1	Experimental evidence for lava-like mud flows under Martian surface conditions							
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26 Large outflow channels on ancient terrains of Mars have been interpreted as the products of catastrophic flood events. The rapid burial of water-rich sediments 27 following such flooding could have led to sedimentary volcanism, in which mixtures of 28 29 sediment and water (mud) erupt to the surface. Tens of thousands of volcano-like landforms populate the northern lowlands and other local sedimentary depocenters on 30 Mars. However, it is difficult to determine whether the edifices are related to igneous 31 or mud extrusions, partly because the behaviour of extruded mud under martian 32 surface conditions is poorly constrained. Here, we investigate the mechanisms of mud 33 34 propagation on Mars using experiments performed inside a low-pressure chamber at cold temperatures. We find that low viscosity mud under martian conditions 35 propagates differently from on Earth, because of rapid freezing and the formation of 36 37 an icy crust. Instead, the experimental mud flows propagate like terrestrial pahoehoe lava flows, with liquid mud spilling from ruptures in the frozen crust, then refreezing 38 to form a new flow lobe. We suggest that mud volcanism can explain the formation of 39 40 some lava-like flow morphologies on Mars, and that similar processes may apply to cryovolcanic extrusions on icy bodies in the Solar System. 41

42 The physics behind igneous volcanism on Mars is better understood [e.g., 1-4] than that of sedimentary volcanism in which mixtures of water and sediment, subsequently referred 43 to as mud, are extruded onto the surface. On Earth, sedimentary volcanism manifests at the 44 surface as eruptions of fluids (water, gas, occasionally oil), fine grained sediments (e.g. 45 clays) and clasts from the country-rock. These geological phenomena are the result of fluid 46 (on Earth typically associated with methane) overpressure [5], generated at several hundred 47 48 to several thousand metres depth, combined with gravitational instability of buoyant sedimentary units buried at deeper stratigraphic levels [6]. The viscosity of ascending mud 49 varies, and affects the shapes, sizes and thicknesses of resulting flows. The higher the water 50

content, the lower the viscosity and vice versa (Figure 1). The focus of our experimental
investigation are water-dominated mud flows that propagate over shallow slopes via
centimetre-thick flows (Fig. 1a,b), as opposed to clay-dominated flows that can be meter(s)
thick (Fig. 1c).

Although the propagation of water at low atmospheric pressure has been previously 55 studied [7-11], there is a lack of theoretical and empirical knowledge about the behaviour of 56 mud at low atmospheric pressure, temperature and gravity, despite an initial study by [12]. 57 This knowledge gap represents an obstacle in the study of landforms interpreted to be the 58 59 result of mud extrusion on Mars [13-24] and other terrestrial or icy solar system bodies. Currently, the low martian atmospheric pressure inhibits the sustained presence of liquid 60 water on the surface [e.g., 7-11,25,26], so evaporation and ice-formation cause the rheology 61 62 of the extruded mud to change rapidly; hence mud flows could propagate differently from on Earth [17]. 63

We used analogue experiments performed in a low pressure chamber to examine how low 64 viscosity, water-dominated mud with a solid fraction of less than ~6.5 wt.% (12.7 mPa.s at 65 276 K and 10.7 mPa.s at 296 K) propagates over a cold surface (244 K to 265 K) under 66 terrestrial and martian (7 mbar) atmospheric pressures. These experiments enabled 67 comparing flow mechanisms at different pressures to be compared and reveal a unique 68 propagation behaviours under martian conditions. Based on these observations we propose 69 that evaporation and surficial freezing would dominate the morphology for relatively thin 70 mud flows (< 1 m) and may influence thicker mud flows hypothesized to be present on the 71 martian surface [e.g. 12,15,17,24,25]. 72

