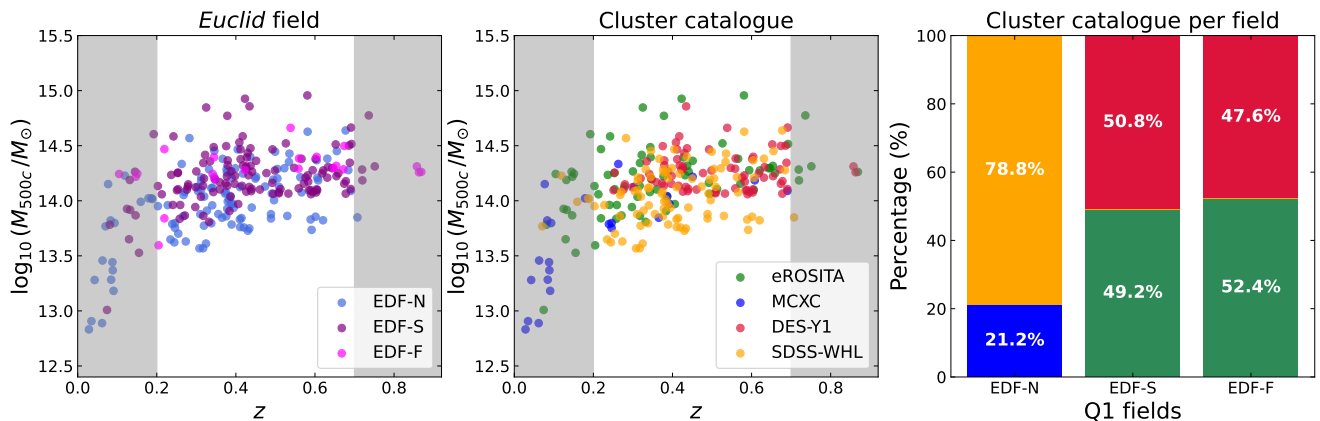
In the left and middle panels of Fig. 1, we show the cluster distribution in mass- z space, colour-coded according to their location in the three Q1 patches (left panel), and in their initial catalogue (middle panel). As shown in the right panel of Fig. 1, this cluster sample can be divided into two distinct subsamples, one containing clusters from MCXC and WHL-SDSS located in EDF-N and the second containing eROSITA and DES-Y1 clusters located in EDF-S and EDF-F. We further excluded clusters at redshifts lower than $z = 0.2$, and higher than $z = 0.7$, to avoid bias due to incomplete sampling of the cluster catalogues. Following this redshift selection, our final cluster catalogue consisted of 219 objects.

2.3. Cosmic web detection

We discuss here our choice of galaxy selection and the construction of the 2D slices centred on each cluster's redshift performed so that we could extract, from the *Euclid* Q1 data, the most robust 2D cosmic web skeletons around clusters. We also discuss

Table 1. Summary of the selected clusters (see text for details).

Catalogues	eROSITA	MCXC	DES-Y1	SDSS	Total	Double	Masked	Selected	$0.2 < z < 0.7$
Clusters in Q1	76	29	123	94	322	-45	-19	258	219


Fig. 1. Selected cluster distribution in mass- z space. The clusters are colour-coded according to the *Euclid* field they fall in and by their native catalogues.

the different filament-finder techniques applied on the 2D slices to measure cluster connectivities and their parametrisation. The methodology described here is the same as that presented in [Euclid Collaboration: Laigle et al. \(2025\)](#). It optimises cosmic web extraction from the Q1 data through Monte Carlo sampling of the galaxy PDF(z), enabling the construction of tomographic slices.

2.3.1. Galaxy selection

The photometric catalogues from Q1 were produced by the OUMER pipeline ([Euclid Collaboration: Romelli et al. 2025](#)), using images from the VIS and NISP instruments ([Euclid Collaboration: Cropper et al. 2025](#); [Euclid Collaboration: Jahnke et al. 2025](#)), processed by OU-VIS ([Euclid Collaboration: McCracken et al. 2025](#)) and OU-NIR ([Euclid Collaboration: Polenta et al. 2025](#)), along with external ground-based datasets such as the Ultraviolet Near- Infrared Optical Northern Survey (UNIONS). Morphological parameters, including Sérsic fits (Sérsic 1963), were measured with SourceExtractor++ (Bertin et al. 2020), and additional visual-like galaxy morphologies were obtained using deep-learning techniques ([Euclid Collaboration: Romelli et al. 2025](#)). For details on the morphology measurements, we refer to [Euclid Collaboration: Quilley et al. \(2025\)](#) and [Euclid Collaboration: Walmsley et al. \(2025\)](#). The redshifts and stellar masses were derived using the OUPHZ pipeline ([Euclid Collaboration: Tucci et al. 2025](#)), using two distinct methods. The first method, Phosphorus, is based on template-fitting models and provides photometric redshifts along with Bayesian posterior distributions. The second method, Nearest-Neighbour Photometric Redshifts (NNPZ), is a machine learning-based approach that computes redshifts and stellar masses based on the 30 nearest neighbours from a calibration sample. It provides redshift mode, median, and percentiles.

For the connectivity analysis and to construct redshift slices centred on cluster redshifts, we used both the Phosphorus redshift posterior distribution of galaxies as their redshift probability distribution functions, called PDF(z), and the NNPZ stellar masses. We start with the galaxy catalogue available in the Euclid Science Archive System and select the galaxies for our subsequent analysis by applying the following steps:

- Selection of mean photometric redshifts from Phosphorus such that $0 < z < 1$;
- Exclusion of artefacts applying the following quality cuts for retained sources:
 - `phz_classification` = 2 (classified as galaxies),
 - `phz_flags` = 0, `spurious_prob` < 0.1, and $M_{\star} < 10^{14} M_{\odot}$ (free of spurious detections).

2.3.2. Construction of 2D slices centred on clusters

For each cluster, we construct a 2D slice centred at the cluster redshift in order to detect the 2D cosmic web skeleton based on the galaxy distribution. To achieve this, we first determine the optimal galaxy mass selection and the slice thickness so that the most probable galaxies are included within the cluster slice taking into account the galaxy redshift uncertainties, illustrated in Fig. 2 across the three Q1 fields (displayed in separate columns) for various galaxy mass thresholds (colour-coded). The 1σ (high-opacity markers) and 2σ (low-opacity markers) redshift uncertainties are both shown. The top panels of Fig. 2 display the redshift uncertainties, while the bottom panels show their corresponding comoving lengths along the line of sight. We conclude that the mean 2σ error is of the order of $170 h^{-1}$ Mpc for galaxies more massive than $10^{10.3} M_{\odot}$. For each cluster, we decided thus to extract the cosmic web on the 2D galaxy distribution, with galaxies more massive than $10^{10.3} M_{\odot}$ and projected in slices of thickness $170 h^{-1}$ Mpc.

We note that in the northern field (EDF-N), the external data come from the Ultraviolet Near-Infrared Optical Northern Survey (UNIONS), whereas in the southern fields (EDF-F and EDF-S), the data originate from the Dark Energy Survey (DES). These differences naturally explain why the photometric redshift uncertainties are very similar between EDF-F and EDF-S, but are noticeably different in EDF-N. In particular, the absence of u -band data in the two southern fields leads to slightly larger photo- z uncertainties at $z < 0.5$. Figure 8 of [Euclid Collaboration: Ausser et al. \(2025\)](#) clearly illustrates the varying depths of these complementary datasets. Our chosen slice thickness seems small for EDF-S and EDF-F in the redshift range $0.2 < z < 0.4$. However, since we focus on cluster environments (typically less

than $5R_{500c}$), where galaxies are on average more massive, and thus are expected to have more accurate redshifts, this limitation might be mitigated. Enlarging the slice, would significantly degrade the cosmic web reconstruction for EDF-N and other redshift ranges of EDF-F and EDF-S. One possible improvement could be to adapt the thickness to the fields; however, for the particular case of $0.2 < z < 0.4$ in EDF-S and EDF-F, a thickness of 2σ corresponds to $300 h^{-1}$ Mpc, and hence is so large that it would enclosed aligned cluster systems. Following the literature, we detail the discussion on our choice of 2D slice thickness in the Appendix A. To account for the photometric redshift uncertainties on the population of galaxies belonging to a slice, we performed 100 realisations of each slice by randomly sampling redshifts from the PDF(z) of each galaxy. The density of galaxies used to trace the cosmic web varies slightly with redshift, but we did not apply a density cut correction in the present study. However, this will be reassessed with future *Euclid* releases (as discussed in Euclid Collaboration: Laigle et al. 2025). Moreover, because the slice thickness is fixed in comoving units, it corresponds to slightly different multiples of the photometric-redshift scatter in different fields and redshift ranges. This may lead to small field- and redshift-dependent biases that we do not attempt to correct for in this first analysis.

