Airflow Measurement across Negatively-Infiltration Processed Porous Aluminium Structures

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Abstract

Detailed structural characterization and experimental measurement of airflow across open-celled aluminium foam structures with near-spherical cells varying in pore sizes and interstices are presented herein. The aluminium foam structures were produced by infiltrating liquid aluminium into convergent gaps created by packed beds of near-spherical hydrosoft salt beds varying in particle sizes, packing densities and infiltration pressures. A quantitative assessment of the unit pressure drops developed across these structures show that viscous and inertial terms of these structures were observed to greatly depend on the shape and structural macroscopic parameters of the porous medium.

Keywords: Porous Metal; Measurement; Pore Sizes; Openings.

Introduction

Engineers and scientists in the past have determined the permeabilities of many soils and sands varying in degree of particle sizes and interstices measured in Darcy (D) or milliDarcy (mD) to avoid the intricacy of prefixes like terms (10^{-09}, 10^{-10}, 10^{-11}, 10^{-12} etc). The numerical value of permeability, k_o, measured in the horizontal or vertical plane for a given packed beds of rocks depends on the degree of its particle sizes and interstices [1] and the addition of a correction factor, Kozeny constant [2-6], was proposed to bridge the gap between analytical model and reality due to the anisotropic and complicated tortuous flow paths of the porous medium.

While the packed beds are often made from packing of naturally occurring rocks and sands, artificially-made porous structures made from metals and alloys such as metal foam (aluminium, nickel, titanium, copper, steel etc) are remarkably useful as structures that interact with a fluid in processes like sound absorption, filtration and catalysis, heat exchange and storage etc. An important feature of this structure is the resistance provided by its skeletal configuration when presented to flowing fluid across the pores and interstices of the porous matrix, often described by the unit pressure drop developed across the inlet and exit of the structure as a function of superficial fluid velocity and properties of the fluid. Measuring or estimating and controlling the unit pressure drop developed across this structure is imperative to the design of porous metallic structures with enhanced performance.

Flow through open-celled metallic foams (Fig 1) are more complicated by the opening of its windows or interconnection of its pores [7]. This interconnection in porous metal foam enabled its high porosity, permeability, large foam surface area per unit bulk volume (specific surface) and thermal conductivity [8-10] which has made it beneficial for use as pressure reduction and heat transfer devices [10]. Understanding the behaviour of fluid flow through commercially available porous metallic structures like Porvair, Recemat, Alantum, Alporas, Duocel etc with open porosities between 80-95% has been proven experimentally in [11-14]. The unidirectional (horizontal) flow permeability and inertial coefficient were calculated by fitting the pressure and velocity data into the so called two-term Hazen-Dupuit-Darcy also known as Darcy-Forchheimer (D-F) model (Eq 1). This two term Darcy-Forchheimer equation relates the developed pressure gradient (∇P) for a moving fluid across porous medium as a function of velocity vector (v_s), permeability (k_o), fluid density (ρ), fluid dynamic viscosity (μ) and the Forchheimer coefficient (C_F).

$$-\frac{\partial p}{\partial x_i} = \nabla p = \frac{\mu}{k_o} v_s + \rho \frac{C_F}{\sqrt{k_o}} v_s^2 \qquad \text{Eq 1}$$
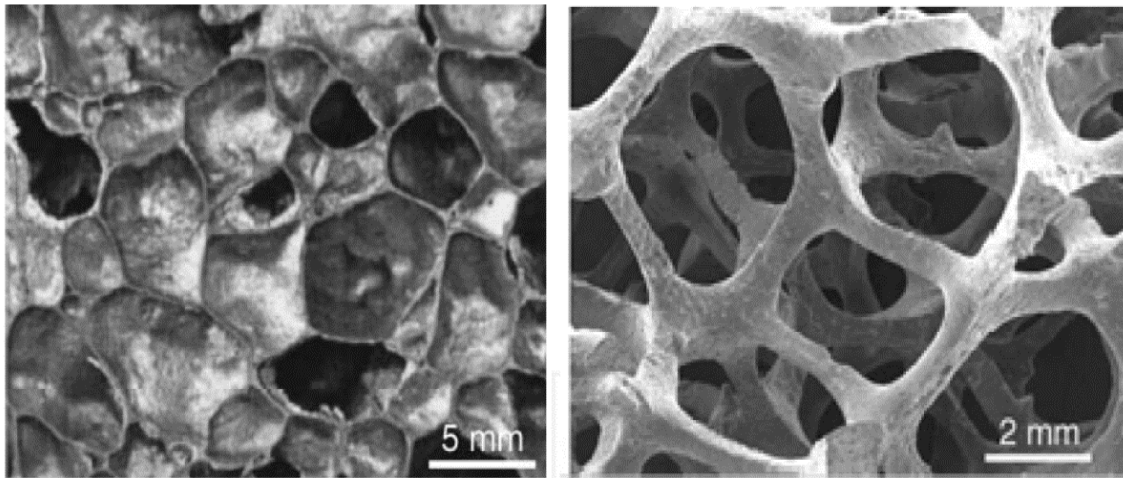

Fig 1. Images of close (Left) and open (Right) celled metallic foam structures (adapted from [17]).

Generally, the pressure drop per unit thickness (∇P) of any porous structure in the Phillipe Forchheimer model can be expressed solitarily in the first term of the LHS of the Eq 1 only at very low velocities (Darcy regime) where losses are mainly due to skin friction whilst the second term predominates at high velocities (laminar-turbulent regime) where viscous forces are negligible while contraction and enlargement are significant [15, 16] (See Eq 2). In addition, extremely high Reynolds number fluid flow could result in the continuous addition of non-Darcy pressure drop and increased drag or inertial forces and the relationship between the pressure drop gradient and superficial fluid velocity, in some

porous structures, can be corrected [30] by the power law or three-term cubic Forchheimer models as shown by Eq 3 and Eq 4 respectively.

$$\nabla p = A v_s \qquad\qquad\qquad \text{Eq 2}$$

$$-\frac{\partial p}{\partial x_i} = \nabla p = A v_s + B v_s^m \;\; (1.6 \leq m \leq 2.0) \qquad\qquad \text{Eq 3}$$

$$-\frac{\partial p}{\partial x_i} = \nabla p = A v_s + B v_s^2 + C v_s^3 \qquad\qquad \text{Eq 4}$$

where A and B are related to the viscous. inertial terms and fluid properties in Eq 1 whilst the last term, C, in the three-term cubic Forchheimer (Eq 4) model was introduced [30] to make the equation reasonably fit to experimental results.

Realistically, the infiltration of liquid metal into near-spherical beds of sacrificial porogens or space fillers has been proven [18, 20] to produce low-density porous "bottleneck-type" materials. Computational fluid dynamics simulation [21] and flow measurement [22] across the low-density porous structure of similar pore sizes and different openings point the permeability of this structure to greatly depend on the sizes of the openings connecting the spherical pores. In this work, a full description of the measurement and flow behaviour of porous aluminium structures having different pore sizes and pore diameter openings is studied to see what benefit therein.