73 Mud flow experiments

74 We performed 21 experiments (Table S1 and Fig. S1 in Supplementary Information) using the Open University (UK) Mars Chamber. During each experiment, 500 ml of mud 75 was poured over a 0.9×0.4 m aluminium tray containing either (a) a ~2 cm deep sand bed 76 77 $(\sim 63-200 \,\mu\text{m} \text{ grain diameter}; 14 \text{ experiments})$ representing a sedimentary surface, or (b) a plastic plate (7 experiments) representing an impermeable icy surface. Fifteen experiments 78 79 were performed at 7 ± 0.5 mbar and six experiments at ~1 bar (Table S1). The mud was released onto the surface from a tilting container situated inside the chamber. This design 80 was chosen for its simplicity and reproducibility, although it represents a simplification of 81 the natural setting. At the beginning of the experiment the mud was above the freezing point 82 of water. The temperatures of the sand bed or plastic plate ranged from ~244 K to 265 K and 83 84 gradually increased with time as no active cooling of the experiment was performed. The aluminium tray was inclined at 5° (18 experiments) or 10° (3 experiments) to force the mud 85 to move in a preferred direction. Each experimental run was performed at least in triplicate 86 and was recorded with four cameras. The experiments did not account for the effect of the 87 lower gravity on Mars as compared to Earth. 88

At the beginning of each experiment the atmospheric pressure was gradually reduced, triggering the boiling of the water in the mud [25,27,28]. The mud within the container cooled by evaporation [25] to almost its freezing point before being poured onto the surface at a pressure of 7 ± 0.5 mbar (Table S1). Once in contact with the cold surface, the mud rapidly began to freeze at the bottom and margins of the flow, and at its upper surface (Fig. 2a and Supplementary Information). The freezing resulted in the formation of an ice-mud crust which modified flow propagation and decreased lateral spreading (Fig. 3).

Mud propagation occurred through an intricate system of narrow flow lobes (Fig. 3b) or several lobate flows (Fig. 3c and 3d). Their formation was controlled by the development of frozen marginal ridges that confined the flow of liquid mud inside a central channel. As 99 freezing continued, ice crystals floated to the surface and started to merge. However, mud 100 still propagated within the crust via a network of "mud tubes", in an analogous way to flow 101 within lava tubes (Fig. 2c). When new pulses of mud arrived, they caused breakouts and the 102 formation of further lobes (Fig. 2b). The newly extruded material rapidly developed a frozen 103 crust.

The presence of internal mud tubes was confirmed by sectioning the frozen mud flows after the re-pressurization of the chamber. A liquid mud core was present even in the experiments where mud was exposed to low pressure for several tens of minutes (Fig. 2d,e). Vesicles ranging in size from 1 to 10 mm were observed within the crusts (Fig. 2d), produced by vapour bubbles that did not escape. The vesicular nature of the crust inhibited the conduction of heat [29] from the interior of the flow, increasing the depth to which vapour bubble nucleation occurred.

111 Comparison with terrestrial lava flows

During flow formation, newly supplied mud was observed to increase the thickness of 112 lobes up to several centimetres (Fig. 3d) via lifting of the protective crust. This occurred 113 when the terminal part of the flow was frozen and the mud release was blocked, but newly 114 supplied mud was still intruding the lobe via mud tubes. This created overpressure within 115 the mud which was able to lift the crust of the lobes. Once a sufficient volume of mud had 116 accumulated within the lobe, the overpressure was able to break the crust and a new lobe 117 formed at the terminus. The mud flow inflated in a manner directly analogous to that of 118 pahoehoe lava flows [30]. This inflation was observed on 5° and 10° slopes, and for 119 120 impermeable and permeable substrates. Under terrestrial atmospheric conditions (i.e. room pressure and temperature) a mud flow moving over a cold surface did not form lobes, did 121 not inflate, and had no icy crust (Fig. 3d) regardless of the temperature of the mud (~274, 122 ~290, and ~293 K were tested). Instead, the mud spread out over the surface in a broad sheet 123

only a few millimetres thick, and was in the liquid phase over the entire length of the flow
(Fig. 3a,d). Only minor freezing was observed in the form of ice crystals on the margins of
the flow after several minutes.