2.3.3. Filament finder techniques

To check the robustness of our results with respect to the filament-finder technique applied in the analysis, we used two different algorithms, `DISPERSE` and `T-REX`. The `DISPERSE` algorithm analyses the topology of the density field (Sousbie et al. 2011; Sousbie 2011). The density field is reconstructed from the discrete galaxy distribution via the Delaunay Tessellation Field Estimator (Schaap & van de Weygaert 2000), where the density is inversely proportional to the area of a triangle in the tessellation. Then, the `DISPERSE` algorithm identifies filaments as ridges topologically connecting pairs of saddle and peak critical points determined through discrete Morse theory. The persistence parameter (σ) sets the significance threshold for filament detection, effectively distinguishing meaningful structures from Poisson noise. In the context of 2D cosmic web reconstruction with photometric galaxies, Sarron et al. (2019) and Darragh Ford et al. (2019) have shown that the optimal persistence value for the `DISPERSE` algorithm to detect filaments connected to clusters ranges between $\sigma = 1.5$ and 2. Following these studies, we set the persistence at $\sigma = 1.5$ and $\sigma = 2$ to capture the large-scale cosmic filaments connected to clusters. These two runs of `DISPERSE`, over the 100 realisations of the 219 cluster slices, are discussed further to explore the impact of the persistence on the overall mass-connectivity relation.

We also utilise the `T-REX` filament finder to detect cosmic web skeletons in 2D slices (Bonnaire et al. 2020; Bonnaire et al. 2022). This complements the `DISPERSE` findings by providing an alternative method to extract the filamentary structure. The `T-REX` algorithm defines filaments as a set of smooth one-dimensional ridges, leveraging a machine learning extension of the minimal spanning tree where nodes of the graph are represented by Gaussian components of a mixture model. The spatial distribution of these nodes is optimised iteratively using the expectation-maximisation algorithm, which maximises a regularised posterior distribution to best fit the galaxy distribution, while preserving a smooth graph representation and a robustness to uniformly distributed noise in the covered area. The trade-off between accuracy and smoothness is governed by the parameter λ , which imposes an indirect constraint on the total length of the

graph during the optimisation. Graph nodes are initialised using a cut in the extremities of the minimal spanning tree, enabling a proper population of the distribution of galaxies initially. For our analysis, we set $\lambda = 5$ to capture a smooth representation of the large-scale cosmic filaments, while avoiding smaller bridges of matter between galaxies (similarly to Gouin et al. 2021, using `T-REX` to compute cluster connectivity in simulations).

To detect the cosmic web skeleton for each cluster, both filament finders are applied to the 100 realisations of their 2D slices. In each realisation, we calculate the cluster connectivity¹ κ , defined as the number of filaments crossing a circle of radius R_k . In the same range of radial distance as Darragh Ford et al. (2019) and Sarron et al. (2019), who used respectively $R_k = 1.5R_{vir}$ and 1.5 cMpc to measure 2D connectivities in COSMOS and CFHTLS, we measure here the cluster connectivity at radial distances of 2, 3, and 4 R_{500c} to investigate possible radial dependence. This yields 100 connectivity values per cluster, from which we compute the mean connectivity and its standard deviation.

For illustration, in Fig. 3 we overlay the 100 2D cosmic web skeletons identified with `DISPERSE` (in blue) and `T-REX` (in green) around three clusters extracted from the eROSITA, DES-Y1, and SDSS-WHL catalogues. The red circles centred on the clusters show the $4R_{500}$ radius environments. We observe that both filament-detection techniques agree well and output coherent skeletons. We also note that filament detection is reliable in dense cluster environments, as indicated by the strong opacity resulting from overlapping skeletons across realisations. However, filament identification becomes less reliable in underdense regions. As demonstrated in Appendix A, we recall that the 2D connectivity estimates are only used in a statistical ensemble, and are not intended for the precise characterisation of individual clusters. As illustrated in Fig. A.2, we show that 2D filaments, detected in 2D slices with $170 h^{-1}$ Mpc thickness along the line of sight, do not accurately match the 3D filaments around clusters.

In Fig. 4, we compare cluster connectivity values derived from filaments extracted with `DISPERSE` and `T-REX`, showing their one-to-one relation across Q1 fields (rows) and redshift bins (columns). In our redshift range $0.2 < z < 0.7$, the connectivities from the two algorithms agree very well for all cluster masses and Q1 fields. We note that in Appendix B, we discuss the tests of different choices of parametrisation for both `T-REX` and `DISPERSE` in order to demonstrate the stability of our results.

3. Connectivity dependence on cluster properties

3.1. Mass-connectivity relation

In Fig. 5, we present the median mass-connectivity relation measured in Q1 at $0.2 < z < 0.7$ by using the `DISPERSE` algorithm (1.5σ and $R_k = 4R_{500c}$) alongside 3D connectivity predictions from the IllustrisTNG (Gouin et al. 2021) and Horizon-AGN (Darragh Ford et al. 2019) simulations, as well as observational 2D connectivity measurements from COSMOS (Darragh Ford et al. 2019) and CFHTLS (Sarron et al. 2019). Our result provides connectivity measurements over a large mass range, $M_{500c}/M_{\odot} \in [10^{13.5}, 10^{15}]$, with a significant correlation coefficient of about 0.5 and in very good agreement with observational measurements from CFHTLS and COSMOS. We note that in this case, we converted the cluster mass given by Darragh Ford

¹ With this definition, we measure a multiplicity, and not connectivity according to Codis et al. (2018).

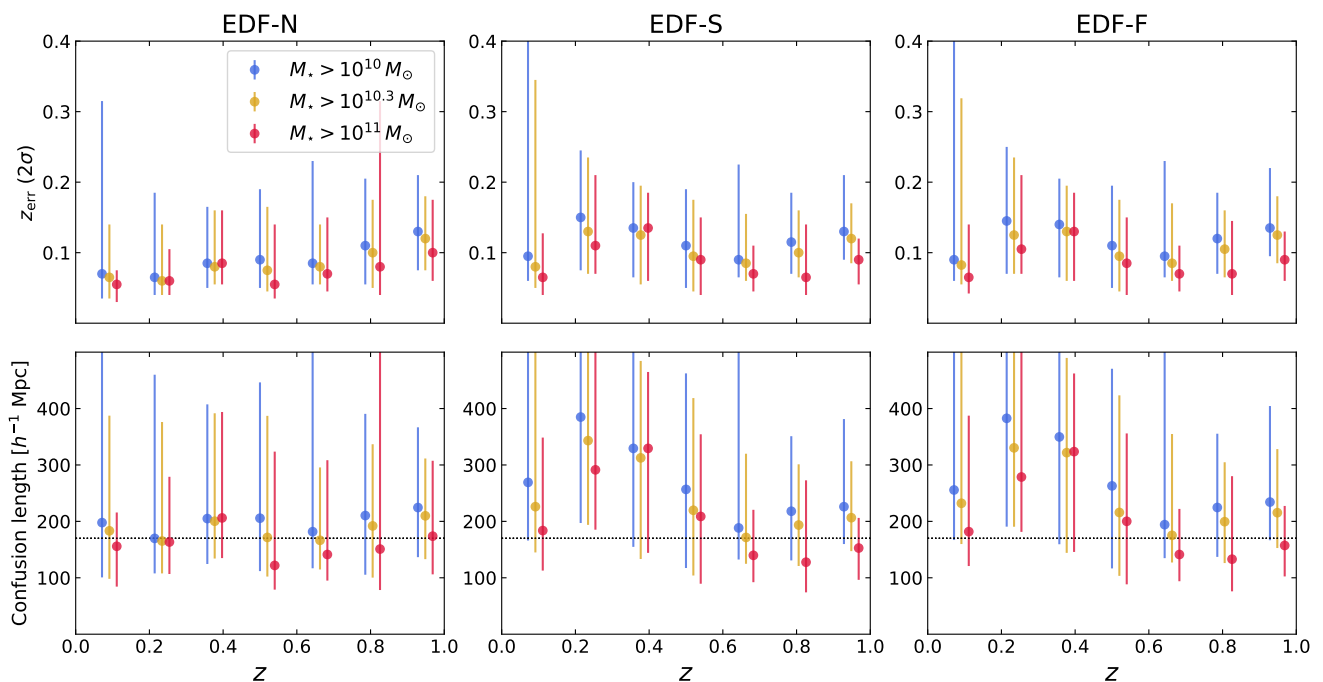


Fig. 2. *Top panels:* Median redshift error ($z_{\text{err}} = 2\sigma$) as a function of the photometric redshift. We consider here three mass selections of galaxies: $M_* > 10^{11} M_\odot$ (red), $M_* > 10^{10.3} M_\odot$ (orange), and $M_* > 10^{10} M_\odot$ (blue points). *Bottom panels:* Median confusion length, i.e. the associated errors on comoving distance as a function of redshift. The horizontal dotted line represents $170 h^{-1}$ Mpc, our choice for the thickness of redshift slices.

