Tropopause and hygropause variability over the equatorial Indian Ocean
during February and March 1999

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Abstract

Measurements of temperature, water vapour, total water, ozone, and cloud properties were made above
the western equatorial Indian Ocean in February and March 1999. The cold-point tropopause was at a mean
pressure-altitude of 17 km, equivalent to a potential temperature of 380 K, and had a mean temperature of 190 K. Total water mixing ratios at the hygropause varied between 1.4 and 4.1 ppmv. The mean saturation water vapour mixing ratio at the cold point was 3.0 ppmv. This does not accurately represent the mean of the measured total water mixing ratios because the air was unsaturated at the cold point for about 40% of the measurements. As well as unsaturation at the cold point, saturation was observed above the cold point on almost 30% of the profiles. In such profiles the air was saturated with respect to water ice but was free of clouds (i.e., backscatter ratio < 2) at potential temperatures more than 5 K above the tropopause and hygropause. Individual profiles show a great deal of variability in the potential temperatures of the cold point and hygropause. We attribute this to short time- and space-scale perturbations superimposed on the seasonal cycle. There is neither a clear and consistent “setting” of the tropopause and hygropause to the same altitude by dehydration processes nor a clear and consistent separation of tropopause and hygropause by the Brewer-Dobson circulation. Similarly, neither the tropopause nor the hygropause provide a location where conditions consistently approach those implied by a simple “tropopause freeze drying” or “stratospheric fountain” hypothesis.

**Introduction**

The stratosphere can be considered to consist of two, very different, regions: the “overworld”, with a lower boundary given by the first potential temperature, θ, surfaces to lie entirely within the stratosphere (ca. 380 K), and the “middleworld” or “lowermost stratosphere”, which lies between the overworld and the tropopause [Holton et al., 1995]. The lowermost stratosphere experiences considerable interchange of air with the troposphere; this interchange being driven by synoptic-scale eddies and folds [e.g. Dethof et al., 2000]. In contrast, the movement of air through the stratospheric overworld is driven by a slow meridional motion, the Brewer-Dobson circulation. Air enters the overworld primarily in the tropics [Brewer, 1949]. This air has well-defined concentrations of trace gases with long tropospheric lifetimes — the CFCs, for example — but has much less well-defined concentrations of compounds with short tropospheric (chemical or physical) lifetimes, including water vapour. Water vapour is important in stratospheric chemistry as a
source of HO₃ radicals and as the major condensate in polar stratospheric clouds, and it plays an essential role in the stratospheric radiation budget [e.g. Dvortsov and Solomon, 2001; Forster and Shine, 1999].

The observed increase of water vapour in the stratosphere over the last decades cannot be explained quantitatively by the atmospheric increase of methane alone [e.g. Rosenlof et al., 2001]. Considerable drying takes place as air moves from troposphere to stratosphere in the tropics. The only plausible mechanism for such drying is “freeze-drying”: the cooling of a cloud-laden air parcel to the temperature at which the local saturation mixing ratio is less than or equal to the observed stratospheric water vapour mixing ratio (to be precise, that part of the observed stratospheric mixing ratio not due to methane oxidation). The excess water, condensed in the cloud elements, is then assumed to sediment from the air parcel in the form of ice crystals.

Several hypotheses have been put forward to explain where and how drying of air entering the stratosphere occurs. The ‘stratospheric fountain’ hypothesis of Newell and Gould-Stewart [1981] suggests that (a) cross-tropopause transport occurs predominantly over Indonesia and the maritime continent in northern winter, and over northern India in northern summer, and (b) large-scale layers of sub-visible cirrus occur over these regions at these times. The detection of large-scale sub-visible cirrus layers in the tropics [Wang et al., 1996] and the results of trajectory-based studies [Fueglistaler et al., 2004; Bonazzola and Haynes, 2004] have rekindled interest in versions of this hypothesis. Holton and Gettelman [2001] suggested that horizontal transport through this cold trap region brings very dry air to other longitudes where it can finally be transported into the stratosphere. Recent trajectory studies using 21 years of reanalyses from the ERA-40 data of the European Centre for Medium-Range Weather Forecasting (ECMWF) have shown that fixing the final water vapour mixing ratio in a trajectory to that given by the saturation mixing ratio at the coldest temperature along the trajectory gives a satisfactory fit to long-term average stratospheric humidity and to the seasonal cycle [Fueglistaler et al., 2005].

Danielsen [1982] proposed an alternative based on the dynamics of individual convective clouds: that tropical convective clouds penetrate the tropopause. Radiatively-driven overturning of the thick cirrus anvils from the convection then allows cloud particles to grow sufficiently that they sediment out, thus drying the
air. There have been a few observations of this mechanism apparently at work [Danielsen, 1993]. It is not clear if there is sufficient tropopause-penetrating convection to allow the Danielsen mechanism alone to dry all the air moving into the stratospheric overworld, as forced by the Brewer-Dobson circulation. Convective cloud at the tropical tropopause implies air masses that are very different from their surroundings — i.e., with trace gas and total water signatures indicative of rapid diabatic transport (convection) rather than cloud formation in situ. Such signatures would include low mixing ratios of ozone, high mixing ratios of total water, and localised ice supersaturation within the cloudy air mass (since cloud formation took place at lower altitudes).

A third hypothesis requires active convection, but not tropopause penetration. Potter and Holton [1995] postulate convectively-induced gravity waves in the lower stratosphere as a source of adiabatic cooling, cloud formation, and, thence, dehydration. Another possible source of adiabatic cooling is from planetary wave activity, such as Kelvin waves [Jensen et al., 1996]. In situ cloud formation at the tropical tropopause implies air masses that are similar to their surroundings — i.e., with trace gas and total water signatures indicative of adiabatic rather than diabatic transport. Such signatures would include average mixing ratios of ozone, average mixing ratios of total water, and ice supersaturation that is not restricted to the cloudy air mass (since cloud formation in the upper troposphere appears to require substantial supersaturation).