Like basaltic lavas on Earth, low viscosity mud flows produce laterally extensive 127 structures with lower relief than those resulting from high viscosity flows [6]. Because the 128 mud used in our experiments is water-dominated, it initially behaves as a Newtonian fluid. 129 In the low-pressure environment, evaporative cooling leads to the formation of ice crystals, 130 which increase the solid content forming a protective crust. The mud evolves into a non-131 Newtonian fluid with non-zero yield strength as the total solid volume fraction increased 132 beyond ~15% [31]. This behaviour is similar to that of low viscosity basaltic lavas whose 133 movement is affected by the formation of an external crust, formed by solidification of the 134 lava due to cooling [32]. In both cases the strength of the visco-elastic part of the crust, 135 between the brittle outer part and the more fluid interior [30], is able to inhibit lateral 136 spreading and allows fluid accumulation and vertical inflation. 137

Over time, the water in the flow (i.e. not lost by evaporation) freezes to form ice in the 138 crust, and hence, the thin crust develops an ever-increasing yield strength due to the ice 139 crystal network becoming increasingly interconnected between the clay particles (see 140 rheology references in [33]). This creates a solid with mechanical strength similar to how 141 mineral crystal frameworks in a cooling silicate rock magma replace the rheological yield 142 strength of the lava. Thus the strength of the crust increases with time, and more inflation 143 should occur before breakouts occur, in agreement with our observation of more inflation in 144 the distal parts of the experimental flows. There is no reason to assume that the mechanical 145 strength of the crust would be different on Mars for mud with the same clay content. Since 146 the rate, as well as the mode, of its formation are controlled by the non-equilibrium 147 thermodynamics of vapour loss, specifically the transfer of sensible heat to latent heat in the 148

liquid surrounding each nucleating vapour bubble [25], the time-scale of crust development
should also be the same on Mars. Calculations describing these processes are reported in the
Supplementary Information.

The above arguments assume that, as in the experiments, the motion of the mud is laminar. 152 However, it is likely that in large flows, that the fluid motion is turbulent. For steady, uniform 153 laminar flows of mud with viscosities of 0.01, 0.1, 1 and 10 Pa.s on 0.6° slopes under martian 154 gravity, the transition to turbulence will occur for mud flows with thicknesses of 17 mm, 155 8 cm, 36 cm and 1.7 m, respectively (see equations in [12]). As the solid content (silicate 156 157 particles and ice crystals) of the mud increases, its viscosity increases, but even for a total solids content of 60% [e.g. 34] the viscosity does not exceed 1 Pa.s. Thus any low viscosity 158 mud flow thicker than a few tens of centimetres on Mars will probably be turbulent and we 159 160 discuss the applicability of our experiments to this case below.

161 Implications for mud flows on planetary surfaces

In a model of mud flow dynamics on Mars, [12] considered the effect of high solid content 162 on the non-Newtonian rheology of mud, assuming the rheological properties were constant 163 everywhere along the flow. Our experiments underline the importance of considering the 164 thermodynamics of the processes that occur when mud is exposed to the martian 165 environment. In a turbulent mud flow, mud from all depths will be exposed to the low 166 atmospheric pressure and will boil, lose vapour, and cool. Thus, it will cool rapidly, 167 increasing the ice crystal content and evolving a yield strength. As the yield strength 168 increases, the critical Reynolds number required to sustain turbulence increases [35], but the 169 170 associated increase in viscosity causes the actual Reynolds number to decrease. When the two become equal this forces a transition to laminar flow, encouraging the formation of lobes 171 and breakouts, as observed in our experiments. 172

Although our small-scale experimental mudflows show characteristics similar to those of low-viscosity lava flows, they have a different heat loss mechanism. For lava flows simulated in the laboratory using wax [e.g., 36-37] a modified Péclet number has been used to distinguish flows that do or do not readily form crusts. This is not appropriate for our flows, where the heat loss mechanism is decompressional boiling, not conduction and convection into the environment (for details see Supplementary Information).