Structural Characterization and Airflow Measurement

Similar experimental procedure on the production of aluminium (Al) foam structure made by replication-casting process presented in [18, 20, 22] was used to produced foam structures used in this work by negatively-infiltrating liquid melts into packed beds of near-spherical hydrosoft-based (sodium chloride) salts. In brief, near-spherical salts of the different sizes typically in the ranges X/1.0-1.4mm; Y/2.0-2.5mm and Z/2.5-3.15mm (Fig 2) were packed in a 40mm diameter and 50mm height vessel and liquid aluminium were driven into the interstices of the beds at some applied differential pressures typically between 0.25 and 0.90 bar. Preheating of the space fillers to within 450-600ºC was done prior to packing to avoid premature solidification of the melt. The pore volume fractions of the porous structures were measured based on their mass-scale difference using digital scales and Vernier calliper whilst optical images of the structures assessed in Image J software reveal in detail, the pore diameter sizes, and openings. Scanning electron microscope (SEM) of the salt samples and optical images of the foam structures produces are presented in Fig 2.
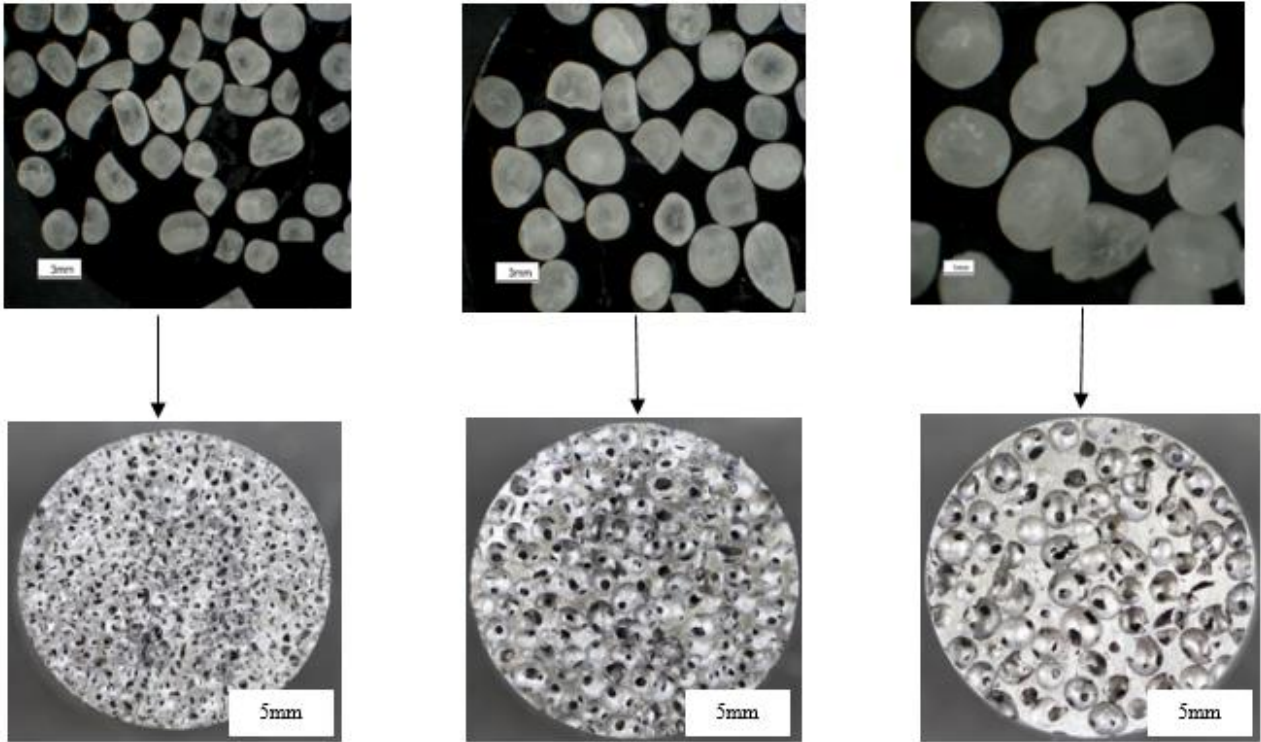
Fig 2. Top, scanning electron microscope (SEM) images of salts used for the casting process for particle size 1.0-1.4mm/X1 (left), 2.0-2.5mm/Y1 (middle) and 2.0-3.15mm/Z1 (right) and bottom, optical microscope images of foam structures produced with particle sizes 1.0-1.4mm/X (left), 2.0-2.5mm/Y1 (middle) and 2.0-3.15mm/Z produced at 0.9 bar.

Advanced imaging techniques working from X-ray computerised tomography (CT) data was also used to enable assessment of the structural properties (pore sizes, openings, roundness and volume) of the porous materials at microscale level as shown in Fig 3. In brief, high resolution (26µm voxel dimension in x, y and z) two-dimensional stack images (Fig 3a) were restored to three-dimensional volume (Fig 3b) in ScanIP module of Simpleware^{TM} (advanced 3D imaging, meshing and physics solving software currently owned by Synopsys computer integrated systems design company, USA). Image processing procedure like thresholding, segmentation, filtering and various image editing tools available within the ScanIP were used to create the three-dimensional volume and to also ensure the right structural material is preserved. A workable representative volume (Fig 3c) was obtained by shrinking large volume of image data until porosity differed by ±2% and the fluid phase (Fig 3d) of the representative matrix was obtained by inverting the masked structural phase. A watershed segmentation was carried out on the fluid domain matrix by disconnecting the particles to enable evaluation of the mean particle/pore sizes of the porous matrix in ScanIP. Additionally, pore diameter openings of the structures were measured by running a centreline across the representative matrix and a mean value of the smaller diameter (openings) was taken and compared to measuring these openings directly in ScanIP as evident by Fig 3g and that obtained using SEM and optical micrograph techniques.
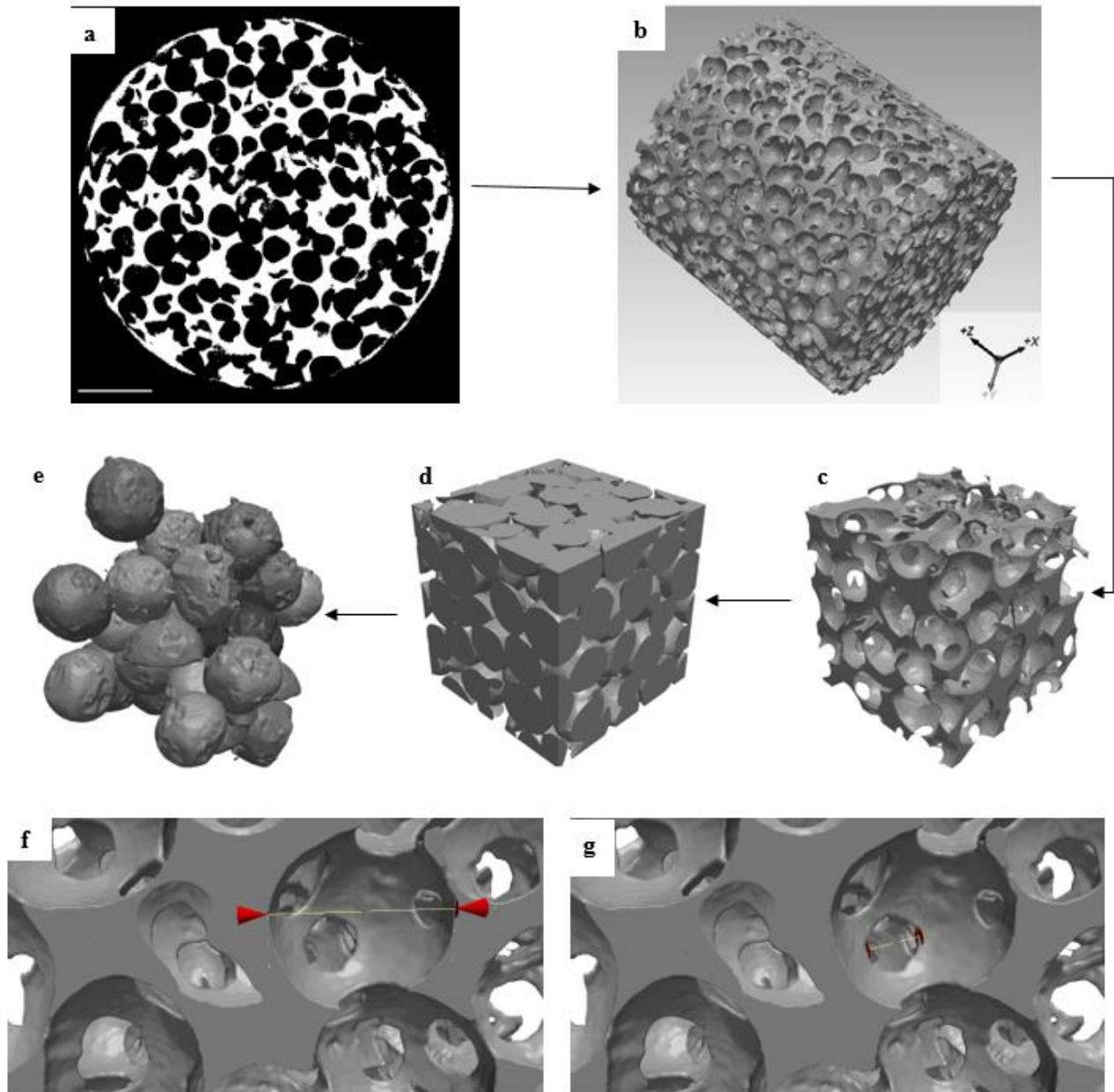
Fig 3 (a) two-dimensional computerised tomography image, (b) three-dimensional reconstructed, (c) representative structural phase, (d) representative fluid phase, (e) segmented pores (f) typical pore sizes measurement and (g) pore openings of measurement of Y1 "bottleneck" structure.