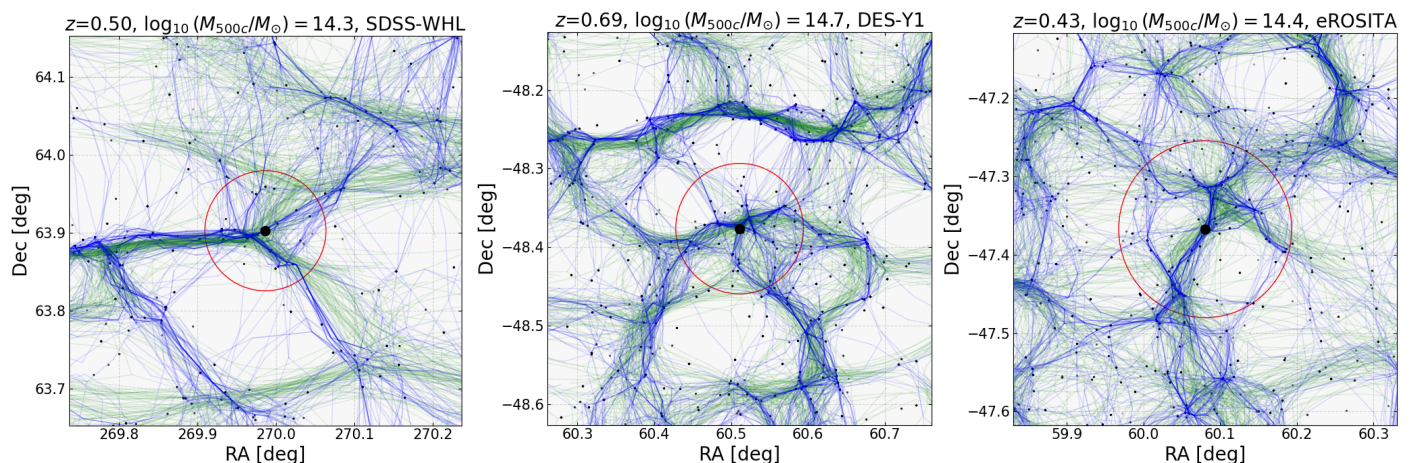


Fig. 3. Illustration of the 100 DISPERSE (blue) and T-REX (green) skeletons (from the 100 realisations) found around clusters from the SDSS-WHL (left), DES-Y1 (middle), and eROSITA (right) catalogues. The red circle is centred on each cluster with a radius of $4R_{500c}$. The patches measure $0.5 \times 0.5 \text{ deg}^2$.

et al. (2019), Sarron et al. (2019), and Guin et al. (2021), from M_{200c} to M_{500c} , by assuming that the ratio $R_{500c}/R_{200c} = 0.7$ as estimated by Ettori & Balestra (2009).

395 We recall that the absolute amplitude and the slope of the mass-connectivity relation, $M_{500c}-\kappa$, are influenced by several factors. First, a key factor is whether the connectivity is measured in 2D projected redshift slices or in full 3D space. As expected, the 2D connectivity measured in Q1 data is slightly
400 lower than the 3D connectivity predicted in simulations. This is explained by well-known projection effects: some filaments may be aligned along the line of sight or overlap with others. Sarron et al. (2019) and Darragh Ford et al. (2019), who investigated the relation between 2D and 3D connectivity measurements, show

that the 2D photometric skeleton (computed with similar slice 405 thickness) leads to an underestimation of the connectivity compared to 3D connectivities. Related to this, Laigle et al. (2018) showed that the 2D segments of filaments that have no counterpart in 3D are less robust, and thus are removed by assuming a persistence threshold. Second, in the case of 2D connectivity 410 measurements, the thickness of the redshift slices varies according to the uncertainties in photometric redshifts. Larger redshift uncertainties lead to thicker slices, which in turn affect the cosmic web reconstruction. Third, the choice of tracer used to reconstruct the cosmic web will affect the amplitude of $M_{500c}-\kappa$ 415 relation. In simulations, cosmic web reconstruction can be performed accurately by using the dark matter or gas density field

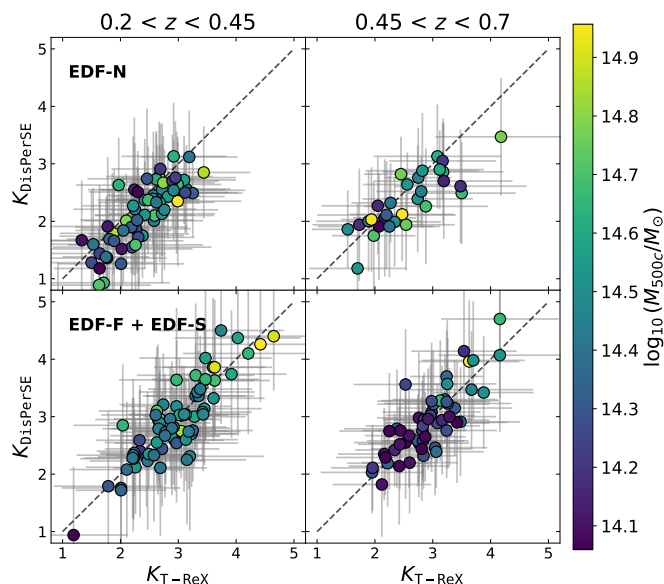


Fig. 4. DISPERSE-T-REX connectivity relation across the Q1 fields. The EDF-N is displayed in the first row and EDF-S+EDF-F in the second row, while the cluster redshifts are presented in the columns. The connectivity points are colour-coded by cluster mass.

(providing the detection of small-scale filaments); in observations, filaments are typically detected on the galaxy distributions with a selection cut, such as a stellar mass threshold. Fourth, the applied filament finder algorithm, its parametrisation, and the radial aperture considered for measuring connectivity might affect the slope of the relation. In Fig. 6 we present the M_{500c} - κ relation in Q1 at $0.2 < z < 0.7$ using DISPERSE and T-REX by considering different radial apertures (R_k) for measuring connectivity and two different DISPERSE parametrisations ($\sigma = 1.5$ and 2). On the one hand, by increasing the persistence value, we increase the density contrast threshold at which filaments are detected, and thus we reduce the number of detected filaments and cluster connectivity. Conversely, a larger aperture increases the connectivity because it enclose a larger number of filament bifurcations compared to connectivity measurements close to density peaks (as shown by Codis et al. 2018). These listed dependences explain the differences between the M_{500c} - κ relations from the literature in Fig. 5. In addition, the scatter of M_{500c} - κ itself is expected to reflect the diversity of cluster mass assembly histories (Cadiou et al. 2020; Gouin et al. 2021). Therefore, rather than focusing on the absolute amplitude of the M_{500c} - κ relation, we examine in the next section how this relation depends on the physical properties of clusters. This approach provides a more robust means of investigating cluster evolution.

