Title: Closed-form approximations to the Error and Complementary Error Functions and their applications in atmospheric science Short title: Closed-form of Error and Complementary Error Functions Authors: C. Ren and A. R. MacKenzie Affiliation: Department of Environmental Sciences Lancaster University, Lancaster LA1 4YQ United Kingdom Corresponding author: C. Ren Address: Department of Environmental Sciences Lancaster University, Lancaster LA1 4YQ United Kingdom E-mail: c.ren@lancaster.ac.uk. Tel: +44-1524-593974 Fax: +44-1524-593985

#### Abstract

The Error function, and related functions, occurs in theoretical aspects of many parts of atmospheric science. This note presents a closed-form approximation for the error, complementary error, and scaled complementary error functions, with maximum relative errors within 0.8%. Unlike other approximate solutions, this single equation gives answers within the stated accuracy for  $x \in [0 \ \infty)$ . The approximation is very useful in solving atmospheric science problems by providing analytical solutions. Examples of the utility of the approximations are: the computation of cirrus cloud physics inside a general circulation model, the cumulative distribution functions of normal and log-normal distributions, and the recurrence period for risk assessment.

#### Keyword

Error function, complementary error function, scaled complementary error function, normal distribution, log-normal distribution, cumulative distribution function, recurrence interval

## Introduction

Error and complementary error functions are extensively used in the fields that employ mathematics and physics, e.g., studies of heat and mass transfer (e.g., Chaudhry and Zubair, 1993; Swartzendruber, 2002). In atmospheric science, as elsewhere, the error and complementary error functions occur when normal or lognormal distributions are expressed as cumulative distribution functions. This note presents a close-form approximation for the error, complementary error, and scaled complementary error functions with maximum relative errors within 0.8%. The benefits of using an analytical approximation for error function in atmospheric sciences are demonstrated in some examples.

The closed-form approximation of error functions

The error function is defined as

$$erf(x) = \frac{2}{\sqrt{\pi}} \int_0^x e^{-t^2} dt$$
, (1)

and the complementary error function is defined as

$$erfc(x) \equiv 1 - erf(x) = \frac{2}{\sqrt{\pi}} \int_{x}^{\infty} e^{-t^{2}} dt \quad .$$
<sup>(2)</sup>

Both functions contain integrals. Sometimes, when one wants to evaluate these functions as accurately as possible, rational Chebyshev approximations (Cody, 1969) can be used. Nowadays, built-in functions are available in several computer languages (Cody, 1990). At other times, however, one may forego some accuracy for the sake of a speedy calculation or in order to gain a clearer insight into the relationships between variables in a problem. Decker (1975) provides and cites approximations that are quick to compute, but all fail to give expressions in closed form. The following achieves closed-form approximations for the error, complementary error, and scaled complementary error functions with maximum relative errors within 0.8%.

The complementary error function can be expanded as

$$erfc(x) = 1 - \frac{2}{\sqrt{\pi}} e^{-x^2} x \sum_{i=0}^{n} \frac{(2x^2)^i}{1 \cdot 3 \cdots (2i+1)}, \quad x < c ,$$
(3)

And

$$erfc(x) = \frac{1}{\sqrt{\pi}} e^{-x^2} \frac{1}{x} \left\{ 1 + \sum_{i=1}^n \frac{(-1)^i \cdot 1 \cdot 3 \cdots (2i-1)}{(2x^2)^i} \right\}, \quad x > c ,$$
(4)

where c is a positive real number with a value around 1. c increases with the increasing n, and Eq. (3) and (4) diverge when x is near to c. However, from Eq. (3), we know

$$e^{x^2} erfc(x) \approx 1 - \frac{2\sqrt{\pi x^2}}{\pi}, \quad x << 1.$$
 (5)

Similarly, from Eq. (4),

$$e^{x^2} erfc(x) \approx \frac{1}{\sqrt{\pi x^2}}, \quad x >> 1$$
 (6)

Both Eq. (5) and (6) suggest that

$$f(x) = \frac{a}{(a-1)\sqrt{\pi x^2} + \sqrt{\pi x^2 + a^2}}$$
(7)

might be a good fit to the scaled complementary error function  $e^{x^2} erfc(x)$ , since f(0)=1 and  $\lim_{x\to\infty} f(x)=\frac{1}{\sqrt{\pi x^2}}$ . Here, a is an adjustable parameter for (7) to match either (5) or (6). When  $\sqrt{\pi x^2 + a^2} = a$  is used for x << 1, the simple match of Eq. (7) to Eq. (5) requires  $a = \frac{\pi}{\pi - 2}$ .

We derive a series of values for a by trial-and-error. These values are given in Table 1, together with their accuracies when used in Eq. (7) to calculate  $e^{x^2} erfc(x)$ .

The errors introduced by using Eq. (7) to estimate erf(x) and erfc(x) are shown in Fig. 1. One can choose a value of a to evaluate the results of the error, complementary error, and scaled complementary error functions easily with a calculator to within the accuracy shown.