Table 1 presents the macroscopic and flow (to be discussed later) parameters for the different porous metallic structures produced using hydrosoft-based salts for particle sizes in the range of 1.0-1.4mm, 2.0-2.5mm and 2.5-3.15mm represented by the following symbols (X, Y and Z) in the table. X1 (1.0-1.4mm) and Z1 (2.5-3.15mm) structures were produced using a liquid metal infiltration pressure of 0.9 bar whilst the Y (2.0-2.5mm) structures were produced at 0.90 (Y1), 0.50 (Y2), 0.33 (Y3) and 0.25 (Y4) bars. Key parameters like porosity, mean pore size, and mean openings or connectivity measured are presented in Table 1. As expected, structural measurements of pore sizes lie closely to within the

mean average salt particle sizes used for the casting process. The mean openings between connecting pores decrease with increased casting pressure differentials and the pore volume fraction (porosity) predominantly controlled by the packing density of the beds. Though, this is also affected by the variation in liquid infiltration pressure with lower pressure differential leading to the less complete filling of the interstices of the packed beds and the higher level of porosity as shown in Y4 structure.

Table 1. Measured structural and flow parameters of the porous Al structure.

Sample	Infiltration pressure difference, P_i (bar)	Porosity, $\emptyset$ (%)	Mean pore size, D_p (mm)	Mean connectivity, D_w (μm)	Permeability, $k_0 / 10^{-09}$ (m^2)	Forchheimer coefficient, C_F [-]
EXPT-X1	0.90	72.03	1.21 ± 0.18	275 ± 10	4.16	0.92
EXPT-Y1	0.90	70.55	2.21 ± 0.15	643 ± 17	13.37	2.03
EXPT-Y2	0.60	72.55	2.23 ± 0.12	728 ± 24	17.00	0.71
EXPT-Y3	0.45	75.09	2.22 ± 0.16	727 ± 15	19.26	0.62
EXPT-Y4	0.25	78.22	2.25 ± 0.24	920 ± 25	31.24	0.29
EXPT-Z1	0.90	75.08	2.75 ± 0.21	710 ± 23	18.98	0.89

Fig 2 confirms the shapes and sizes of the salt beads to dictate the morphology (Fig 3) and pore sizes of the structures respectively. These also show smaller and multiple rounded "windows" or connectivity between near-spherical connecting pores and the size of these windows increases with increased pore/particle sizes and infiltration pressure differential used as presented in Table 1. Optical, scanning electron microscope (SEM) assessment and tomography imaging of these structures reveal that not only the "window" sizes of the structures are affected with variation in pressure differentials but also the number of contacts (coordination number) between connecting pores is also affected. Previous work in [20, 21] has shown that coordination number per cell is dependent on that obtained for packing of salt beds affected by variation in packing density and particle shapes and sizes. For monosized "virtual" near-spherical monosized structures of roundness close to unity, typical values between 6 and 7 are attained in [20, 23] and higher for distorted near-spherical shaped salt beads in [22] used herein.

The experimental setup in Fig 4 was designed to accurately measure pressure drops across an open-celled aluminium foam structure at a defined flow rate of air movement. The experimental setup and measurement methods were designed like those reported in [13, 14] and in brief, in [22]. In detail, it consists of compressed air supply unit, a Norgren F74G 40um G 1/2" automatic pneumatic filter (working pressure 0/10bar), a nickel-plated brass hydraulic needle valve series (FT 1251/2 & connection thread BSP 3/8"), standard (stable, easy-to-read, ±3% of full scale) precision-machined acrylic Flo-Rite^{TM} air flow meters (0-140LPM) a flow straightener, a mid-flange assembly of foam sample holder with a GEMS sensors pressure transducer (0-2.5/25bar) and a data-log PC installed with 2016 NI LabVIEW signal for data acquisition.
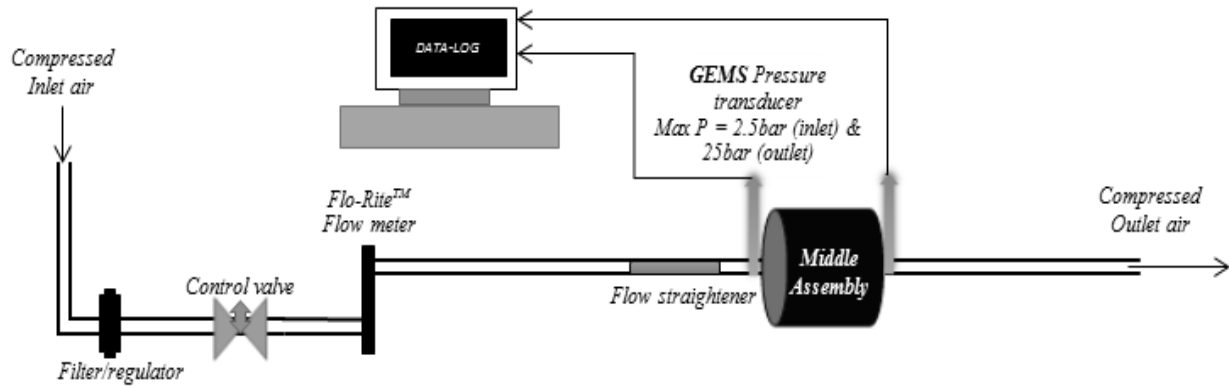
Fig 4. Schematic representation rig set up to measure air flowing across porous metallic structures (adapted from [22]).