On Earth, low effusion rates and volumes can lead to the formation of mud flows 179 (Fig. 1a,b) with similar magnitudes to those in our experiments. Therefore, such terrestrial 180 flows exported to Mars would be strongly affected by the mechanisms observed in our 181 experiments. We expect these mechanisms to also be important at the margins of kilometre-182 scale flows where their thicknesses decreases and numerous overlapping centimetre-thick 183 184 flows occur. Once mud is extruded into the martian surface pressure environment, it would cool due to evaporative cooling (Fig. 4a) and freezing of the flow surface would eventually 185 cause a change in the flow regime. From that point, the mud would no longer propagate via 186 open channel(s) (Fig. 4b), but instead via mud tube(s) (Fig. 4c) and lobes (Fig. 4d). The 187 distance such a transition occurs from the source would depend on the effusion rates, 188 189 volumes, viscosity and temperature of the extruded mud.

Our study shows that inflation of decimetre-thick and meter-sized lobes could occur at the margins of both mud flows and lava flows, thus both igneous and mud volcanism surface flows could have similar morphological characteristics [e.g., 12,17]. Our calculations show that the morphologies of mud flows at scales larger than those covered by our experiments could be affected by the same processes, particularly at their margins. Hence, as their mechanisms of mud propagation would differ, martian mud volcanoes may be substantially different in shape from terrestrial ones [17]. 197 Sedimentary volcanism has also been proposed for the dwarf planet Ceres [38,39], which may have a water-muddy ocean beneath a crust made of clays, salts, clathrates and ice 198 [39,40]. The process of evaporative cooling and associated freezing should also occur there, 199 200 affecting the morphologies of resulting effusive landforms; even more so on Ceres than on Mars as Ceres lacks an atmosphere. The same principles would apply to other Solar System 201 bodies and icy moons, and so our experimental results should be considered when 202 interpreting effusive cryovolcanic surface features on these bodies [e.g., 41,42]. Our results 203 show that it is vital to consider the effects of the differing environmental conditions on other 204 planetary surfaces when comparing analogue landforms observed on Earth with apparently 205 similar effusive morphologies on other bodies. 206

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334 Author contribution

The experimental set-up and the methodology were conceived and designed by P.B. and O.K. with the help and advice of S.J.C., J.R., M.R.P., M.R.B., A.M., and E.H. The technical support was provided by M.R.S. The data analysis was done by P.B. with significant feedback from O.K., L.W., S.J.C., E.H. and A.M. The DEM production was done by O.K. and the theoretical considerations associated with scaling were done by L.W. All authors contributed to discussion, interpretation and writing of the manuscript.

341 Competing interests

342 The authors declare no competing interests.

343 Figure captions

Fig. 1. Examples of surface expressions of terrestrial sedimentary volcanism caused by muds of various viscosity. (a) A water-dominated mud flowing from the crater of Bakhar mud volcano in Azerbaijan (39°59'55.7"N, 49°28'29.9"E). (b) An individual mud flow from a 'gryphon' on top of Dashgil volcano in Azerbaijan (39°59'48"N, 49°24'11"E). (b). (c) A kilometre-sized highly-viscous mud flow outgoing from Koturdag mud volcano in Azerbaijan (39°58'30"N, 49°21'36"E). Note minibus for scale.

Fig. 2. Examples of morphologies and interior structures of mud flows formed in a lowpressure environment. Panel a shows three frames from video taken by Cam #2 covering experiment #16 in which the formation of a narrow, thick mud flow occurred. Panel b shows in detail the formation of the icy crust and outbreaks of new mud pulses from beneath the icymuddy crust. When the resulting mud flow features were sectioned, large cavities filled with liquid mud (c), or voids (d) in the ice were observed. A liquid mud core was commonly observed in the flow interior (e). Fig. 3. Timeline maps of modelled mud flows derived from the videos and final topographic cross sections. Flows formed in terrestrial (a) and low-pressure environments when mud was poured by high (b) or low (c) release rates. The numbered bold lines represent the margins of the flows at 5 s intervals, the other lines are 1 s intervals. (d) The position of each topographic profile is marked in panels a, b and c. Topography was measured after the chamber was re-pressurized.

Fig. 4. Hypothesised development of a low viscosity mud flow on Mars. (a) The mud boils and self-cools through evaporative cooling. Once a mud flow develops, the mud is transported via an open channel (b) evolving into a mud tube due to freezing of the flow surface (c). The mud propagates via mud tubes to the flow front. As the crust prevents free movement of the mud, it spreads via lobe breakout (d) which occurs when the pressure inside the frozen flow is large enough to break the crust or to lift it up; exposing the mud again to the martian environment.