A fourth hypothesis [Sherwood and Dessler, 2000] combines the overshooting dehydration of Danielsen [1993] with slow upward motion through a thick tropopause layer. The existence of a tropopause layer several kilometres thick has been suggested at various times in the literature (see the review in Highwood and Hoskins [1998]). The tropopause layer is sometimes referred to as the sub-stratosphere, since it is above the mean level of tropospheric convection [Thuburn and Craig, 2000]. It is also known as the tropical tropopause transition layer (TTL), a region that is not efficiently flushed by convection and that, for the most part, is subject to radiative heating [Folkins et al., 1999].

Analysis of data obtained at different tropical stations and seasons show, that the relative contribution of the aforementioned dehydration processes varies with longitude and season [Vömel et al., 2002]. However in
that study, deep convection was found to be important only in setting up the tropical tropopause layer which is then subject to large-scale wave activity and wave breaking at the tropopause. In general, progress in distinguishing the effectiveness of the different dehydration mechanisms has been slow, due to the sparseness of observations, particularly coincident observations of temperature, humidity, water vapour, condensed water, and gas-phase tracers of atmospheric transport.

During February and March 1999, the European APE-THESEO mission took place, based on Mahé on the Seychelles (4° 42’ S, 55° 30’ E) in the Indian Ocean, where the UTLS region had not been the target of a major scientific campaign involving aircraft or balloons before. For this project, the Russian high-altitude research aircraft Geophysica was equipped with a comprehensive in-situ payload, accompanied by the DLR Falcon carrying a lidar and radiometers, which acted as a pathfinder for cloud studies with the Geophysica. Flight paths consisted of a combination of long-range transects, nominally directed normal to the inter-tropical convergence zone (ITCZ), and more complicated interceptions of cloud systems.

APE-THESEO stands for: Airborne Platform for Earth Observation – Third European Stratospheric Experiment on Ozone. APE-THESEO aimed to study the transport, chemistry, and cloud physics of the tropical tropopause region [MacKenzie et al., 2000; Stefanutti, et al., 2004], particularly with respect to factors affecting the concentrations of trace gases in the stratosphere. Below, we use the data collected during APE-THESEO, along with radiosonde and meteorological analyses, to describe the structure and variability of the tropopause over the equatorial Indian Ocean during northern winter/spring 1999. A summary of the characteristics of all the ascents and descents across the tropopause, made during the campaign, is given first. Individual profiles are then discussed in detail, in the light of the hypotheses above. Finally, the mean water vapour profile is compared to mean profiles from previous aircraft missions and discussed in terms of the tropical “tape recorder” paradigm. A previous analysis of the total water data used below can be found in Beuermann [2000].
Instruments and Data

The following analysis uses in-situ data of water (total and vapour), ozone and particles, measured from the *Geophysica* aircraft [Stefanutti et al., 1999] over the Indian Ocean in February and March 1999. We also use routine measurements of temperature, pressure, position, etc., from the aircraft sensors.

The Fast In situ Stratospheric Hygrometer (FISH), developed at the Forschungszentrum Jülich (Germany), is based on the Lyman-α photofragment fluorescence technique. It has flown previously on other aircraft [Zöger et al., 1999]. The overall accuracy of this hygrometer is 6%, or 0.3 ppmv in the case of the very low mixing ratios that occur in the tropics. FISH was shown to yield results that are consistent with other stratospheric water vapour measurements [Kley et al., 2000]. In the presence of clouds, FISH measures total water, with an over-sampling of cloud elements [Schiller et al., 1999]. For typical Geophysica cruising altitude and speed, the over-sampling factor for particles with radii larger than 4 μm is 5. Thus FISH measurements are very sensitive to cloud water content.

The FLuorescent Airborne Stratospheric Hygrometer FLASH instrument developed by the Central Aerological Observatory (Moscow) is an aircraft version of the water vapour instrument that has previously been deployed on balloon [Merkulov and Yushkov, 1999], and also uses the Lyman-α fluorescence technique. During APE-THESEO, the inlet was designed to measure gas-phase water providing complementary information to the FISH total water measurement inside clouds. Out of cloud, FISH and FLASH gave generally consistent data. Since FISH was calibrated between all flights during APE-THESEO, and FLASH only once before the campaign, for this study FLASH was recalibrated against FISH during out-of-cloud measurements.

Temperature onboard the Geophysica aircraft during the APE-THESEO campaign was measured by the on-board system whose accuracy is specified to ± 0.5 K (B. Lepouchov, Myasishchev Design Bureau, Russia, private communication, 2000). Comparison with independent temperature probes (Rosemount TDC probe, microwave temperature profiler) available on this aircraft during following projects confirmed this accuracy. For the RHi calculated in the discussion chapter (in particular Figure 7), this uncertainty translates
in an uncertainty of approximately ± 15% RH, assuming 3 ppmv H₂O mixing ratio and conditions close to saturation.

The Electro-Chemical Ozone Cell ECOC is a modified electrochemical ozone sonde. Ozone measurements from ECOC have been validated against ozone sondes [Kyrö et al., 2000] and shown to yield a small negative bias of -5.7 ± 2.8 %. Such a bias is not significant in the context of the following analysis. FOZAN (Fast OZone ANalyzer) is a chemiluminescent ozone sensor [Yushkov et al., 1999] based on a solid-phase dye. FOZAN can detect fast ozone variations, but needs frequent calibration, and ECOC has been used as a reference instrument for the FOZAN. In flights where both ECOC and FOZAN were operating, ECOC and FOZAN showed generally close agreement, with an average correlation coefficient of 0.94 and no significant bias. Here we use ECOC data throughout except for the last scientific sortie (11 March 1999) when they are absent and FOZAN data are used in their place.