A precisely machined foam sample of diameter, 25±0.2mm, and foam thickness 34.5±0.3mm (conferred to be the maximum length of the sample holder in the mid-assembly without the need of a spacer) was wrapped with a polytetrafluoroethylene (PTFE) tape to prevent air-bypass by ensuring total control of air passage only through the pores and "windows" of the foam structure and as well maintain a tighter fit and greater coefficient of friction with the wall of the mid-flange typically, 21.183mm internal diameter (flow diameter). The filter regulator made up of element material sintered polypropylene was used to remove water vapor and impurities (filtration size, 40μm) from the air supply. Two pressure transducers were connected to both the inlet (0-2.5bars) and outlet (0-25bars) of the mid-assembly. Pressure-transducer voltage calibration was done using a digital pressure indicator (DPI 601) for the inlet transducer (0-2.5bars) whilst the outlet pressure was conferred to be atmospheric.

An unmitigated power supply was ensured by using 3 x 9Volts non-rechargeable batteries supplying 8-30volts GEM transducers. This was done to minimize noise effects and as well minimize the gap in error band. The needle valve with a minimum and maximum operating temperature of -20 to +100ºC and 210 bars maximum operating pressure was used to control the outflow of air into the flow meter from the filter regulator operating at 2.5bars. This nickel-plated brass hydraulic needle valve consists of a double-acting valve for controlling and shutting off the flow and a single-action valve for flow control with free counter-current. A plot of pressure and voltage difference measured across the foam structure shows a perfectly fitted (P-ΔV) linear correlation (Fig 5) of mean standard deviation 1.63 and 0.77% determined for the inlet and exit transducers respectively. A minimum steady time of 2 minutes was allowed before taking the average sum of 2,000-3,000 voltage data recorded in a space of 2-3 minutes by the data acquisition PC device translating to upstream and downstream gauge pressures.
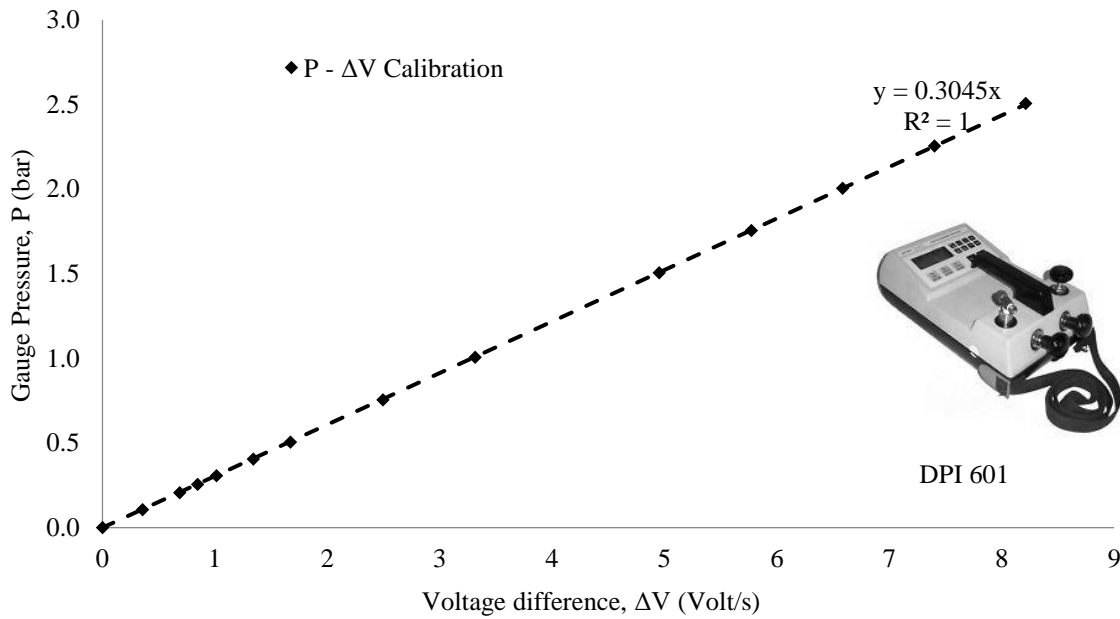
Fig 5. Inlet pressure-voltage calibration a digital pressure indicator (DPI 601).

The flowmeter available in metric scales was used to determine the air flow rate (0-140LPM) before it flows through the mid-flange (sample holder) via the flow straightener. Upstream and downstream air pressures drop across the foam sample were taken with the help of the GEMS pressure transducers and data-log for the given range of flow rates. The measured upstream and downstream gauge pressures were converted to absolute pressures by simply adding the measured gauge pressures with atmospheric pressure [8, 13, 24, 25]. Due to the fractional change in the air as opposed to liquid, care was taken to account for compressibility effects to avoid significant underestimation of the static pressure variation caused by an alteration in the density of the flowing fluid (air) across the foam cross section [8, 14, 26, 28]. Eq 5 was applied to determine the "real" pressure drop (ΔP) across the structure as a function of the inlet (P_i) and outlet (P_0) absolute pressures while P_R was taken as reference or atmospheric pressure.

$$\Delta P = \frac{P_i^2 - P_o^2}{2P_R} \qquad \text{Eq 5}$$

To ensure accuracy of the experimental setup in Fig 4 is preserved and the right data were obtained, care was taken to the first test for leak detection and an airflow measurement on a 20mm thick Inconel 450μm (metal alloy) foam sample (Fig 6) of density and pore volume fraction, 828kg/m^3 and 84% respectively was carried out and compared to existing data available in the literature [14]. A plot of pressure drops per unit length (Pa.m^{-1}) against superficial air inlet velocity (m.s^{-1}) for the Inconel foam

structure followed a second order Forchheimer relation (Eq 1) and a reasonable agreement of 99.7% was attained in comparison with the literature data.