370 Methods

Experimental setup. The low viscosity mud used in our experiments was a mixture of water 371 that contains 0.5 % w/w dissolved magnesium sulphate (MgSO₄) salts corresponding to the 372 average river water salinity and clay obtained from the claystone named after the Rokle locality 373 operated by the company Keramost, which is situated near the town of Kadaň in the Czech 374 Republic. The clay was a bentonite composed of 76 % montmorillonite, 23 % illite, and 1 % 375 kaolinite which has been formed by alteration of pyroclastic rocks. As there is no direct in-situ 376 knowledge of which types of clays could be present on Mars during subsurface sediment 377 378 mobilization and as explosive volcanism was once present on Mars [42], to the first approximation this material seems to be a suitable analogue. To exclude the presence of 379 potentially more lithified clayey aggregates, the clay was mixed with water and salt and 380 homogenized in a blender for 3 minutes. Adding the small amount of the salt was necessary to 381

allow submillimetre particles to get into suspension within the mixture [43] and also realistic
for the martian surface [e.g., 44-46]. The resulting viscosity of the mud was 12.7 mPa.s at ~276
K and 10.7 mPa.s at ~296 K and the average density of the mixture was 1037.5 kg.m⁻³. The
viscosity was measured with Haake Rotovisco RV 20 and Viscotester VT 550 rheometers
(Institute of Hydrodynamics of the Czech Academy of Sciences) with ledges on the MV2
cylinder to prevent slip of the measured material on its walls.

Each experimental run (for details see Table S1 in the Supplementary Information) started 388 by inserting the 0.9×0.4 m aluminium tray filled with (a) a ~2 cm deep substrate bed (natural 389 390 sand, $\sim 200 \,\mu\text{m}$), or b) a plastic plate, and similarly sized copper plate inside the freezer to pre-cool the tray and the plate to temperatures around 238 K. Once the required temperature 391 was reached, the plate and the tray were inserted inside the vacuum chamber. At the same 392 393 time the 500 ml of liquid mud was poured inside the tilting container equipped with one thermocouple to record the temperature of the mud, and the container was installed inside 394 the chamber. The temperature of the mud varied from 274 to 297 K before the pressure drop. 395 The temperature within the chamber was also monitored by another thermocouple. 396 Additionally, five thermocouples were set in a grid (see Fig. 1 in the Supplementary 397 398 Information for details about the positioning of thermocouples) within the tray in order to monitor the temperature of the surface over which the mud propagated. 399

Once the tray was in place inside the chamber, a series of images were taken by a singlelens reflex camera from different angles to obtain the digital elevation model of the pristine surface before the experimental ran. Subsequently the chamber was closed and the process of depressurization started. To achieve the pressure drop from ambient terrestrial pressure to 7 mbar took usually around 6 minutes. Once the pressure started to drop, the decrease in the temperature of the mud within the container was measured. Every time the mud self-cooled close to 273 K during the pressure drop, but it remained liquid. When the pressure of ~7 mbar 407 was reached, the container was manually flipped by the operator and hence mud was poured 408 from the height of \sim 5 cm to the surface. The mud flux was not directly measured, we recorded 409 only the time for how long the mud was extruded (Table S1), as the intense boiling occurred 410 within the container and hence caused irregularities within the flux.

The movement of the mud over the surface was recorded by four video cameras. Once the 411 mud propagation stopped, the resulting mud flow feature was left in the low pressure 412 environment for various lengths of time ranging from several minutes to about one hour. After 413 that the process of re-pressurization of the chamber to terrestrial values started, typically 414 415 before temperature of the tray surface rose above ~273 K. Once the pressure inside the chamber reached atmospheric pressure, the chamber was opened and the resulting flow features were 416 documented by taking images from different angles to acquire data for subsequent DEM 417 production. Ultimately the mud flows produced were sectioned and their inner structure was 418 investigated and documented. 419