3.2. Estimation of galaxy members in clusters

To estimate galaxy members within a given cluster, we refined our selection criteria described in Sect. 2.3.1. In addition to the initial cuts applied to exclude spurious sources, we identified galaxy members based on two key conditions. First, galaxies must reside within the cluster projected sky area, specifically within a projected distance of $2R_{500c}$ from the cluster centre. Second, galaxies must have a high probability of being near the cluster redshift. Therefore, we computed for each galaxy the probability that it is at the cluster redshift using the general formalism of George et al. (2011) (Castignani & Benoist 2016; Sarron

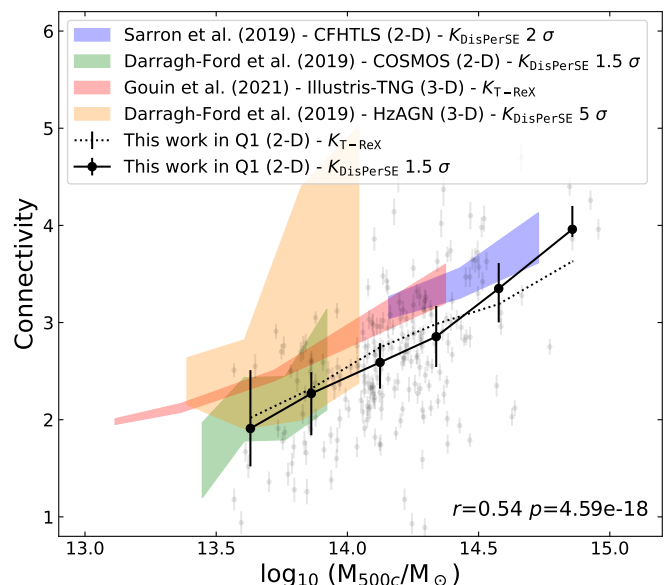


Fig. 5. Median mass–connectivity relation measured in Q1 at $0.2 < z < 0.7$ (in black) compared with results obtained, from the Horizon-AGN simulation (in pink; Darragh Ford et al. 2019) and the IllustrisTNG (in red; Gouin et al. 2021), and with observational results, from CFHTLS (Sarron et al. 2019) and COSMOS (Darragh Ford et al. 2019). The sigma values in the legend refer to the persistence threshold applied with the DISPERSE algorithm. The Pearson correlation between cluster connectivity and its mass is given at the bottom of the panel, with the correlation coefficient (r) and the p -value.

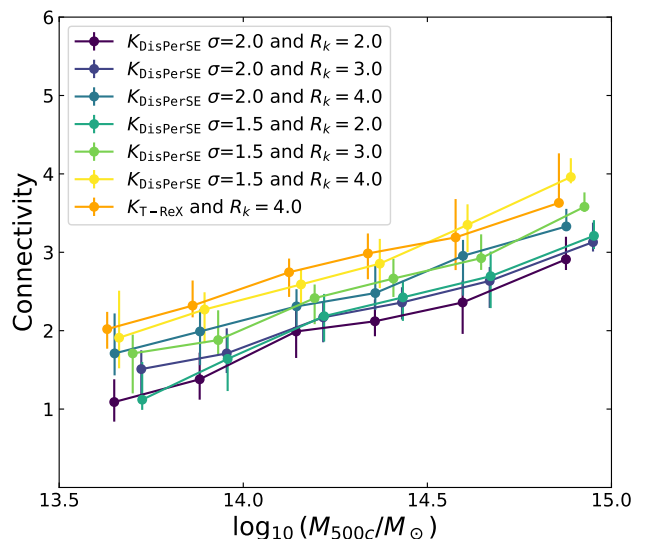


Fig. 6. Median mass–connectivity relation measured in Q1 at $0.2 < z < 0.7$, by considering two different DISPERSE skeletons with a persistence $\sigma = 1.5$ and 2, and with three different radial apertures for measuring connectivity $R_k = 2, 3$, and 4.

& Conselice 2021). In this formalism we model the expected redshift probability distribution of cluster galaxies using a normal distribution, $\mathcal{N}(z|z_c, \sigma_p)$, centred at the cluster redshift, z_c , and with standard deviation on galaxy photometric redshifts σ_p . We note that this method neglects the uncertainty in the cluster redshift itself, which is appropriate for our data as the typical redshift uncertainty for spectroscopically confirmed clusters is $\sim 0.001(1+z)$, an order of magnitude smaller than σ_p (see Sarron & Conselice 2021, for a discussion). Hence, the likelihood

of observing a cluster galaxy with $P(z)$ given this model is

$$p(P(z)|\text{gal} \in C) = \int P(z)\mathcal{N}(z|z_c, \sigma_P) dz. \quad (1)$$

Using Bayes' theorem, we can write the probability that a galaxy belongs to the cluster given its $P(z)$:

$$p(\text{gal} \in C|P(z)) \propto p(P(z)|\text{gal} \in C) p(\text{gal} \in C). \quad (2)$$

465 We consider uninformative priors $p(\text{gal} \in C) = 1$, meaning that in practice we compute the relative probability that the galaxy is at the cluster redshift, assuming the model described for cluster redshift distributions. Following the arguments in [Castignani & Benoist \(2016\)](#), this probability is rescaled such that the maximum achievable probability is one. This is done as in [Sarron & Conselice \(2021\)](#),

$$p(\text{gal} \in C|P(z)) = \frac{p(P(z)|\text{gal} \in C) p(\text{gal} \in C)}{p(P(z)|\text{gal} \in C, \sigma_P = 0.01)}, \quad (3)$$

such that a galaxy with photometric redshift distribution $P(z) = \mathcal{N}(z|z_c, \sigma_P = 0.01)$ has a probability of one. We assume that galaxies are identified as cluster members when their probability is higher than 0.5. This relative probability threshold is used as a ranking criterion and does not correspond to a fully calibrated Bayesian membership probability; in particular, it should not be interpreted as a literal 50% probability of cluster membership. For this initial exploration, we neglected any dependence on magnitudes and radius, such as cluster profiles and the segregation of bright galaxies in the cores. In future more in-depth analyses, we will use the cluster probability memberships computed by the RICH-CL processing function from the *Euclid* LE3 official galaxy cluster detection and characterisation pipeline, which improves on these limitations ([Castignani & Benoist 2016](#)).

3.3. Galaxy morphology estimation

For each cluster, the galaxy members are identified following the procedure described in Sect. 3.2. We further apply a mass selection such that $M_\star > 10^{10.3} M_\odot$, to be similar to the galaxy selection used to trace the cosmic web skeleton (see Sect. 2.3). According to [Euclid Collaboration: Quilley et al. \(2025\)](#), this stellar mass selection should not be affected by mass incompleteness, given that a much more restricted sample with $I_E < 23$ is 90% complete above $M_\star > 10^{10} M_\odot$ at $z = 0.6$. Moreover, we explored the galaxy bi-modality by plotting galaxies according to their location in the $n_{\text{Sersic}}-M_\star$ diagram in Fig. 7 for the three redshift bins from 0.2 to 0.7. As shown in this figure, there is a galaxy morphology bi-modality such that early-type galaxies (ETG) are defined by $n_{\text{Sersic}} > 1.75$ and late-type galaxies are represented with $n_{\text{Sersic}} < 1.75$. We note that the n_{Sersic} threshold slightly evolves with redshift, but we ensure that a fixed threshold did not significantly affect our result. In general, we found that the fraction of early-type galaxies in clusters is overall consistent with those reported in the literature (e.g. [Simard et al. 2009](#)), and is slightly evolving with redshift. Therefore, we later investigated the environmental impact on galaxy member morphologies as depending on both cluster mass and redshift.

3.4. Relation between connectivity and cluster galaxy morphologies

510 By using these methods to characterise cluster members and galaxy morphology, we now explore the relation between cluster

connectivity and the morphology of galaxies inside clusters. In Fig. 8 we present the $M_{500c} - \kappa$ relation colour-coded by f_{ETG} , the fraction of ETG inside clusters. Our cluster sample is divided into three different redshift bins: $0.2 < z < 0.4$ (101 clusters; left panel), $0.4 < z < 0.6$ (82 clusters; middle panel), and $0.6 < z < 0.7$ (35 clusters; right panel). We note that, for consistency, we used the same radial aperture of $2R_{500c}$ to identify cluster galaxy members and compute cluster connectivity. In addition, to quantify the correlation between the galaxy morphologies inside clusters and the connectivity, beyond mass-driven effects, we used the partial Pearson correlation which measures the degree of association between these two variables, after removing the effect of one (here the cluster mass). We found that f_{ETG} tends to correlate with connectivity, only for clusters at $0.2 < z < 0.4$, with a moderate partial correlation coefficient of 0.34 and a low p -value. Our result tends to suggest that, beyond the first-order mass dependence, the more a cluster is connected, the more it is populated by early-type galaxies. Similarly, we present in Fig. 9 the connectivity-mass relation colour-coded by the median Sérsic index $\langle n_{\text{Sersic}} \rangle$ of cluster galaxy members. We can see that, on average, clusters populated by higher Sérsic index tend to present a higher connectivity. This trend is also only weakly significant for the first redshift bin, with a correlation factor of 0.28 between connectivity and median Sérsic index. In Appendix B, we verify that this trend is robust for different DISPERSE persistence threshold, and we confirm its consistency when using the T-REX algorithm with different values of λ . This appendix therefore demonstrates the stability of our results with respect to variations in filament detection settings.