Also, airflow measurement through a 12mm thick Recemat RCM-NCX 1723 sample shows reasonable agreement with numerically simulated (using X-ray tomography images of the "real" structure) and measured values of unit pressure drop developed across the foam sample reported in [19]. Though, flow length reported in [19] was said to have a considerable influence on the flow performance due to the inhomogeneous topology (combination of irregular dodecahedron and tetrakaidecahedron cell with pentagonal faces described as the [29]) of the sample geometry resulting in a non-uniform connectivity and variable length thereby necessitating a more developed flow field (increased computational time) to fully describe the entire flow behaviour in a large sample especially at higher flow velocities. Flow behaviour through porous aluminium structures herein may differ from the Inconel 450μm and Recemat RCM-NCX 1723 samples but the observable agreements between the measured and literature data confirm the precise experimental setup and the sureness of any measured data of unit pressure drop across structures of any morphology, cell sizes, and openings.
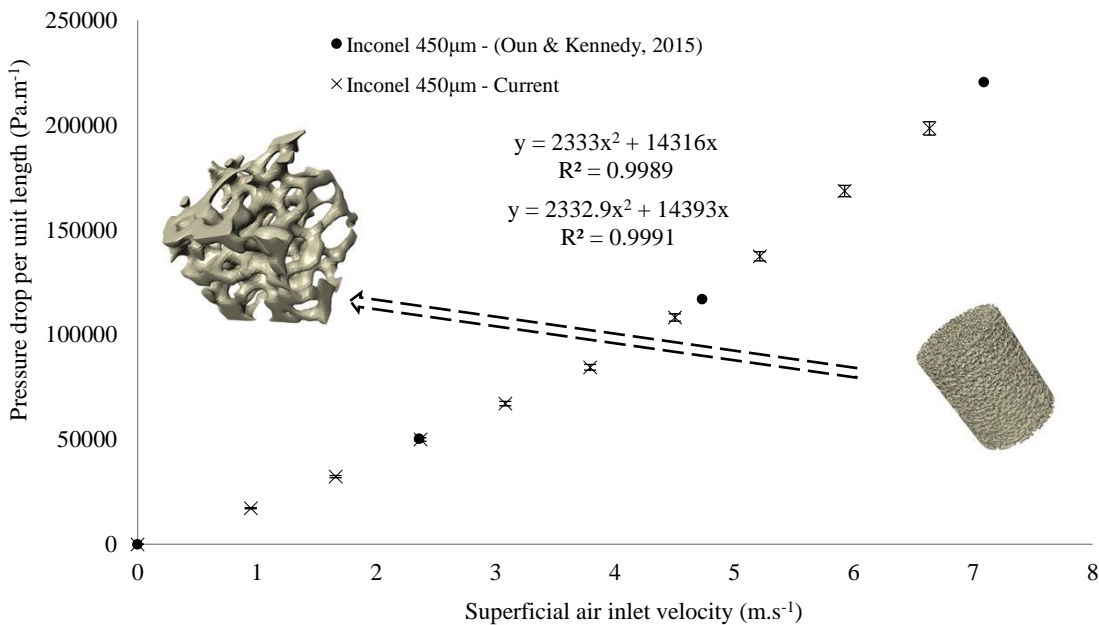

Fig 6. X-ray CT images of Inconel 450μm foam structure and a plot of pressure drops per unit length (Pa.m^{-1}) against superficial air inlet velocity (m.s^{-1}) for the structure.

Experimental [8, 25, 30, 31, 39], analytical [16] and computational models [21, 32, 33, 34, 40, 41] available in the literature have shown that the pressure drop developed across a porous body presented to a flowing fluid, is a function of its thickness and the superficial inlet fluid velocity. The unidirectional, steady-state and fully developed pressure drop of airflow across Z1 (2.5-3.15mm)

aluminium foam structure was measured repeatedly (reassembling the experimental setup and flipping both sides of the foam structure) at different air superficial velocity typically between 0.6 and 6.63m.s^{-1} (14–140LPM) and was divided by its porous body thickness (L=34.5mm) to determine the pressure drop per unit length ($\Delta P/L$) presented in Fig 7. Also, a reduced pressure drops ($\Delta P/LV$) was determined by diving the pressure drop per unit length against superficial air inlet velocity presented in Fig 8.

For the range of fluid velocity (0.6-6.63 m.s^{-1}) studied, the relationship between the developed pressure drop gradient across the porous structure (Z1) and superficial fluid velocity can simply be described by the third-term cubic Forchheimer relation (Eq 4) to extract the viscous and inertial terms. Fig 8 shows a clear distinction between the range of fluid velocity which clearly follows 2nd and 3rd order Forchheimer relation. This shows that fluid velocity ranging between 0.6-2.36 m.s^{-1} (14-50LPM) obeys the 2nd order Forchheimer relation with a correlation of 99.33%. Hence, the viscous (Permeability) and inertial (Form drag) terms of this fluid flow behaviour can easily be determined using the 2nd order quadratic model of Hazen-Dupuit-Darcy (Eq 1) fit of a pressure gradient to superficial velocity. Above this velocity range, a third order Forchheimer relation is observed for a correlation of 99.43%. A detailed understanding of the flow behaviour and the regime of occurrence is done using Darcy-Weisbach approach.
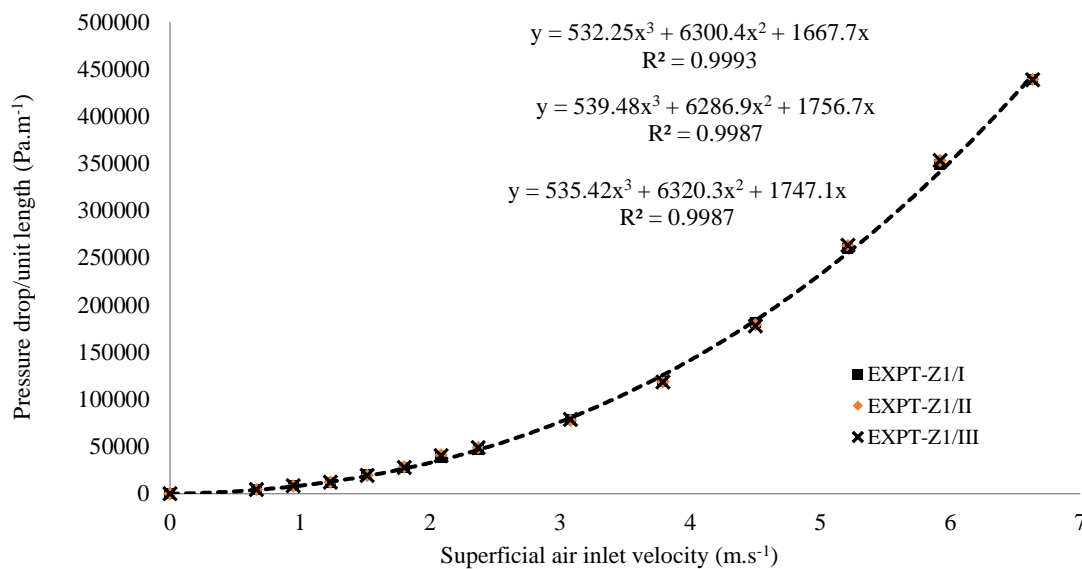