420 **DEM production.** To compare the elevation profiles along and perpendicular to the flow directions of the mud flows we calculated a series of digital elevation models (DEM). The 421 sedimentary bed was photographed after each experimental procedure ~30-70 times from 422 multiple viewpoints. The reconstruction of a 3D model surface was produced by using the 423 'Structure-from-Motion' [47] commercial software Agisoft PhotoScan. For image orientation 424 correlation and scaling of the 3D models we used twelve fixed black-on-white printed markers 425 which were affixed onto flat topped cylindrical posts. The posts had two different elevations 426 (4.6 cm and 9.6 cm) and the markers were ~2.67 cm in diameter. Typical discrepancies between 427 actual and calculated marker positions were ~ 0.8 - 1.6 mm. Exported DEMs and orthophotos 428 (TIFF format) were imported to QGIS for further analysis and production of the elevation 429 profiles. 430

431 Cooling and freezing of mud when the external pressure is less than the saturation vapour pressure of the water. Both the silicate component and the water that has not vet 432 vaporized will cool from the initial temperature until the freezing point is reached. After the 433 freezing point is reached, the temperature remains constant and vapour continues to be lost until 434 all of the remaining water has been converted to ice. If the external pressure is less than the 435 saturation vapour pressure of the ice, evaporation continues, the frozen mud surface cools, and 436 a wave of cooling propagates into the frozen mud. The bulk density of the mud changes 437 continuously throughout these stages as a function of the initial mass fractions of silicate and 438 439 water in the mud. Let the initial masses of water and silicate in a given sample of the mud be m_w and m_s , respectively, the volumes of water and solid be v_w and v_s , respectively, and the 440 corresponding densities be ρ_w and ρ_s . Let the bulk density of the mud be β . Then 441

442

443
$$\beta = (m_w + m_s) / (v_w + v_s) = (m_w + m_s) / [(m_w / \rho_w) + (m_s / \rho_s)]$$
(1)

444

446

447
$$m_w / m_s = (\rho_w / \rho_s) [(\rho_s - \beta) / (\beta - \rho_w)]$$
 (2).

448

We next find the mass of vapour, M_{vc} , that must be lost from the water to cool the remaining water and silicate to any given lower temperature. Let the initial mud temperature be θ_i and the final temperature be θ_f . The heat H_v removed from the water by the formation of the vapour is

453

454
$$H_v = M_{vc} L_v$$
 (3)

liquid water mass decreases from its initial value m_w to a smaller final value m_f, where by 457 definition 458 459 $m_f = m_w - M_{vc}$ 460 (4). 461 Thus, the average mass of water, m_a, during the cooling process is 462 463 $m_a = 0.5 (m_w + m_f) = 0.5 (m_w + m_w - M_{vc}) = (m_w - 0.5 M_{vc})$ (5). 464 465 As long as the specific heat of the water can be approximated as a constant, the heat the water 466 loses while cooling is H_c where 467 468 $H_c = m_a C_w (\theta_i - \theta_f)$ (6) 469 470 and Cw is the specific heat of water, 4186 J kg⁻¹ K⁻¹. The silicate mass ms also cools, and loses 471 an amount of heat equal to H_s where 472 473 $H_s = m_s C_s (\theta_i - \theta_f)$ (7). 474 475 Here C_s is the specific heat of the silicate, say 1000 J kg⁻¹ K⁻¹. Equating the sum of H_c and H_s 476 to H_v, 477 478 $(m_w - 0.5 M_{vc}) C_w (\theta_i - \theta_f) + m_s C_s (\theta_i - \theta_f) = M_{vc} H_v$ (8) 479 480

where L_v is the latent heat of vapourization, 2.46×10^6 J kg⁻¹. During this cooling process, the