3.5. Interpretation and discussion

Our results appear to be consistent with a scenario in which high connectivity is associated with clusters predominantly populated by early-type galaxies. Supporting this result, [Darragh Ford et al. \(2019\)](#) found that galaxy groups with a passive central galaxy tend to have higher connectivity on average than those with a star-forming central galaxy in COSMOS observations (see also [Einasto et al. 2014](#)). Analysing hydrodynamical simulations, they suggested that different connectivity levels might trace distinct mass assembly histories, with highly connected groups and clusters having typically undergone their last major merger more recently. Such past merging activity could, in turn, contribute to quenching and morphological transformations of the central galaxy (and its members). Conversely, it may be easier to preserve late-type morphologies in haloes that have not merged. Additionally, [Kraljic et al. \(2020\)](#) found that less star-forming and less rotation-supported galaxies in SDSS tend to exhibit higher galaxy connectivity, a result further supported by simulations. In highly connected clusters, galaxies are subject to intensified environmental effects, including tidal and ram-pressure stripping, harassment, and strangulation ([Moore et al. 1996](#); [Gay et al. 2010](#); [Mastropietro et al. 2005](#); [Wetzel et al. 2013](#)). These mechanisms inhibit star formation, leading to galaxy quiescence and a higher proportion of elliptical galaxies. High-connectivity clusters also tend to be dynamically unrelaxed ([Gouin et al. 2021](#)), i.e. with higher velocity dispersions, and hence stronger environmental quenching, which leads to the secular disruption of disk structures (see [Hong et al. 2024](#), Appendix D, and reference therein). The multiple infalling directions around highly connected clusters may reinforce such mechanisms. From a theoretical point of view, [Aragon Calvo et al. \(2019\)](#) also proposed the Cosmic Web Detachment model, suggesting that as galaxies accrete into filaments, shell-crossing occurs ([Laille et al.](#)

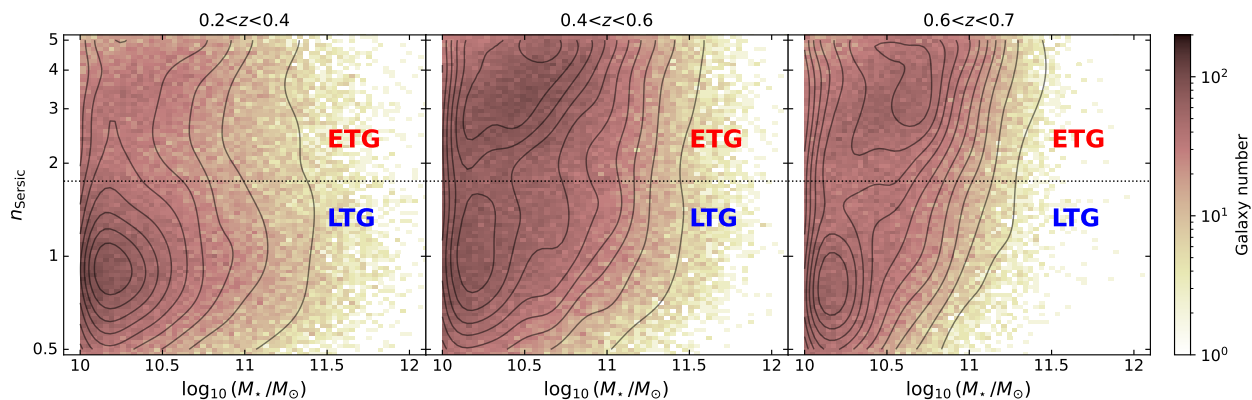


Fig. 7. Sérsic index–stellar mass diagram for galaxies in three different redshift bins: $0.2 < z < 0.4$ (left panel), $0.4 < z < 0.6$ (middle panel), and $0.6 < z < 0.7$ (right panel). We divided the galaxies into two types: early-type galaxies with $n_{\text{Sérsic}} > 1.75$ and late-type galaxies with $n_{\text{Sérsic}} < 1.75$.

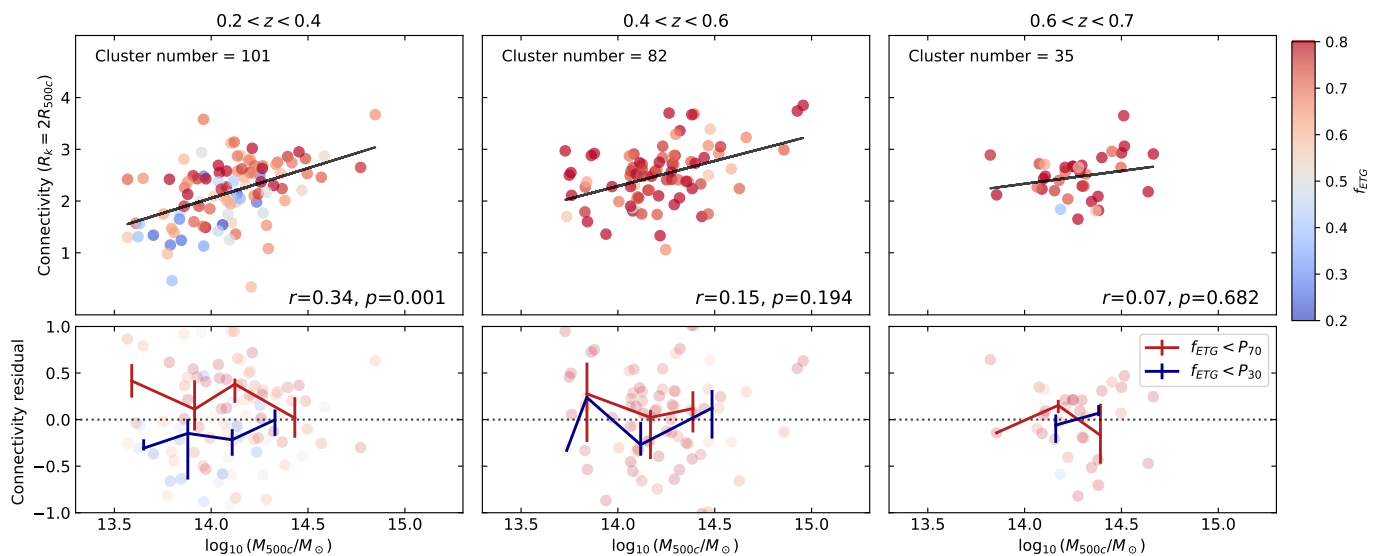


Fig. 8. Mass–connectivity relation from Q1 data by considering clusters in three different redshift bins: $0.2 < z < 0.4$ (left panel), $0.4 < z < 0.6$ (middle panel), and $0.6 < z < 0.7$ (right panel). The points are colour-coded by f_{ETG} fraction of early-type galaxies inside clusters ($R < 2R_{500c}$). The black solid lines show the connectivity–mass relation on average, with the connectivities measured at $R = 2R_{500c}$. In the bottom panel we show the connectivity residual, defined as $\kappa - \langle \kappa(M_{500c}) \rangle$ to remove mass dependence. The red (and blue) solid lines represent the average profile for clusters with f_{ETG} lower than (higher than) the 30th(70th) percentile. The partial Pearson correlation between cluster connectivity and ETG fraction, beyond first-order mass dependence, is given at the bottom of the panels, with the correlation coefficient (r) and the p -value.

2015), cutting off their cold gas supply and ultimately quenching star formation. This model might explain their possible pre-processing, even before they enter clusters (as observed by Conzelice et al. 2001; Sarron et al. 2019; Gouin et al. 2020). Both pre-processing and processing could explain our findings, weak but statistically significant, that the high fraction of ETGs in clusters is correlated with high connectivity values.

This scenario appears to contrast with several past observational studies that reported enhanced star formation in clusters (see e.g. Porter & Raychaudhury 2007; Fadda et al. 2008; Biviano et al. 2011; Darvish et al. 2014; Lee et al. 2019; Ko et al. 2024). However, a number of factors can naturally explain these differences. First, the redshift range probed in some of these works is significantly higher. For example, Lee et al. (2019) analysed eight groups and clusters at $0.6 < z < 1.3$, i.e. systems in an early assembly phase where cold gas accretion through filaments is still efficient, potentially triggering starbursts during the proto-cluster stage. A similar effect may apply to the $z = 1$ sample of Darvish et al. (2014). Second, the nature of the structures examined varies. Porter & Raychaudhury (2007) focused

on filaments within the Pisces–Cetus supercluster, whereas our study investigates the statistical behaviour of a large population of clusters rather than individual superstructures. Along the same line, Fadda et al. (2008) found that the fraction of starburst galaxies in filaments is more than twice that in the cores of Abell 1770 and Abell 1763. Finally, none of these examples directly quantify connectivity; instead, they rely on proxies for the large-scale environment, such as the friend-of-friend fraction used in Ko et al. (2024).