$$y = 532.25x^3 + 6300.4x^2 + 1667.7x$$
$$R^2 = 0.9993$$

$$y = 539.48x^3 + 6286.9x^2 + 1756.7x$$
$$R^2 = 0.9987$$

$$y = 535.42x^3 + 6320.3x^2 + 1747.1x$$
$$R^2 = 0.9987$$

Fig 7. Experimental pressure drops per unit length (Pa.m^{-1}) against air superficial inlet velocity (m.s^{-1}) for a repetitive test on aluminium foam sample Z1 (2.5-3.15mm).
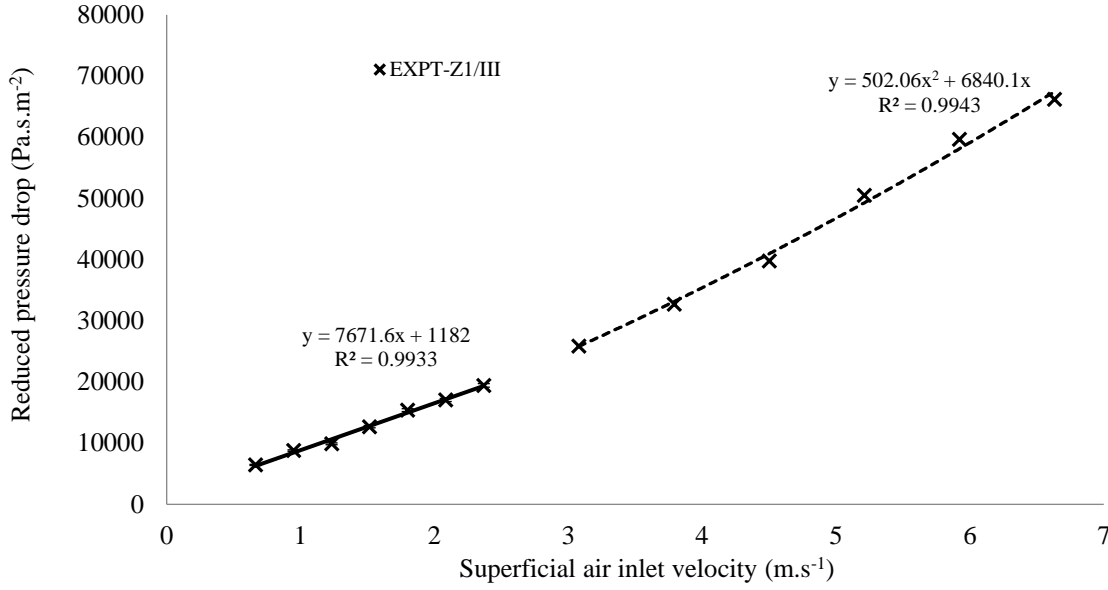
Fig 8. Experimentally reduced pressure drops (Pa.s.m^{-2}) against air superficial inlet velocity (m.s^{-1}) for a repetitive test on aluminium foam sample Z1 (2.5-3.15mm).

Reynolds Number and Friction Factor Relation

In practically all cases, the term Darcy's law for flow through porous structures is valid providing its mean pore/particle diameter-based Reynolds number (Re_D) does not exceed some value between 1 and 10 [30]. Lage et al. (2002), specified this regime to be a transition from Darcy to Forchheimer whilst complete Darcy regime was reported to exist at $0 \leq Re_D \leq 1.0$. Similarly, complete laminar, laminar-turbulent transition and fully developed turbulent regimes were reported for Reynolds number between 10-150, 150-300 and > 300 respectively. A more simplified approach to understanding the pattern of flow and regime of manifestation is the use of Darcy-Weisbach approach usually done by plotting a dimensionless Fanning friction factor (f_F) against mean pore diameter-based Reynolds number [21, 30] as shown in Fig 9 for Z1 structure, typically with the largest cell sizes and highest differential pressure. Optical images of this structure revealed a mean pore size of 2.77mm measured and was taken as the characteristic linear dimension (D_P) of the "bottleneck-type" structure as used in [29]. Mathematical representations of the Fanning friction factor (f_F) and pore-diameter based Reynolds number (Re$_D$) are presented in Eq 5 and Eq 6 as a function of pressure drop (ΔP), superficial velocity (V), mean pore diameter or characteristic linear dimension (Dp), fluid density (ρ) and fluid dynamic viscosity (μ).

$$f_F = \frac{\Delta P}{\left(\rho V^2 / 2\right)} \qquad \text{Eq 5}$$

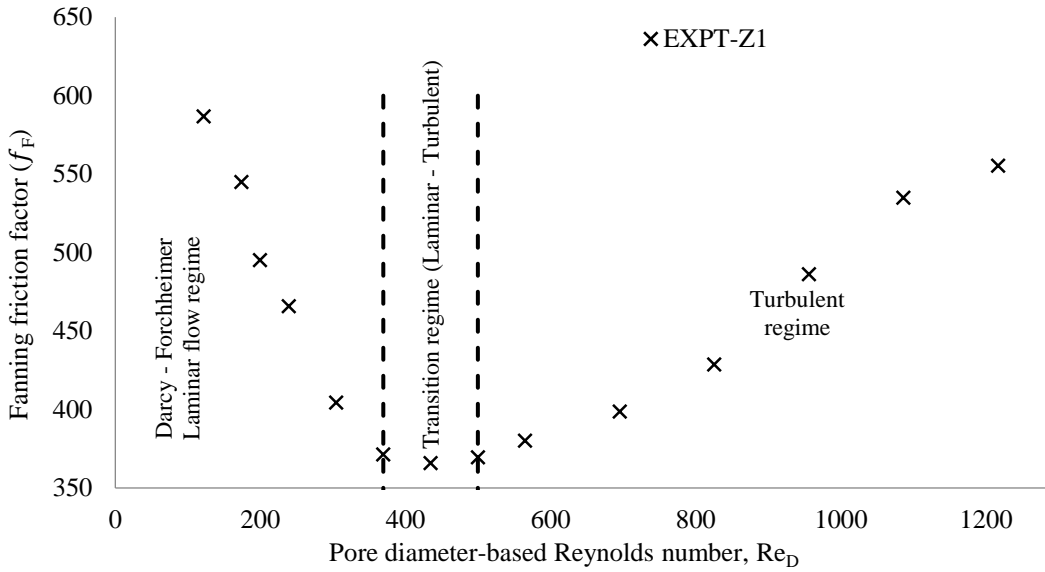$$Re_D = \frac{\rho V D_P}{\mu} \qquad\qquad\qquad \text{Eq 6}$$

Fig 9. Plots of dimensionless Fanning friction factor (*f*) against pore diameter-based Reynolds number (Re_D) for airflow through the lowest window-sized (Z1) porous Al foam structure.

Notable observations from the plot are the range of values estimated for the Reynolds number, typically to within 122-1221 for the specified velocity range (0.6-6.6 m.s^{-1}) and the regimes of manifestation with increased fluid velocity in the porous structure. Below 435 value of Reynolds number ($v_s \leq 2.36$m/s), there exists a power law relationship between the Fanning friction factor and Reynolds number. Viscous forces predominate inertial forces at the lower end of this laminar flow regime whilst inertial forces predominate viscous forces at the upper end of this regime. This regime is termed as the non-linear laminar flow regime and can be simply be described by the 2nd order Forchheimer relation (Eq 1). The permeability (k_0), Form drag (*C*) and Forchheimer coefficient (C_F) determined in this regime are 18.98x10^{-09}m^2, 6474.63m^{-1} and 0.89 respectively.

A significant contribution of wall function (increased Form drag) deviates its behaviour away from power law and this regime is termed as the transition regime. In this case, the transition from laminar to turbulent flow exists for velocity and Reynolds number range of $2.36 \leq v_s \leq 2.72$m/s and $435 \leq Re_D \leq 494$ respectively. As Re_D increases beyond 494, the turbulent flow is observed (Darcy's law not valid) and inertial forces completely dominate the fluid flow across the porous structure. This is an indication that high flow velocity of the fluid through a porous media is a major contributory factor in attaining turbulence [35] while the non-linearity deviation from Darcy is caused by inertial effects [8, 27] and not turbulence even though the flow is laminar for $Re_D \leq 435$.