and regrouping,	
$M_{vc} \left[H_v + 0.5 C_w \left(\theta_i - \theta_f\right)\right] = (m_w C_w + m_s C_s) \left(\theta_i - \theta_f\right)$	(9)
ог	
$M_{vc} / m_s = \{ [(m_w / m_s) C_w + C_s] (\theta_i - \theta_f) \} / [H_v + 0.5 C_w (\theta_i - \theta_f)]$	(10).
Using equation (2) for (m_w/m_s) we can also find (M_{vc}/m_s) as a function of the assum	ed value
of ρ_s . Finally, the ratio, R, of the mass of water converted to vapour to the initial wat	er mass,
i.e. $R = (M_{vc} / m_w)$, is equal to $[(M_{vc} / m_s) (m_s / m_w)]$ or more conveniently	
$R = [(M_{vc} / m_s) / (m_w / m_s)]$	(11).
The mass of water remaining in the mud after the cooling phase, m _f , is therefore	
$m_{f} = (1 - R) m_{w}$	(12).
The above equations apply between any pair of temperatures θ_i and θ_f until θ_f l	becomes
equal to the freezing point, θ_{fr} . After the freezing point is reached, the temperature	remains
constant while liquid water continues to evaporate and the latent heat of vaporiz	zation is
extracted from the remaining water, progressively freezing into ice all of the water th	at is not
lost as vapour. The latent heat of vaporization in the 283-293 Kelvin range is 2.46×1	0 ⁶ J kg ⁻¹
and the latent heat of solidification is 3.34×10^5 J kg ⁻¹ . The ratio of these is Q = (3.34)	4×10^{5} /
	and regrouping, $M_{ve} [H_v + 0.5 C_w (\theta_i - \theta_f)] = (m_w C_w + m_s C_s) (\theta_i - \theta_f)$ or $M_{ve} / m_s = \{[(m_w / m_s) C_w + C_s] (\theta_i - \theta_f)\} / [H_v + 0.5 C_w (\theta_i - \theta_f)]$ Using equation (2) for (m_w / m_s) we can also find (M_{ve} / m_s) as a function of the assum of ρ_s . Finally, the ratio, R, of the mass of water converted to vapour to the initial wat i.e. R = (M_{ve} / m_w), is equal to [(M_{ve} / m_s) (m_s / m_w)] or more conveniently $R = [(M_{ve} / m_s) / (m_w / m_s)]$ The mass of water remaining in the mud after the cooling phase, m _f , is therefore $m_f = (1 - R) m_w$ The above equations apply between any pair of temperatures 0 _i and 0 _f until 0 _f ! equal to the freezing point, θ_{fr} . After the freezing point is reached, the temperature constant while liquid water continues to evaporate and the latent heat of vaporiz extracted from the remaining water, progressively freezing into ice all of the water th lost as vapour. The latent heat of vaporization in the 283-293 Kelvin range is 2.46 × 1 and the latent heat of solidification is 3.34×10^5 J kg ⁻¹ . The ratio of these is Q = (3.3)

 2.46×10^6) = 0.13577. Thus, to produce 1 kg of ice we would have to evaporate 0.13577 kg of 505

506	water into vapour from an initial total mass of 1.13577 kg of water. A fraction $[Q/(1 + Q)]$ of							
507	the water mass remaining after cooling must become vapour and a fraction $[1/(1 + Q)]$ of the							
508	water mass remaining after cooling becomes ice. The final ice mass, m _i , is							
509								
510	$m_i = [1/(1 + Q)] [(1 - R) m_w]$	(13)						
511								
512	and the mass of water converted to vapour during the freezing phase is $M_{\rm vf}$ where							
513								
514	$M_{vf} = [Q/(1 + Q)] (1 - R) m_w$	(14).						
515								
516	The total mass of vapour generated by the whole process is $M_v = M_{vc} + M_{vf}$.							
517								
518	As a result of the loss of vapour, the bulk density of the frozen mud will be diffe	erent from						
519	the density of the initial mixture. The ice has a density ρ_i of 916.8 kg m ⁻³ so the mass m _i of ice							
520	has a volume of $v_i = (m_i / \rho_i)$. The silicate volume is still v_s and so the final bulk density is β_f							
521	where							
522								
523	$\beta_{f} = (m_{i} + m_{s}) / (v_{i} + v_{s}) = (m_{i} + m_{s}) / [(m_{i} / \rho_{i}) + (m_{s} / \rho_{s})]$	(15).						
524								
525	Equation (12) gives (m_i / m_w) and equation (2) gives (m_w / m_s) so in terms of these,							
526								
527	$\beta_{f} = [(m_{i} / m_{w}) + (m_{w} / m_{s})^{-1}] / \{[(m_{i} / m_{w}) / \rho_{i}] + [\rho_{s} (m_{w} / m_{s})]^{-1}\}$	(16).						
528								
529	The density, ρ_s , of the clay minerals in the experimental mud was ~2500 kg m ⁻³ and	d the bulk						
530	density of the mud was ~ 1040 kg m ⁻³ , implying that the clav component formed \sim	-6.5 % of						