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а	$a^2$	Error in	Error in	Best choice for			
		erf(x)	erfc(x)				
π	7.5732	-0.65-0.00%	0.00-0.92%	errors with			
$\pi - 2$				known sign			
2.7749	7.7000	-0.47-0.47%	-0.01-0.82%	erf(x)			
2.7889	7.7780	-0.38-0.75%	-0.03-0.76%	single value of $a$ for			

Table 1 Values of a, errors in erf(x) and erfc(x), and advantages of each formulation

				both $erf(x)$	and	erfc(x)
2.9110	8.4740	-0.04-3.11%	-0.34-0.34%	erfc(x)		
3	9	0.00-4.70%	-0.65-0.12%	simple for	mula	

Examples of applications in atmospheric science

As indicated in Table 1, using a=3 results in the simplest formula:

$$e^{x^2} erfc(x) \approx \frac{3}{2\sqrt{\pi x^2} + \sqrt{\pi x^2 + 9}}$$
 (8)

which is never in error by more than 0.65% for complementary error and scaled complementary error functions, and which provides a neat closed solution that can be incorporated into analytical solutions for a broad range of physical and engineering problems. An example of one such problem, and the driver for the development discussed above, is the nucleation and growth of cirrus cloud particles (Ren and MacKenzie, 2005). The approximation allowed us to describe the behaviour of cirrus clouds under all conditions, avoiding an unwieldy and unhelpful description based on asymptotic expansions to both ends when it is, in fact, the middle range that is most interesting (Kärcher and Lohmann, 2002; Ren and MacKenzie, 2005). The cirrus parameterisations are designed for implementation in global climate models (GCMs) (Lohmann et al., 2004), where the error associated with the closed-form approximation is small compared to uncertainty in the model output resulting from missing processes and other simplifications. There are clear advantages — in calculation speed and interpretation of results — in the use of closed-form approximations to the error and related functions within these very large GCM computer codes.

Other examples relate to normal or log-normal distributions. The size distributions of aerosols and clouds, and the parameters of turbulent processes are often log-normally distributed. Shoji and Kitaura (2006), for example, found that hourly, daily, and annual precipitation distributions were fitted well with log-normal distributions. The cumulative distribution function, D(I) - which indicates the probability that rainfall amount, I, will not be exceeded within period of time (hourly, daily, or annual), T - can, therefore, be given by

$$D(I) = \frac{1}{2} \left( 1 + erf\left(\frac{\ln I - \ln I_m}{\sqrt{2}\ln\sigma}\right) \right) = 1 - \frac{1}{2} \frac{ae^{-\lambda^2}}{(a-1)\sqrt{\pi\lambda^2} + \sqrt{\pi\lambda^2 + a^2}},$$
 (9)

where  $I_m$  is the geometric mean of rainfall amount,  $\sigma$  is the geometric standard deviation of rainfall amounts, and  $\lambda = \frac{\ln I - \ln I_m}{\sqrt{2} \ln \sigma}$  is a convenient measure of the position of a particular rainfall amount in the rainfall distribution.  $\lambda = 1/\sqrt{2}$  when  $I = \sigma I_m$ ,  $\lambda = 2/\sqrt{2}$  when  $I = \sigma^2 I_m$ , and so on. You can calculate D(I) with a given  $\lambda$  by a calculator, even by hand.

Having derived an analytical expression for the cumulative distribution function, the recurrence interval is then

$$R(I) = \frac{T}{1 - D(I)} = \frac{2T}{a} \left[ (a - 1)\sqrt{\pi\lambda^2} + \sqrt{\pi\lambda^2 + a^2} \right] e^{\lambda^2} .$$
(10)

This relatively simple expression is much easier to "read" than the equivalent retaining the error function. For instance, for  $I = I_m$ ,  $\lambda = 0$  and so  $R(I_m) = 2T$ , which confirms that there is a 50:50 chance that rainfall exceeds the geometric mean,  $I_m$ . For  $\lambda = 1/\sqrt{2}$ ,  $R(\sigma I_m) \approx 6.3T$ ; for  $\lambda = 2/\sqrt{2}$ ,  $R(\sigma^2 I_m) \approx 43.6T$ . Using Eq. (10), values of  $\lambda$  for the 50-, 100-, and 200-year events are given in Table 2. This time, a = 2.7889 is used, as this value of a guarantees the approximation having relative errors within 0.8%. Beyond their intrinsic interest, log-normal rainfall statistics also propagate into hydrological theory – theoretical treatments of slope stability for example (Iida, 2004) – and engineering design, where, again, avoiding the use of error functions makes model building and theoretical interpretation easier (Swamee, 2002).

λ	R(I)			
	50yrs	100yrs	200yrs	
Hourly	3.2417	3.3427	3.4409	
Daily	2.7361	2.8533	2.9663	
Annual	1.4543	1.6468	1.8230	

Table 2 Values of  $\lambda$  for calculating the rainfall amount at a given recurrence interval

# Conclusion

A closed-form approximation for the error, complementary error, and scaled complementary error functions with maximum relative errors within 0.8%. Unlike other approximate solutions, this closed-form equation gives answers within the stated accuracy for  $x \in [0 \infty)$ . It is very useful when one wants to gain a clearer insight into the relationships between variables in a problem involving error functions.

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Figure 1 Relative errors of using Eq. (7) with  $a = \pi/(\pi - 2)$ , 2.7749, 2.7889, 2.9110, and 3 to calculate error functions. Each line is for a single value of a. For erf, a increases from bottom up; for erfc, a increases from top down. Errors in erf(x) with  $a = \pi/(\pi - 2)$  and in erfc(x) with a=3 are almost identical when x<1.