Other foam structures conformed to having a similar trend of the Darcy-Weisbach approach, but their Reynolds number and friction factor differ from each other due to the variation in the structural parameters resulting from the use porogens of different sizes and the important applied differential pressure used during the casting process. A typical example of this is the flow behaviour through foam structure X1 made with packed porogens with the least particle sizes in the range of 1.0-1.4mm and at 0.9 bar liquid foam infiltration pressure. Darcy-Forchheimer, transition regime and turbulent regime for flow through this structure were attained for pore diameter-based Reynolds number, typically in the range of $48 \leq Re_D \leq 152$, $152 \leq Re_D \leq 173$ and $173 \leq Re_D \leq 485$ respectively (Fig 10). Similarly, the ranges of Reynolds numbers attained for these regimes for flow through the Y4 structure with the largest openings are $97 \leq Re_D \leq 453$, $453 \leq Re_D \leq 663$ and $663 \leq Re_D \leq 977$ presented in Fig 11.
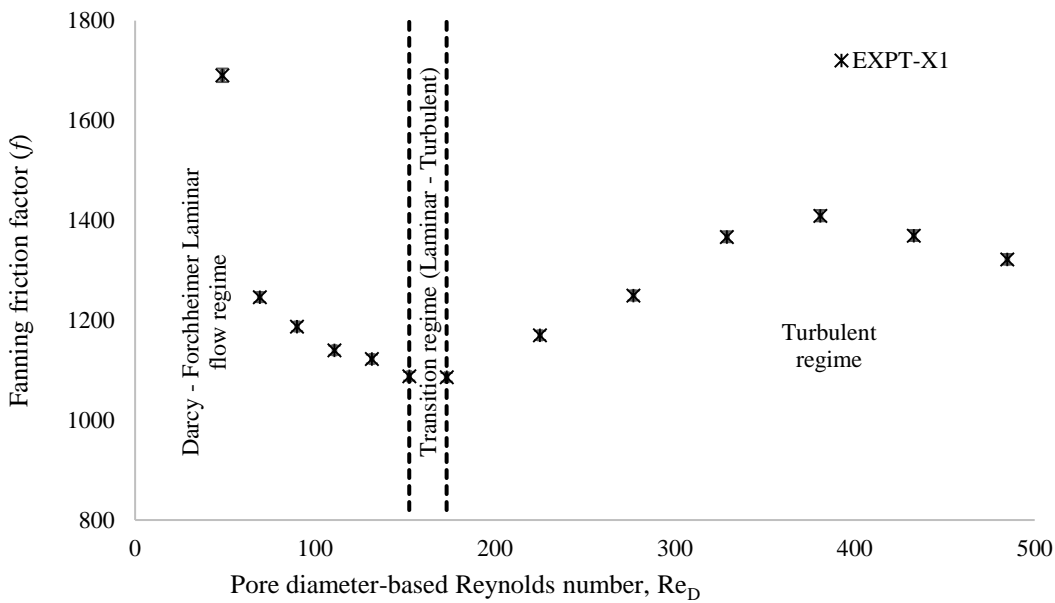

Fig 10. Plots of dimensionless Fanning friction factor (f_F) against pore diameter-based Reynolds number (Re_D) for airflow through the lowest window-sized (X1) porous Al foam structure.
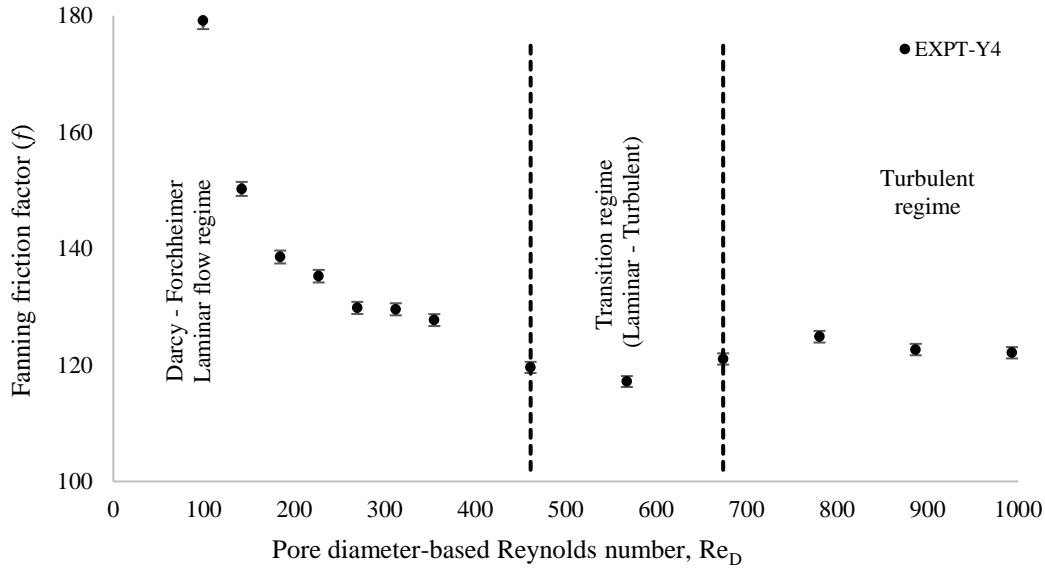
Fig 11. Plots of dimensionless Fanning friction factor (f_F) against pore diameter-based Reynolds number (Re_D) for airflow through the largest window-sized (Y4) porous Al foam structure.

Understanding this flow pattern is necessary towards the determination of viscous (permeability) and inertial (Form drag and Forchheimer coefficients) of fluid flow through porous structures of any topology and cell sizes. Though, the square root of permeability ($\sqrt{k_0}$) has been reported in the literature [36-38] to be useful in defining the true characteristic linear dimension of the structure due to fact that the permeability and not cell size represents more the inherent morphology of the porous structure. A plot of the Fanning friction factor against the permeability-based Reynolds number may indicate low values of Reynolds number but the behavioural trends from Laminar to Turbulent regime remain the same irrespective of the choice of characteristic linear dimension chosen.

Pressure Drop Data

A plot of the unit pressure drop (Pa.m^{-1}) measured for the six porous structures for the different range of pore sizes and differential pressures against the superficial air inlet velocity (m.s^{-1}) is presented in Fig 12. This polynomial trend can be fully described by the 2nd order Darcy-Dupuit-Forchheimer model with flow regime particularly to be within Laminar using Fig 7-11. Dividing the pressure drop per unit structural thickness measured for all the six foam samples by the superficial velocity (reduced pressure drop) and plotting against the superficial air inlet velocity as reported in [8, 15], the laminar flow unit pressure drop (Fig 12) can be further substantiated and to ensure the accurate values of the viscous and inertial terms are obtained. Such a plot is presented in Fig 13 for all the structures with a linear dependence fit (R^2 > 99%) indicating the measured pressure drops to be within the Forchheimer

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regime except for the X1 structure with the lowest cell sizes. Flow through this X1 structure indicate the presence of Darcy regime for the first two velocities with little or insignificant difference between their reduced pressure drops and Forchheimer regime exist for flow between 0.9-2.36m.s^{-1} with a linear dependence of very close fit.