The Multiwavelength Aerosol laser Scatterometer MAS first flew on the Geophysica during the APE POLECAT mission [Adriani et al., 1999]. In daylight, MAS measures backscatter and depolarisation at 532 nm. The sun is required to be outside the field of view of the instrument (i.e., a solid angle of 20°), looking horizontally through a shutter on the starboard side of the aircraft. Since sorties were often conducted around sunrise, there were several occasions when the MAS shutter was closed to prevent direct sunlight entering the instrument.

Besides MAS there was other in-situ cloud and aerosol instrumentation on board the Geophysica (FSSP-300, Mini Copas, CVI). Descriptions can be found in Stefanutti et al. [2004]. We do not present data from these instruments here, but we do draw on the analysis of FSSP data given in Thomas et al. [2002]. The in-situ measurements from the Geophysica are complemented by lidar observations of clouds and aerosols from the Falcon aircraft. A recent description of the lidar is given in Flentje et al. [2002].
Results

During APE-THESEO, 11 flights of the Geophysica aircraft were carried out in the tropical region over the Indian Ocean providing 37 individual vertical profile observations. Details of the locations of the profiles are given in Table 1. The first of these profiles was obtained at 31 °N, and so, strictly, outside the tropics. Some of the flight legs included long periods of horizontal flight or slow climbs, so that some of the vertical sections are composed of data taken over a wide area (as short as 250 km and as long as 2000 km). The vertical profiles taken over a long duration did not show any systematic difference to those taken over a shorter duration. Nevertheless, these aircraft-derived vertical profiles should not be over-interpreted, because the air masses at different altitudes have only a limited common history as a result of vertical shear in the horizontal winds. This implies, inter alia, lower correlations between measures of TTL structure with increasing vertical separation.

We determined for all tropical ascents and descents of the campaign the cold-point tropopause, the lapse rate tropopause, and the highest pressure-altitudes, where 100 and 200 ppbv ozone thresholds were found, in order to provide insight into the vertical structure of the TTL. Ozone mixing ratios around 100 ppbv have been used in the extra-tropics to define the “chemical tropopause” [e.g., Bethan et al., 1996]. The variation of these various tropopause-like quantities is displayed in Figure 1. The analysis to derive the lapse-rate tropopause in the aircraft profiles is somewhat subjective, since the ascent has a substantial horizontal component that can produce a noisy lapse-rate profile in the presence of even modest horizontal temperature gradients [Danielsen, 1993]. A 90-second running average was used to smooth the data in the lapse rate calculation. The lapse-rate tropopause is generally lower than the cold-point tropopause, by up to 1 km for individual profiles, and by 500 m in the mean. The cold-point and the lapse rate tropopause heights are only weakly correlated ($r^2 = 0.53$). Reid and Gage [1996] found that the monthly mean lapse-rate tropopause at Truk (7.5 °N, 151.8 °E) for 1980 was 150 m below the monthly mean cold-point.

In the TTL, the ozone profile often contains a local maximum, or even several local maxima. The lowest altitude, at which an ozone mixing ratio of 100 ppbv was observed during APE-THESEO, varies from 13.3
to 17.6 km (mean: 16.1 km, or 373 K potential temperature). This altitude is usually located over a kilometre beneath the cold point tropopause and varies independently of the cold-point tropopause ($r^2 = 0.06$), as also found by Folkins et al. [1999]. The highest altitudes with an ozone mixing ratio of 100 ppbv are located closer to the cold-point tropopause (13.4 to 18.1 km, mean: 16.7 km, or 377 K potential temperature, Figure 1), but are also not correlated ($r^2 = 0.06$). In several profiles, even ozone mixing ratios of 200 ppbv were found below the cold point. These high ozone values indicate ongoing ozone production during the slow ascent of air in the TTL [Folkins et al. 1999] and/or transport from the stratosphere into the TTL as suggested by Tuck et al. [1997; 2004]. The highest altitude with a 200 ppbv ozone mixing ratio (16.8 - 18.7 km, mean: 17.7 km, or 398 K potential temperature) shows a relatively strong positive correlation with the cold-point tropopause ($r^2 = 0.64$). This suggests that the highest altitude at which a 200 ppbv ozone mixing ratio occurs is in the lower stratosphere, aligned with a potential temperature surface. The good correlation between cold point and the highest altitude with a 200 ppbv ozone mixing ratio implies that much of the observed variability of the tropopause height in Figure 1 is associated with reversible wave motion.

Figure 2 shows the potential temperature of the cold-point tropopause, the hygropause, and the range of potential temperatures for which each profile was saturated with respect to water ice. The tropopause and hygropause data are given with respect to pressure-altitude in Table 1. From Figures 1 and 2 and Table 1, we can deduce the following. In the deep tropics in the Indian Ocean, i.e., latitudes equatorward of 20º, the cold-point tropopauses observed vary in temperature and height from 194 K at 16.6 km to 182 K at 18.1 km. There was variation in the observed cold-point potential temperature from 365 K to 404 K. The mean of the cold-point temperatures is 188 K, at a mean potential temperature of 380 K. As will be discussed further below, the cold point and hygropause can be rather poorly defined sometimes, due to small vertical gradients in temperature and water vapour (or total water). However, taking the data in Figure 2 at face value, 11 profiles show the hygropause at a potential temperature more than 5 K above the tropopause, 5 show the hygropause more than 5 K below the tropopause, and 21 show the tropopause and hygropause within 5 K of each other. There is neither a clear and consistent setting of the tropopause and hygropause to the same
altitude by dehydration processes, nor a clear and consistent separation of tropopause and hygropause by the Brewer-Dobson circulation. Of the 37 profiles, 6 are saturated at the tropopause but not at the hygropause, none are saturated at the hygropause but not at the tropopause, and 13 are saturated at both the tropopause and hygropause. 11 profiles are saturated to potential temperatures more than 5 K above the hygropause. This variability of saturation conditions at the hygropause and tropopause underlines the complex processes occurring during the observation period. Neither the tropopause nor the hygropause provide a location where conditions consistently approach those implied by a simple "stratospheric fountain" hypothesis (i.e., a high frequency of cloud and saturation with subsaturation above and below), and where simple deductions of water vapour mixing ratios could be made from temperature soundings.