Taken together, these apparent discrepancies likely reflect differences in redshift, structure type, analysis scale, and galaxy populations considered. Rather than being contradictory, they point towards a more complex evolutionary picture in which (i) proto-clusters may experience an early phase of enhanced star formation, and (ii) supercluster environments or recently merged structures may temporarily boost star formation before quenching becomes dominant. Overall, we argue that the global framework for galaxy and cluster evolution within the large-scale structure is still under development, and that additional simulations and observations will be required to establish a unified scenario.

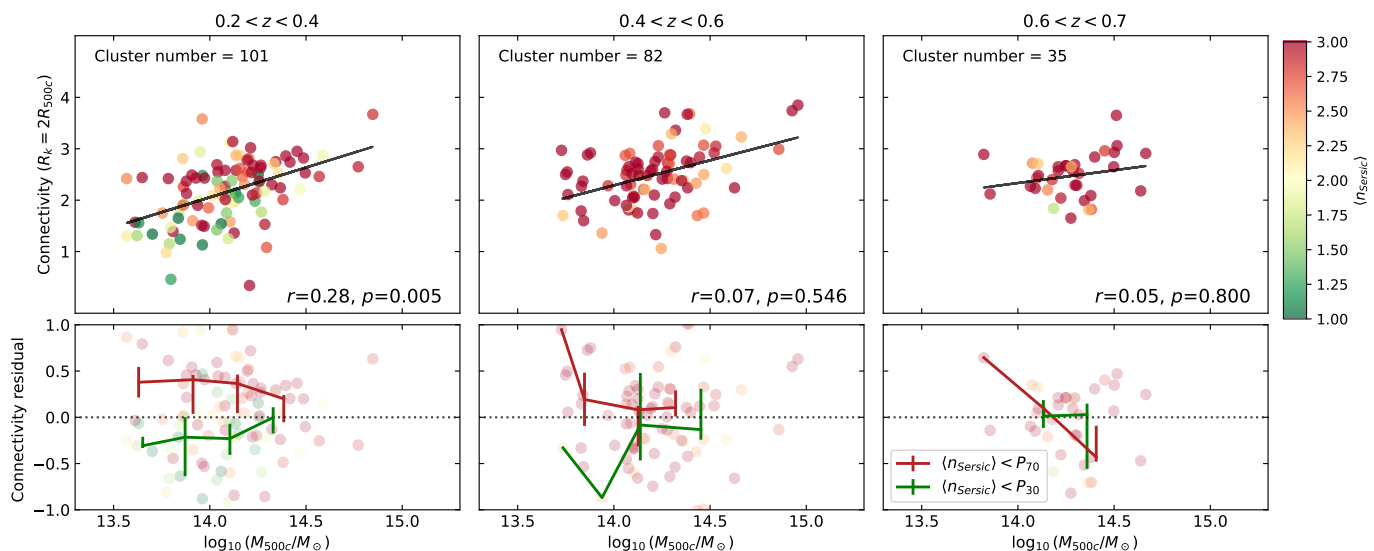


Fig. 9. Same as Fig. 8, but colour-coded by the median Sérsic index of galaxies inside clusters $\langle n_{\text{Sérsic}} \rangle$.

It should be noted that in this proof of concept investigation we have investigated the dependence of the $M_{500c}-\kappa$ relation on the morphology of galaxy members. This study serves as a first step towards a more comprehensive analysis; we will later extend the approach to other galaxy properties such as star formation rate and the fraction of quenched galaxies in clusters. A relation between cluster member morphology and star formation activity has been established in the literature, showing a clear trend in which more massive haloes host a larger fraction of quenched galaxies (e.g. Paccagnella et al. 2016; Reeves et al. 2021). However, at this first stage, we focused solely on morphology, as Q1 data provides accurate morphological measurements.

4. Conclusions

This study investigated the role of the cosmic web in shaping galaxy clusters using the first *Euclid* Quick Release 1 data. For this work we used an ensemble of 219 clusters at $0.2 < z < 0.7$ from the eROSITA (Bulbul et al. 2024), MCXC (Piffaretti et al. 2011), DES-Y1 (Rykoff et al. 2016), and SDSS (Wen et al. 2012) catalogues. By using the photometric redshift posterior distributions of galaxies provided by Q1, we performed a Monte Carlo resampling of the galaxy PDF to build 100 realisations of each 2D slice centred on each cluster. This statistical procedure allowed us to accurately estimate the connectivity for each cluster. By using two different filament-finder algorithms (T-REX and DISPERSE), we ensured the robustness of our connectivity measurements. Even if these 2D connectivity measurements are not suitable for the precise characterisation of the cosmic web environments of individual clusters, we demonstrate in Appendix A that our 2D methodology reasonably recovers statistical information on the connectivity of the large-scale structure in dense environments.

We confirmed the expected mass-connectivity relation predicted by hierarchical structure formation models (Codis et al. 2018). Our result provides 2D connectivity measurements over a wide mass range, and is found to be in very good agreement with past observational measurements (Darragh Ford et al. 2019; Sarson et al. 2019). Moreover, we explored the relation between the connectivity and morphology of galaxy members. By accurately identifying galaxy members of clusters in Q1 data, we found

a moderate correlation, suggesting that the higher the fraction of early-type galaxies, the higher the average connectivity, but only for low-redshift clusters ($0.2 > z > 0.4$). Finally, investigating the median Sérsic index of galaxy members, we found a weak correlation, indicating that for low-redshift clusters, that the higher the median Sérsic index of galaxies, the higher the average connectivity. These weak but statistically significant findings, are consistent with a scenario in which high cluster connectivity is associated with clusters predominantly populated by elliptical galaxies. These results are in agreement with the trend found by Darragh Ford et al. (2019) on the impact of connectivity on the star formation activity of group central galaxies in COSMOS, and with the results from Kraljic et al. (2020) on relations between galaxy connectivities and their properties.

This work demonstrates the capabilities of Q1 data to investigate the impact of the cosmic web's filaments on cluster evolution. The results pave the way for more comprehensive analyses with future *Euclid* data releases, including higher redshift ranges and deeper spectroscopic datasets. At the end of the *Euclid* mission, the EDF will have been visited 40 times and will provide a novel spectroscopic sample, including galaxies with an $H\alpha$ flux above $5 \times 10^{-17} \text{ erg cm}^{-2} \text{ s}^{-1}$ with 60% completeness. In parallel, using the *Euclid* spectroscopic sample will allow us to reduce the slice thickness (to $25 h^{-1} \text{ Mpc}$ comoving), and to investigate the galaxy cluster accretion properties in greater detail. In a future study, we will extend the present proof of concept analysis using the catalogue of clusters detected in the Q1 data with their identified galaxy members (*Euclid* Collaboration: Bhargava et al., in prep.) and later the DR1 *Euclid* cluster catalogue. *Euclid* will provide galaxy cluster samples identified using the AMICO (Bellagamba et al. 2018; Maturi et al. 2019) and PZWav (Werner et al. 2023; Thongkham et al. 2024) algorithms, extending the overall observed area and reaching higher redshifts, up to $z \sim 2.0$, containing hundreds of thousands of sources (Sartoris et al. 2016; *Euclid* Collaboration: Adam et al. 2019).

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Appendix A: Discussion on the slice thickness

The slice thickness results from balancing two competing needs: limiting confusion from stacking filaments of different intrinsic 3D scales, and accommodating photometric redshift uncertainties. This choice is consistent with previous studies that examined projection effects on the 3D filamentary skeleton (Tab. A.1). For instance, Sarron et al. (2019) used $300 h^{-1}$, Mpc slices and Darragh Ford et al. (2019) used $120 h^{-1}$, Mpc, both with photometric redshift uncertainties comparable to ours. Using mock catalogues, they showed that although 2D projections lower the absolute connectivity amplitude, the trends with galaxy properties (Sarron et al. 2019) and with BCG–connectivity correlations (Darragh Ford et al. 2019) remain robust and physically consistent with their 3D counterparts.