531 the mud mass. The mud was released into ambient experimental chamber pressures in the range 650-700 Pa, and initial mud temperatures were up to ~275.5 K. Taking account of the 532 weight of the overlying mud and the experimental chamber pressure, the saturation vapour 533 pressure of water, with values up to \sim 730 Pa [48], would have been reached at depths up to 534 5-6 mm in the experiments. The relative values of the specific heat and the latent heat of 535 evaporation of water are such that while the mud was cooling from ~275.5 K to its freezing 536 point, ~0.34 % of its water would have been lost, having a trivially small effect on its 537 essentially Newtonian rheology. By the time subsequent vapour loss had frozen the 538 remaining water, ~88 % of the initial mass of water would have been converted to ice at the 539 expense of losing ~ 12 % of the initial water mass as vapour, leaving solid mud with a density 540 of ~961 kg m⁻³, slightly less dense than the original liquid mud. This should have produced 541 542 a frozen crust, again of thickness 5-6 mm, beneath which the mud would have been partially liquid, as observed in the experiments. In similar scale flows under martian gravity, the 543 thermodynamics of this process would have been the same, as it involves only heat transfer 544 545 by conduction, but the thickness of the frozen outer crust would have been greater, 14-17 mm, because the pressure in the mud depends on the acceleration due to gravity [49]. 546

547 **Data availability**

The movies, photos, pressure and temperature logs generated during and analysed during the current study that support our findings are available in the Zenodo repository with the identifier DOI: 10.5281/zenodo.3457148 (<u>https://doi.org/10.5281/zenodo.3457148</u>).

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a) ambient pressure, variable release rates



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ridges /

/ lobes



578

ι



Exp #	Pressure range [mbar]	Inclination [°]	Release time* [s]	Surface T** [°K]	Mud T*** [°K]	Type of surface	Salinity****	
5	6.55 - 7.33	5	15	265	294	sand	D.I. water	
6	7.16 - 7.66	5	26	264	292	sand	tap water	
7	6.43 - 7.02	5	27	258	296	sand	saline water	
8	7.11 - 7.37	5	25	261	294	sand	saline water	
9	7.09 - 7.63	5	27	262	294	sand	saline water	
10	1000	5	21	258	290	sand	saline water	
11	1000	5	19	264	274	sand	saline water	
15	1000	5	15	258	293	sand	saline water	
16	6.32 - 6.54	5	46	256	290	sand	saline water	
17	6.57 - 6.94	5	34	252	278	sand	saline water	
18	6.37 - 6.58	5	15	260	290	sand	saline water	
19	6.06 - 6.81	10	35	258	286	sand	saline water	
21	6.55 - 6.77	10	34	253	293	sand	saline water	
22	7.08 - 7.36	10	37	254	294	sand	saline water	
23	6.43 - 6.97	5	40	265	293	plastic plate	saline water	
24	6.66 - 7.11	5	28	254	295	plastic plate	saline water	
29	6.55 - 7.08	5	40	259	295	plastic plate	saline water	
34	5.29 - 6.55	5	failed exp.	259	297	plastic plate	saline water	
41	1000	5	42	244	295	plastic plate	saline water	
49	1000	5	34	247	290	plastic plate	saline water	
54	1000	5	34	249	277	plastic plate	saline water	

 $\ensuremath{^*}$ Time period over which the mud was poured from the container

 $\ast\ast$ Temperature of the surface before the release of the mud from the flipping container

 $\space{1.5}\space{1.$

**** Saline water refers to a mixture of water that contains 0.5% w/w dissolved magnesium sulphate (MgSO₄)