Figure 3 shows the total water as measured by FISH, saturation water vapour mixing ratio (with respect to ice) and relative humidity (for total water) at the hygropause for each ascent and descent. The hygropause total water mixing ratios, measured by FISH, range from 1.4 to 4.1 ppmv. The mean hygropause total water mixing ratios is 2.4 ppmv, at a mean pressure altitude of 17.2 km (383 K potential temperature).

Also shown in Table 1 are cloud-top pressure-altitudes for clouds above 14 km and a 532-nm backscatter ratio greater than 2, as derived from MAS measurements. This backscatter ratio threshold for cloudiness will miss ultra-thin cirrus clouds (UTTCs), which were frequently observed during APE-THESEO. The 532-nm backscatter ratio for UTTCs is about 1.2 only [Peter et al., 2003; Luo et al., 2003a, b]. Detection of such thin clouds is not usually possible from MAS data alone, and is outside the scope of the present discussion. The statistics below on cloud occurrence are, therefore, lower limits. 17 of the vertical profiles show clouds, 8 show no clouds, and 12 have no cloud data available. When clouds are present with backscatter ratio greater than 2, cloud-top pressure-altitudes range from 14 km to 18.1 km. 10 of the 17 vertical profiles with clouds have cloud-tops within 200 m of the cold point.

The distributions of values around the mean for most of the parameters in Table 1 are generally near-normal, given the small sample (n ≤ 37). The distributions of saturation mixing ratios are log-normal, as expected from the exponential dependence of the saturation mixing ratio on temperature.
Discussion

*Individual profiles: Clouds, water vapour, and meteorological conditions*

Examples of specific vertical profiles of water, ozone, clouds and temperature measured during APE-THESO are discussed in this section. They include two cases involving cirrus clouds whose characteristics, origin and impact on the water vapour budget are investigated. The third and fourth profiles are clear-sky profiles, demonstrating the vertical structure of the TTL over the Indian Ocean without the direct influence of clouds. Profiles are numbered as per Table 1.

**Profile # 29: 11 March 1999**

The first example has relatively thick (but still barely visible) cirrus clouds at the tropopause, and a relatively low cold-point potential temperature. Figure 4 shows the lidar cross section through this cloud along the flight path of the Falcon, which was obtained approximately 1 h before *in-situ* measurements from the Geophysica were made. The vertical extent of this cloud varied between 1 and 4 km with a cloud top around 17.2 km.

The cirrus cloud in Figure 4 was profiled by the Geophysica at different locations about 1 h after the Falcon measurements (the flight path is overlaid on the lidar curtain plot). The individual profiles, i.e., the vertical extension of the cirrus, are consistent in both measurements. As an example, Figure 5 shows a profile from 11 March 1999 (profile No 29 in Table 1, marked by an arrow in Figure 4). The hygropause is at 16.9 km ($\theta = 374$ K), and occurs at the top of a cloud, as indicated by the backscatter data. Consistent with the lidar data, the cloud at the hygropause is relatively thick geometrically and optically, indicative of a (barely) visible cirrus cloud: backscatter ratios, at 532 nm, often exceed a value of 10. The condensed water mixing ratios, calculated as the difference of total water measurement of FISH and the saturation mixing ratio, are up to 3.5 ppmv. Total water mixing ratios of 7-8 ppmv above 16 km are significantly higher than in clear-sky profiles during APE-THESO and thus a clear indicator for injection of water from lower altitudes. The cold-point tropopause is coincident with the hygropause. There is a stable and saturated layer, of several
hundred metres, above the cold-point. This layer is free of cloud. At the top of the saturated layer, the
gradient of the ozone profile increases significantly. In the core of the clouds, however, ozone values are only
50-70 ppbv and thus significantly lower than observed elsewhere in the TTL (see below).

The other profiles on 11 March, but also those of 9 March (12 profiles in all), have very similar
characteristics: barely visible cirrus of more than 1 km vertical extension, with ice water content up to
6 ppmv, and cloud top and hygropause close to the rather low tropopause. Ozone mixing ratios of 50 to
70 ppbv were measured in the clouds of these profiles which are significantly lower than their environment
and do not show the increase of ozone usually observed in the TTL [e.g. Folkins et al., 1999]. These air
masses are therefore likely transported rapidly from lower altitudes probably in convective systems [see also
Sherwood and Dessler, 2000].