In addition to previous studies, we test the robustness of our 2D-slice methodology using the FLAGSHIP simulation (Carretero et al. 2017; Tallada et al. 2020). Specifically, we use galaxies with DR1-like photometry from FLAGSHIP and apply the same pipeline to a 15×15 , deg² mock of the South field from Euclid Collaboration et al. (2025). We note that we have performed the same test on the North field mock, but here we show the worst-case scenario, i.e. the South field with higher photo- z uncertainties, to test the cosmic web reconstruction methodology. As shown in Fig. A.1, the resulting mock galaxy sample reproduces the photometric redshift uncertainties of the Q1 Euclid Deep Field South data presented in Fig. 2. Following Sect. 2, we adopt the same galaxy selection ($M_{\star} > 10^{10.3} M_{\odot}$ and $0.1 < z < 0.8$) to construct mock 2D slices of $170 h^{-1}$, Mpc, obtained via Monte Carlo resampling of the galaxy photo- z PDFs.

To further validate this method, we also construct a 3D map around each cluster using the true redshifts of the mock galaxies. This allows us to obtain the corresponding true 3D skeletons within the same 2D slices. The 3D skeletons are computed 100 times to estimate the intrinsic error on the 3D connectivity: for each iteration, we bootstrap 99% of the true galaxy distribution and compute the 3D skeleton and connectivity. The mock study contains 1380 clusters, detected with the ROCKSTAR halo finder (Behroozi et al. 2013), which provides halo masses M_{halo} and radii R_{halo} (dark matter halo properties later used for painting galaxies through abundance-matching techniques). Here, we assume galaxy clusters to be haloes with masses $M_{\text{halo}} > 10^{14} M_{\odot} h^{-1}$. We note that, in this mock case, the 3D (or 2D) connectivity is defined as the number of filaments intersecting a sphere (or circle) of radius R_{halo} .

In Fig. A.2, we show two examples of mock clusters, with the 2D (blue) and 3D (red) skeletons overlaid on the galaxy distribution around the clusters. True galaxies spatially close to the cluster along the line of sight are marked with red circles. Qualitatively, while the 2D and 3D skeletons appear quite different, they tend to indicate similar directions around clusters.

In Fig. A.3, we explore the connectivity–mass relation for both 2D and 3D skeletons. Both capture the same trend with halo mass, suggesting that 2D connectivity can statistically recover the impact of the large-scale environment on cluster properties. We note that the 3D skeleton (persistence of 3) is computed with twice the persistence of the 2D skeleton (persistence of 1.5, the same as in Q1 observations), since the 3D case is essentially noise-free and requires higher persistence to avoid generating filaments at overly small scales (similarly to Darragh Ford et al. 2019; Sarron et al. 2019, for a 2D and 3D DISPERSÉ comparison).

In Fig. A.4, we examine the correlation between 2D and 3D connectivity and find a Pearson correlation coefficient of 0.28.

The substantial scatter highlights the large uncertainties associated with estimating connectivity in 2D. Nevertheless, a positive correlation is still present. To complement this analysis, Fig. A.5 shows the probability distribution function of the 2D connectivity for different ranges of 3D connectivity. We observe that the peak of each PDF follows the expected ranking of 3D connectivity, although 2D connectivity is generally lower than its 3D counterpart (as expected from Sarron et al. 2019, Fig. 4).

This Appendix analysis demonstrates that our 2D methodology reasonably recovers statistical information on the connectivity of the large-scale structure in dense environments, even though the presence of photometric redshift uncertainties strongly affects the reconstruction of the cosmic web around individual clusters.

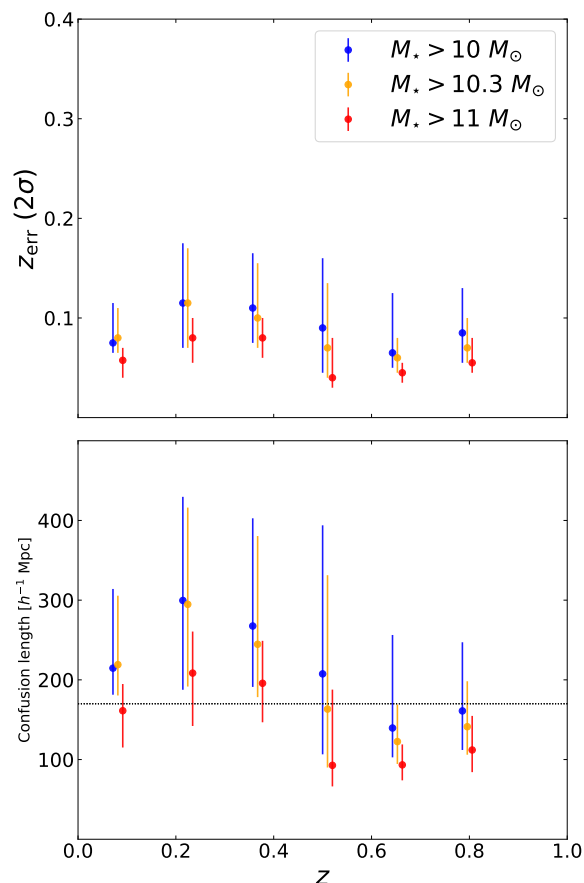


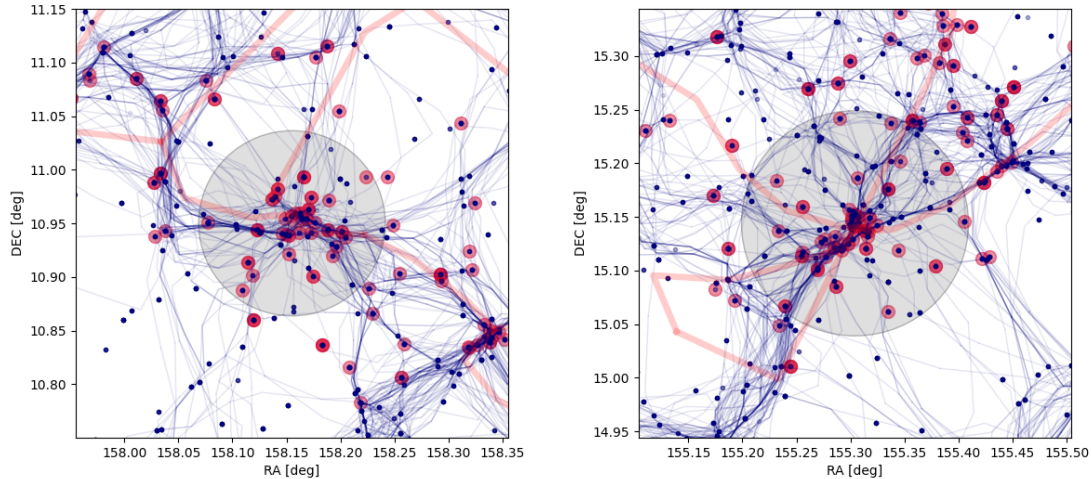
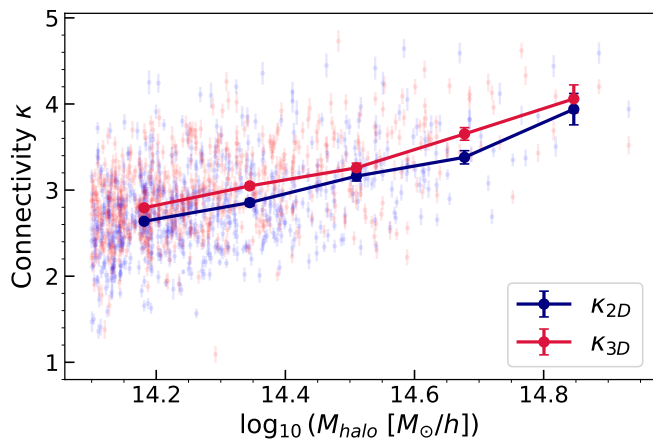
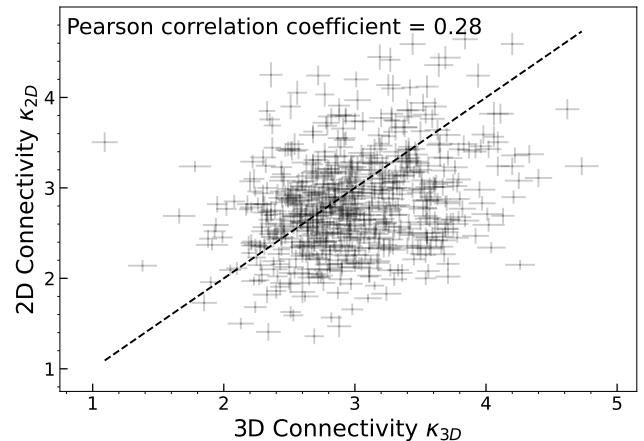
Fig. A.1. Same as Fig. 2, but considering mock galaxies from the FLAGSHIP simulation (using DR1-like photometry).