The difference and order of magnitude in the measured unit pressure drop for all the foam structures can be classified into low, mid and high-pressure drop. Higher values of pressure drops are observed for samples with the lowest porosities (X1 and Y1) whilst the lowest pressure drop is attained for the sample (Y4) with largest openings and conferred to have the highest porosity value (Table 1). The similarity in the intermediate pressure drop attained for the Y2, Y3, and Z1 structure is greatly influenced by the minimal difference in the changes associated with porosity and openings between the connecting cell sizes within the microstructures.
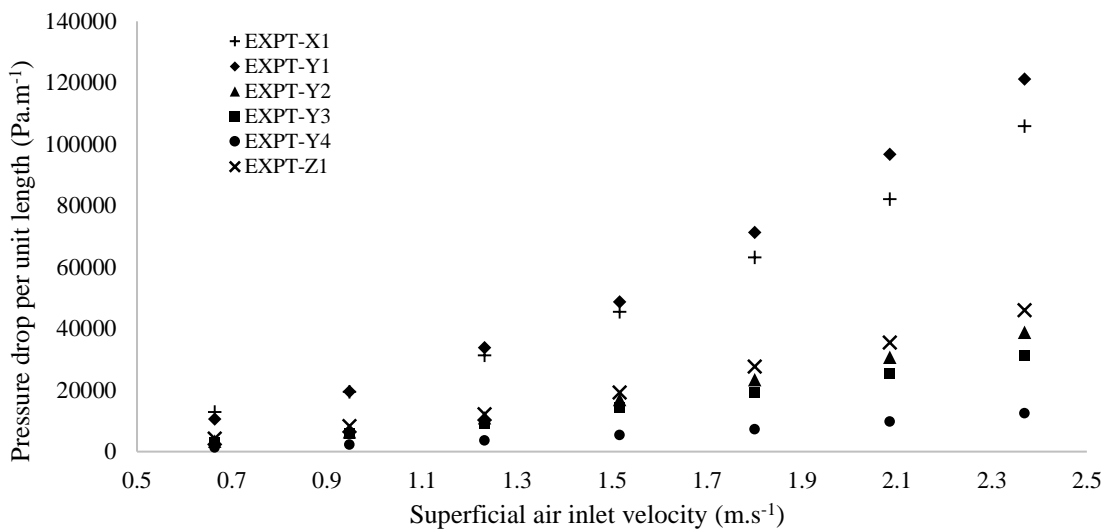

Fig 12. Plots of measured pressure drop per unit length (Pa.m^{-1}) against superficial air inlet velocity (m.s^{-1}) six foam structures.
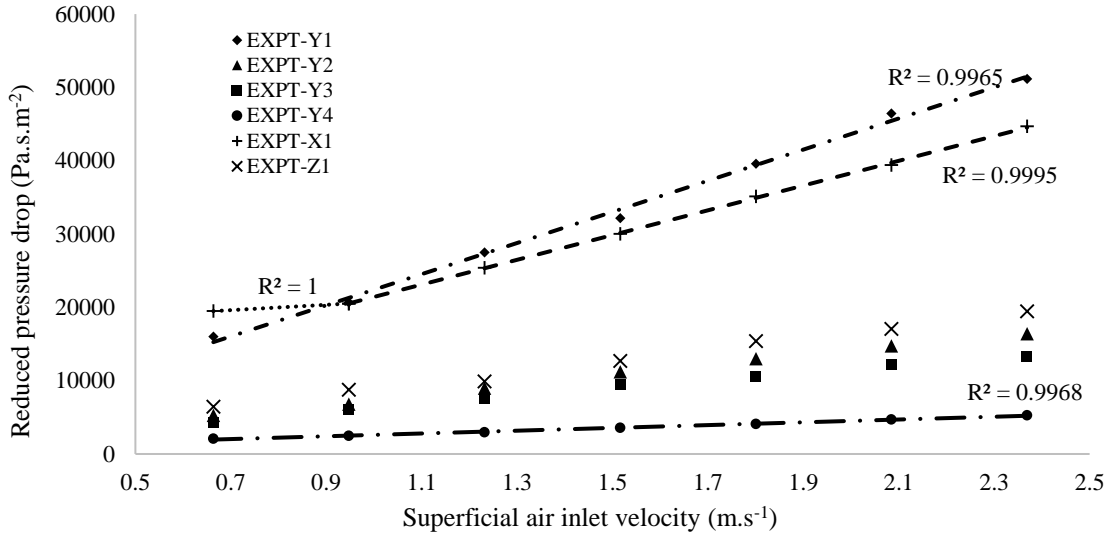
Fig 13. Plots of reduced pressure drop per unit length (Pa.m^{-1}) against superficial air inlet velocity (m.s^{-1}) six foam structures.

The size of the pores within the porous structure influences the tortuous path (tortuosity) and interstices available as pore volume for the flow of fluid (air). Larger pore sizes (opposed to smaller pore sizes) yielded a low tortuous path (tortuosity), low specific surface and large fluid volume which invariably resulted in low-pressure drops and high permeability of fluid measured across the foam structures. This flow trend is observed for foam X1, Y1 and Z1 structures (produced using similar liquid pressure differential of 0.9 bar) at very low velocity were the flow behaviour of the structure are largely dependent on the openings of the structures [36]. At high velocity greater than 1m.s^{-1}, the high pressure drops attained for Y1 structure when compared to the X1 structure are largely due to thickened ligaments resulting in the low porosity value of the Y1 structure and thereby providing more resistance to the flowing fluid across the structure. Optical imaging of these structures (Fig 1) revealed a near-spherical structure for the Y1 structure whilst the X1 structure consists of highly irregular-shaped, inhomogeneous cells and large "window" to pore ratio (resulting from more half sized spherical salts used during casting), thereby providing lesser resistance to flowing fluid at high velocities.

The behaviour of airflow through this structure is also affected by the number of "windows" (coordination number) existing with the porous matrix, typically, flow through Y structures made for different liquid differential pressures. At high infiltration pressures, liquid foams (aluminium) are forced into spaces created by packed beds of porogens thereby reducing the size and the number of "windows" available for permeation of fluid. Reported in [22], for samples having similar pore sizes and differential pressures, a plot of the ratio of "window" to pore size against the liquid pressure differential indicates a power inverse relation between these parameters. It is evident from this Table

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1 that at low differential pressure, the connectivity (openings) between connecting pores within the porous structure is larger than that produced at high infiltration pressure provided the similar pore sizes are maintained. Lower pressure drops per unit foam thickness and the largest "window" sizes are associated with the lowest infiltrated pressure Y1 sample (high capillary interaction between connecting pores) whilst a reverse order of macroscopic and flow data was obtained for the highest liquid infiltrated pressure Y1 structure (low capillary interaction between connecting pores. Tabular representation detailing values of the structural and flow parameters of the all the six porous structures studied herein are presented in Table 1.

Conclusion

Structural characterization and measurements of airflow across porous metallic structures typically, of the bottleneck-type thereof, have been presented. The permeability of the structures is said to greatly depend on the sizes of the "bottleneck" between connecting pores, typically, higher for larger apertures and lower for smaller apertures. The inertial term increases for the low-density structures and decreases for high-density structures having wider openings. Lower pore sizes were reported to give low permeability with greater resistance to flowing fluid, else, for the larger pore size structures. Computational fluid dynamics (CFD) modelling and simulation of laminar airflow across "virtually-created" structures imitating the fluid-structural domain of this "real" structures and an extension to bimodal structures could be considered to ease rapid understanding of flow behaviour of this material at low operating cost and flexibility resulting from changes pore diameter sizes and openings.

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