Meteosat 5 images combined with wind analysis (Figure 6) imply that these clouds were likely to be the
direct result of convective-cloud outflow. Several hours before the airborne observations, cloud top
temperatures of less than 193 K (− 80°C) were observed east of the southern part of the flight track in several
extended convective cells. Cloud top temperatures below 193 K in these days, corresponding to a pressure
altitude of 100 hPa, were among the lowest observed in the APE-THESEO period over the Indian Ocean,
when convection usually did not exceed 16 km [Stefanutti et al., 2004]. Wind velocities at 100 hPa were
easterly and about 30 m s⁻¹ (UKMO output, not shown), and so allow the outflow from these systems to be
transported in 4 to 12 hours to the region of the flight track (as sketched in Figure 6). To estimate the
radiative effect of the cirrus cloud, the heating rate has been calculated, based on the actual observations of
O₃, H₂O, and cloud extinction and following the method of Corti et al. [2005]. Mean heating rates for the
cloud are from 12 to 18 K day⁻¹. With a vertical gradient in potential temperature of 10 K km⁻¹, the cloud
could be lifted by several hundred meters on its way from the convective area to the observation region.
Consistently, ECMWF analyses show rather high vertical transport velocities in the area of the cirrus cloud
areas, of around 2 cm s⁻¹. The cloud-radiation interaction at least compensates for cloud particle
sedimentation, which, for particles with radii of 15 to 20 \( \mu \text{m} \), is similar in magnitude to the rate of cloud lofting.

Figure 7 (solid lines) shows the frequency distribution of relative humidity over ice (RHi) for the profiles of the flight on 11 March. Outside clouds, RHi usually does not exceed 50 \%, peaking at 20 \%, and is thus not close to saturation conditions to form new cirrus. In the clouds themselves, the RHi distribution peaks at 100 \% with only a few observations of moderate supersaturation. This suggests that the clouds in Figure 4 and 5 are ‘aged’ to near-equilibrium conditions or already in the process of evaporation, which corroborates the finding that they originate directly from convective activity and were transported to the position of the observation.

Profile # 6: 19 February 1999

Figure 8 shows a profile measured on 19 February 1999. The hygropause is at 18.1 km (\( \Theta = 387 \text{ K} \)), and occurs towards the top of a cloud, as indicated by the backscatter data. The cloud at the hygropause is relatively thin: backscatter ratios are less than 10, and condensed water mixing ratios are below 3 ppmv. Santacesaria et al. [2003] have analysed this case, and calculated optical depths on or slightly above the threshold for visibility for different parts of the cloud field around the sounding shown in Figure 8. The cold-point tropopause is coincident with the hygropause. There is a saturated, near-adiabatic layer above the cold-point, with potential temperature increasing much more slowly between 18 km and 18.5 km than above or below. The backscatter data at this altitude are noisy, but show no clear signs of a cloud. Ozone mixing ratios are between 70 ppbv and 100 ppbv for altitudes above 16.9 km, \( i.e. \), from well below the cold-point.

The profile shown in Figure 8, with very tenuous subvisible cirrus clouds near the tropopause and a relatively high cold-point potential temperature, is similar to all those observed on 19 February. In the absence of strong convection near the flight path, the observed cooling and cloud condensation must be due to processes other than convection, such as gravity or planetary waves. Whilst convection reaching the tropopause was not frequent during the APE-THESO period, less energetic convection, capable of generating gravity waves, was widespread [Stefanutti et al., 2004]. A mechanism for the formation of the
profile shown in Figure 8 is discussed in Santacesaria et al. [2003]; a related example, but with
direct injection of water by convection into a cloud originally formed above convection, is discussed in
Garrett et al. [2004]. Trajectory calculations place the air observed on February 19 over convection less than
24 hours previously. Gravity-wave-induced turbulence above the convection probably initiated the transport
of water to the cold point. On subsequent cooling, this transported water vapour condensed to form a
subvisible cirrus. The UTTCs reported by Peter et al. [2003] and Luo et al. [2003a,b] (e.g. profile number 12
in Table 1) occur in profiles similar to that shown in Figure 8 and might originate from these subvisible cirrus
clouds. Ozone mixing ratios are close to 100 ppbv (Figure 8) which is a typical TTL value and thus
significantly higher compared to the profile in Figure 5 discussed above. This makes a recent direct injection
of lower tropospheric air unlikely.

The corresponding frequency distributions of RHi on 19 February (Figure 7, dashed lines) are different
than for 11 March. Outside clouds, much higher humidity is observed, including supersaturated air up to
130%. Inside the cirrus, also a higher fraction of air is supersaturated up to RHi = 170%, i.e. further from
equilibrium. This could suggest that a smaller fraction of their growth time-scale has elapsed since nucleation
compared to the fraction elapsed in the thicker clouds. Since the thick, anvil outflow, cirrus have number
densities about 1 order of magnitude larger than the thin cirrus [Thomas et al., 2002], the growth time-scale
of the thin clouds will be 1 order of magnitude longer for the thin clouds (i.e., 100 s compared to 10 s; see
e.g., Kärcher and Lohmann, 2002), observations of the thin clouds far from equilibrium will be more
frequent. Further, the coldest temperatures on 11 March were 185-190 K and thus the cloud among the
coldest ever observed. Under these conditions, supersaturation can be maintained even for hours. The cloud
might then have gone through several warming and cooling cycles due to changing temperatures, which
change the relative humidity significantly, but not necessarily the absolute amount of water vapor. The
findings support the idea that these thin clouds close to the tropopause are unlikely to originate from the
direct outflow of convection.
Both types of clouds affect the water budget of their environment in a different way: the thick cloud — in Figures 4, 5, and 7 — carries additional water into the TTL and is thus likely moistening the ambient air. However, the thin cloud close to the tropopause, in Figures 8 and 7, which is formed \textit{in situ} by local cooling in air masses close to saturation, has a greater chance of dehydrating the air to the local saturation water vapour mixing ratio, as discussed in \textit{Luo et al.} [2003b].