Appendix B: Filament finder parametrisation

In this Appendix, we discuss the robustness of our results with respect to the parametrisation of both T-REX and DISPERSÉ. Our goal is to verify that the observed trends in cluster connectivity and their correlation with galaxy morphology are not driven by specific choices of parameters in the filament detection algorithms. We focus on the regime where the signal is statistically significant, namely low-redshift clusters ($0.2 < z < 0.4$), in order to assess the impact of cluster connectivity on the morphology of their member galaxies. Cluster member galaxies are identified within $2R_{500c}$, and connectivity is measured at $R_k = 2R_{500c}$, ensuring a consistent definition of the cluster environment across the sample.

Table A.1. Table of the main studies comparing 2D and 3D connectivity measurements, with predictions applied to mock observations.

Citations	Observations	Simulations	2D thickness
Laigle et al. (2018)	COSMOS2015	HorizonAGN	75 $\text{Mpc}.h^{-1}$
Darragh Ford et al. (2019)	COSMOS2015	HorizonAGN	120 $\text{Mpc}.h^{-1}$
Sarron et al. (2019)	CFHTLS	Deep EUCLID lightcone	300 $\text{Mpc}.h^{-1}$
Euclid Collaboration: Malavasi et al. (2025)	-	GAEA model on Millennium	75 $\text{Mpc}.h^{-1}$
EC: Sarron et al. (<i>in prep</i>)	-	GAEA model on Millennium	25 $\text{Mpc}.h^{-1}$


Fig. A.2. Visualisation of 100 2D skeleton realisations (dark blue lines). Also shown are 3D true skeletons (red lines; projection for 3D filaments between $\pm 3R_{\text{halo}}$). The same galaxies in the 100 realisations of the 2D slices are blue points and the galaxies that are actually enclosed in $\pm 3R_{\text{halo}}$ along the line of sight are circled in red. The black circles are centred on the cluster halo, with radius equals R_{halo} .

Fig. A.3. Mass dependence of 2D (red) and 3D (blue) connectivities.

Fig. A.4. One-to-one relation between 2D and 3D connectivities.

1355 In Fig. B.1, we examine the connectivity–mass relation,
colour-coded by the ETG fraction, using the DISPERSE con-
nectivity for three different persistence thresholds, σ : 0.5 (right
panel), 1.5 (middle panel), and 2.5 (left panel). We find that
1360 skeletons with $\sigma = 0.5$ generally produce noisy filaments, which
slightly weakens the observed correlation. In contrast, a persis-
tence threshold of 2.5 yields highly robust filaments but reduces
the number of detected structures, thereby narrowing the range of
connectivity. Overall, the trend is observed across all persis-
tence values, with $\sigma = 1.5$ providing an optimal balance between
1365 noise suppression and filament detection.

A similar test is performed in Fig. B.2 for T-REX, where we
vary the parameter λ , which controls the trade-off between accu-
racy and smoothness in the filament reconstruction. We adopt
three values of λ (1, 5, and 10), spanning from high sensitiv-
ity to more strongly smoothed skeletons. We find that the level
1370 of smoothing applied during the filament detection process has
a minimal impact on the observed trends between connectivity
and ETG fraction.

Overall, these tests demonstrate that our results are stable
across different parameter choices for both T-REX and 1375
DISPERSE. These findings support the robustness of our con-

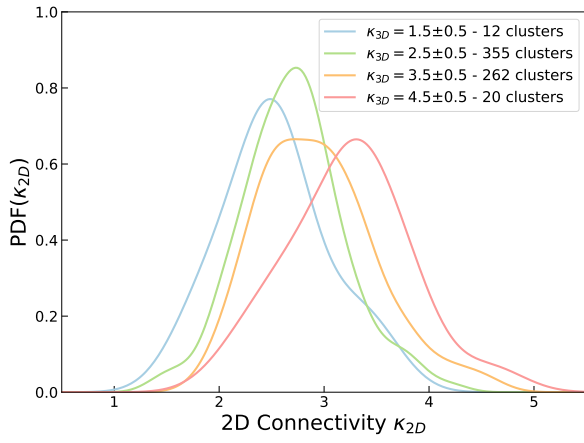


Fig. A.5. Probability distribution function of 2D connectivity κ_{2D} for four ranges of 3D connectivity κ_{3D} : $\kappa_{3D} = [1 - 2]$ (blue), $\kappa_{3D} = [2 - 3]$ (green), $\kappa_{3D} = [3 - 4]$ (orange), and $\kappa_{3D} = [4 - 5]$ (red).

clusions, with no strong dependence on the filament detection parameters.

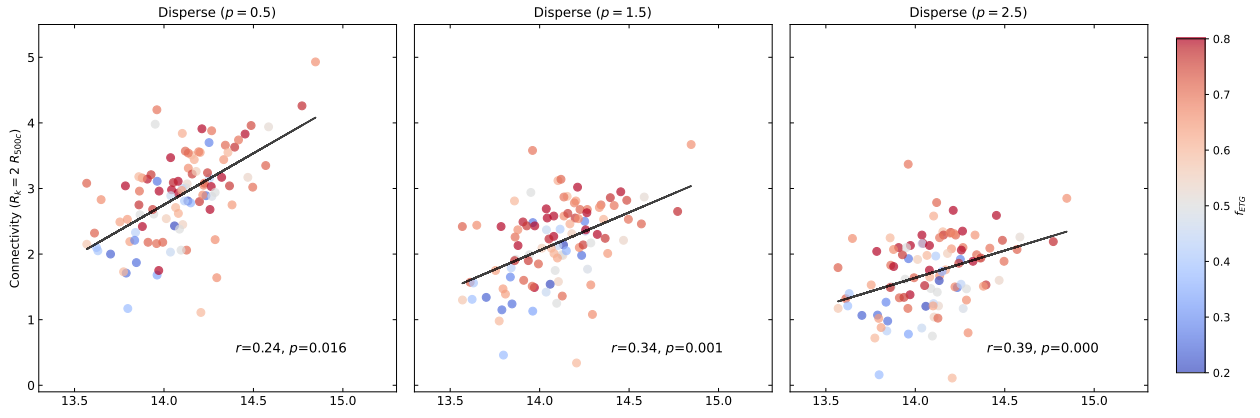


Fig. B.1. Same as Fig. 8, but considering only clusters at $0.2 > z > 0.4$, and testing different values of σ persistence on our connectivity computed by the DISPERSE algorithm.

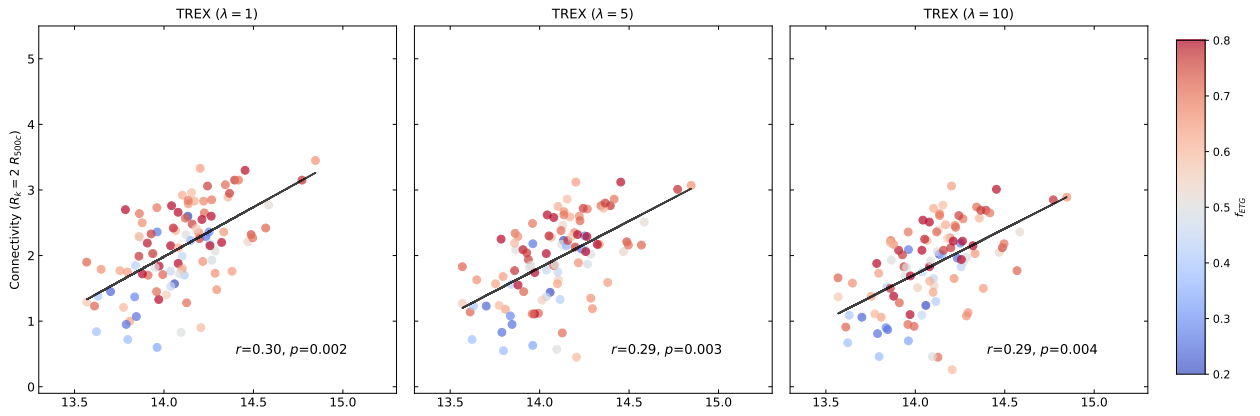


Fig. B.2. Same as Fig. 8, but considering only clusters at $0.2 > z > 0.4$, and testing different values of λ parameter on our connectivity computed by the T-REX algorithm.