**Profile # 19: 6 March 1999**

Figure 9 shows a profile from 6 March 1999 (profile number 19 in Table 1). The hydropause is not well pronounced: there is a broad region (16.5 to 17.6 km) with mixing ratios below 4 ppmv, which is a much higher value than for the previously discussed profiles; the actual total-water minimum is at 16.8 km. The cold-point tropopause is coincident with the hydropause (16.8 km, $\theta = 377$ K) and occurs at the top of a moderately stable layer (lapse rate approximately 6 K km$^{-1}$), extending from a shallow inversion at 13.1 km. There is a marked increase in stability above the cold point, but a 500 m thick neutral layer occurs between 17.1 and 17.6 km. In the TTL, the water vapour mixing ratio is between 4 and 10 ppmv, the ozone mixing ratio is close to 50 ppbv, and the backscatter ratio near unity, i.e., no clouds (including UTTCs) are detected. Convective clouds with cloud-top temperatures above – 65 °C were present below the aircraft during this part of the flight and there are indications of small-scale waves that may have been induced by the convection, as discussed in \textit{Stefanutti et al.} [2004], especially their figures 7 and 11.

There are two regions of water vapour saturation in the profile shown in Figure 9: between 15.7 km and the cold point, and between 17.3 and 17.6 km (although there are a range of temperatures, and hence saturation mixing ratios for the portion of the flight near 17.5 km, the time series of data (not shown) demonstrate that there are periods of saturation). The coincidence of cold point, hydropause, and saturation indicates that the tropopause-level water vapour mixing ratio in this profile could be correctly deduced from a temperature-only sounding (e.g. a radiosonde), but the lack of cloud in the profile indicates that no dehydration is occurring at this place and time. However, since the profile is saturated near the tropopause,
the water vapour mixing ratio might have been "set" by the local temperatures just recently before the observation.

Profile # 21: 6 March 1999

Figure 10 shows a profile from 6 March 1999 (profile number 21 in Table 1). The hygropause is not well pronounced: there is a broad region with mixing ratios below 4 ppmv which is a much higher value than for the first two example profiles, the actual minimum being at 18.2 km. The cold-point tropopause is at 16.6 km (θ = 382 K), and occurs at the base of a broad (1-km deep) layer of cold temperatures. Above this layer there are two further near-adiabatic layers producing secondary temperature minima, each minimum matching features in the total water profile, including the hygropause at 18.2 km. In the TTL, the water vapour mixing ratio is between 4 and 8 ppmv, the ozone mixing ratio is close to 100 ppbv, and the backscatter ratio near unity, i.e., no clouds (including UTTCs) are detected.

In the cloud-free profile shown in Figure 10, multiple maxima/minima can be identified in the temperature, water vapour, and ozone profiles, and there is no simple relation between cold-point, hygropause, and cloud-top. These multiple maxima/minima may be due to differential advection of air that has experienced different degrees of dehydration, or may be due to the production of mixed layers when gravity waves reach their critical levels [Teitelbaum et al., 1999]. This type of profile could also be the result of projecting aircraft data from a sloping vertical travel into the height-temperature and height–water vapour planes, in which case the water vapour and ozone layers are indicative of mesoscale horizontal variations in the tropopause and hygropause height. However, since we are concerned with the large-scale ascent of air into the stratosphere, it does not matter greatly whether we consider these features to be directly on top of each other or not. Similar profiles were observed on 4 March, 15 March, and at other times on 6 March. The frequent occurrence of such profiles, which have no simple relationship between the temperature and water vapour profiles, questions the usefulness of applying a simple dehydration hypothesis to all vertical temperature soundings through the tropical tropopause region.
General view of the TTL above the Indian Ocean in early 1999

Out of 36 profiles during APE-THESEO, 7 show cold point, hygropause, and cloud top within 200 m of each other. A further five profiles show cold point and hygropause within 200 m of each other, but no clouds at that altitude. Two further profiles show cold point and cloud top within 200 m of each other. The frequent coincidence of cold point and hygropause suggests that dehydration is occurring, or has just occurred, in situ at the cold point over the Indian Ocean. The observations of clouds at the cold point, as indicated by FISH measurements, MAS measurements (Table 1), and other measurements aboard the Geophysica and from the Falcon [Thomas et al., 2002; Santacesaria et al., 2003], support this view.

Comparison with radiosonde measurements across the equatorial Indian Ocean show that the Geophysica data are consistent with the radiosonde data, and that the high, cold tropopause seen during APE-THESEO is representative of the entire region (Table 2). The tropopause temperatures observed by the Geophysica correspond to a mean saturation water vapour pressure of 3.0 ppmv. However, the mean measured hygropause mixing ratio was even lower, at 2.4 ppmv. This discrepancy with the mean saturation mixing ratio is due to the large number of unsaturated profiles. This clearly demonstrates that discussions of tropical troposphere-stratosphere dehydration and entry-level water should not be based solely on instantaneous temperature analyses [Vömel et al., 1999].

As expected for the season of APE-THESEO, the mean H₂O mixing ratios at the tropopause are below the averaged ‘entry level’ of stratospheric water vapour mixing ratios estimated from stratospheric measurements of water vapour and methane [e.g. Dessler and Kim, 1999, and references therein]. Figure 11 compares the range of APE-THESEO total water profiles with those from previous missions. Diagrams of this kind have been presented as evidence of the operation of the tropical ‘tape recorder’ — i.e., the upward transport of relatively isolated tropical air by the Brewer-Dobson circulation [Mote et al., 1996]. In the profile from Darwin in January-February 1987 [Kelly et al., 1993], the mean hygropause, shown in Figure 10, coincides with the mean tropopause. In the profile from Panama in September 1980 [Kley et al., 1982], the hygropause occurs at potential temperatures about 100 K above the tropopause. The data from APE-
THESEO show a broad minimum, at potential temperatures about 20 K higher than the Darwin minimum, but, like the Darwin minimum, coincident with the mean tropopause. Since we have presented evidence above for the occurrence of dehydration at potential temperatures above 370 K, passive (non-dehydrating) vertical advection by the ‘tape recorder’ is not the only mechanism determining the water vapour profile, for potential temperatures between 370 K and 400 K (i.e., above 17.5 km). The Darwin and APE-THESEO profiles have a similar shape between 390 K and 470 K, although the data from APE-THESEO are lower than the Darwin data. However, Vömel et al. [1995] also report for the same season — but for the Central Pacific Ocean — single profiles with water vapour mixing ratios at the hygropause of less than 1.5 ppmv.

As a whole, the results show the western equatorial Indian Ocean to be a site of active dehydration during northern winter/spring 1999. Comparing our results to the climatology of Newell and Gould-Stewart [1981], a number of possibilities arise. In this climatology, localised regions with high frequencies of monthly-mean 100-mbar temperatures of −82.4 °C or below were found. In Northern Hemisphere summer, these localised regions were the Indian and SE Asian monsoon regions. In Northern Hemisphere winter, the localised region was the western Pacific Ocean, with “an extension of this low temperature area into the Indian Ocean and over to Africa albeit with lower frequency” [Newell and Gould-Stewart, 1981]. APE-THESEO may have been fortunate to sample such a less frequent active dehydration period in the western Indian Ocean (cf. Figure 2 of Newell and Gould-Stewart [1981]). Bonazzola and Haynes [2004] and Fueglistaler and Haynes [2005] show, based on ECMWF operational analysis data and ERA-40 data respectively, that the APE-THESEO period was indeed characterised by lower tropopause temperatures over the western Pacific Ocean and the Indian Ocean region than, for example, the period December 1997 to February 1998. The mean entry-level water vapour mixing ratio across the tropics reported by Fueglistaler and Haynes [2005] is 2.2 ppmv, which agrees well with the mean hygropause total water mixing ratios of 2.4 ppmv as derived from our measurements, while those of the previous year are higher by approximately 1 ppmv [Bonazzola and Haynes, 2004]. Rather than APE-THESEO having sampled an anomalously cold year, another explanation could be that the Newell and Gould-Stewart climatology did not consider cold points at
pressures below 100 hPa, and this may have significantly affected their analysis, since the higher mode in the distribution of measured cold-point potential temperatures occurs at a pressure of about 80 mbar. Or, lastly, the apparent cooling of the tropopause and stratosphere since 1979 [e.g., Simmons et al., 1999] may have caused the area of the ‘stratospheric fountain’ to spread.

**Summary and Conclusions**

APE-THESEO made *in-situ* measurements of ozone, water vapour, total water, and cloud properties above the western equatorial Indian Ocean in February and March 1999. We have combined measurements of temperature, total water, water vapour, clear-sky relative humidity, and cloud to show evidence of active dehydration as air is transported from troposphere to stratosphere. The tropopause, as indicated by the temperature minimum or ‘cold point’, was high (i.e., a mean pressure-altitude of 17 km, equivalent to a potential temperature of 380 K) and cold (i.e., 190 K (−83 °C) in the mean). The mean measured water vapour minimum (i.e., the hygropause) was 2.4 ppmv, at a mean altitude close to the mean cold-point altitude. The mean saturation water vapour mixing ratio does not accurately represent the mean of the measured water vapour mixing ratios, since the air was unsaturated in the region of the cold point for about 40 % of the measurements. The very low mixing ratios observed during APE-THESEO are comparable to those found in previous studies in the ‘fountain region’ over Micronesia.

The lapse rate tropopause, although difficult to determine from aircraft data, was generally lower than the cold point, rather than at the same altitude as would be the case for radiative-convective adjustment [Thuburn and Craig, 2002]. The highest altitude of 100-ppbv ozone mixing ratios was also generally lower than the cold point. Clouds were observed up to the altitude of the cold-point. We identified cirrus clouds with differing characteristics: first, visible or just sub-visible cirrus clouds at relatively low potential temperatures, which stem from the direct outflow of anvils (with subsequent advection and lofting), and second, sub-visible cirrus and ultra-thin cirrus occurring at higher potential temperatures. These latter clouds are triggered either by vertical motion of air above convective systems or wave activity in the TTL, but are
not directly linked to the outflow of Cumulonimbus clouds, which do not often reach these altitudes over the Indian Ocean. The different origins of the two classes of clouds are indicated by the different frequency distributions of relative humidity — both inside and outside of the clouds — for each class. While the first class of clouds injected from the outflow of anvils into the dry TTL will moisten its environment, the second class of cirrus has the potential for effective dehydration of the air masses.

In single events, ongoing dehydration was observed in air parcels at potential temperatures as high as 390 K. Ongoing dehydration at potential temperatures this high will smear out the zonal-mean ‘tape recorder’ signal. Water vapour profiles from different regions of active dehydration, when plotted against potential temperature, will have minima at different potential temperatures — between 370 K and 400 K — and only above 400 K will passive vertical advection of the water vapour minimum by the Brewer-Dobson circulation dominate everywhere. Detailed transport studies are required to decide whether air passes through multiple active (i.e., dehydrating) cold points between 370 K and 400 K. Initial studies [Bonazzola and Haynes, 2004] suggest that transport through the layer between 360 K and 380 K is relatively localised, but that horizontal transport in the layer does bring air parcels through the coldest regions. In summary, the data from APE-THESEO demonstrate the complicated behaviour of water vapour, clouds, and ozone in the region of the tropical tropopause. The TTL over the western equatorial Indian Ocean was found to be a region with active dehydration down to very low water vapour mixing ratios, acting on air masses before they reach the stratosphere.

**Acknowledgements**

The authors gratefully acknowledge the help and advice of their APE-THESEO colleagues, including S. Bormann, R. Carla’, K. S. Carslaw, D. Lowe, B. P. Luo, P. Mazzinghi, V. Mitev, O. Riediger, G. Toci and M. Volk. We would like to thank the pilots and ground crew of the M55 Geophysica for the flexible and safe operation of the aircraft in a very difficult environment. This work was carried out as part of EC contract ENV4 CT97 0533, NERC contract GST/02/2210, NERC contract NER/T/S/2000/00977, BMBF contract 01 LA 9829/3 within the programme “Angewandte Klima- und Atmosphärenforschung”, with funding from the
Italian Space Agency (ASI), and with help-in-kind from the UK Meteorological Office. The authors are grateful to the Seychelles Directorate of Civil Aviation for substantial help with mission logistics. Paul Berrisford of the UK Universities’ Global Atmospheric Modelling Programme, and the British Atmospheric Data Centre, provided ERA climatologies, for which we thank them.

References


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Table 1: Individual profiles measured by the Geophysica during APE-THESO. All the profiles were made in the longitude sector 48°E - 60°E; more details, including plan views of flight paths, are available in Stefanutti et al. [2004].

<table>
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<th>Day</th>
<th>UCSE</th>
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UCSE = Unit for Connection of Scientific Equipment, which records aircraft parameters; Time = duration of flight segment, kilo-seconds since 00Z; Lat = latitude, °N; T<sub>cp</sub> = cold-point temperature, K; Z<sub>cp</sub> = cold-point pressure-altitude, km; θ<sub>cp</sub> = cold-point potential temperature, K; e<sub>cp</sub> = cold-point saturation mixing ratio, ppmv; χ<sub>H</sub> = hygropause water vapour mixing ratio, ppmv; Z<sub>H</sub> = hygropause pressure-altitude,
km; $e_H$ = hygropause saturation mixing ratio, ppmv; $Z_{CT}$ = cloud-top pressure-altitude for backscatter ratio > 2 at 532 nm; x = not reached; - = no data; nvc = no cloud, with backscatter ratio above 2 at 532 nm, observed.

**Table 2: Monthly mean cold-points and cold-point potential temperatures during 1999, for stations across the Indian Ocean**

<table>
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<tr>
<th>Station</th>
<th>Location</th>
<th>Cold-point temperature, K</th>
<th>Potential temperature of cold-point, K**</th>
<th>Number of ascents to pressures &lt; 100 mbar</th>
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<td>Kenya</td>
<td>17° S, 39° 23’ E</td>
<td>191 185 188</td>
<td>368 382 382 -</td>
<td>33 35 7</td>
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<td>Serge-Frolow</td>
<td>15° 20’S, 54° 20’E</td>
<td>191 193 190 190</td>
<td>388 372 398 382 22 27 27 29</td>
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<td>Seychelles</td>
<td>4° 42’S, 55° 30’E</td>
<td>189 189 187 189</td>
<td>392 388 382 388 28 59 54 56</td>
<td></td>
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<tr>
<td>Mauritius*</td>
<td>20° 11’S, 57° 20 E</td>
<td>198 197 192 196</td>
<td>382 378 378 378 17 16 17 20</td>
<td></td>
</tr>
<tr>
<td>E. Indian Ocean</td>
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<tr>
<td>Cocos Is.</td>
<td>12° S, 97° E</td>
<td>189 189 187 188</td>
<td>378 382 382 362 38 36 41 38</td>
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<tr>
<td>Sumatra</td>
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<td>186 187 186 191</td>
<td>398 392 392 392 11 16 41 54</td>
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*No ascents above 390 K are recorded in the database.

**Cold-point potential temperatures are for the mid-point of 5-K potential temperature bins.
**Figure 1**

Cold-point tropopause (crosses), lapse-rate tropopause (squares) and the highest pressure-altitudes where 100 and 200 ppbv ozone thresholds were found (diamonds), for all tropical ascents and descents of the APE-THESAO campaign. Profile numbers are as in table 1.

**Figure 2**

Potential temperature of the cold-point tropopause (crosses) and the hygropause (i.e. the minimum measured total water by FISH - diamonds) for each ascent and descent of the Geophysica across the hygropause. Additional minima in total water are shown as diamonds. Also shown is the section of each profile that is saturated with respect to water ice (grey bars).

**Figure 3**

Top panel: total water mixing ratio (diamonds) and saturated vapour mixing ratio over ice (squares) at the hygropause, for each ascent and descent of the Geophysica across the hygropause. The profiles are the same as those listed in Table 1.

Bottom panel: relative humidity (RH) over water ice at the hygropause, for each ascent and descent of the Geophysica across the hygropause. The region above 100 % RH is shaded.

**Figure 4**

Backscatter ratio (ratio of air plus aerosol to air molecular backscatter intensity) at 1064 nm of the cirrus cloud on 11 March 1999, measured by the lidar onboard the Falcon aircraft approximately 1 h prior to the Geophysica observations; Geophysica flight track is given by the blue line. The Falcon turned back at 4:30 in 14 S; hence the symmetry in the cirrus about the turn point. The aspect ratio is 1:250, i.e. the cloud structures are compressed horizontally by a factor of 250. The vertical profile shown in Figure 5 is marked by the arrow.
Figure 5

Vertical profiles from an ascent near 14 °S on 11 March 1999. Left panel: MAS backscatter ratio at 532 nm, FISH total water (ppmv), saturation vapour mixing ratio (ppmv). Note that MAS data extend off-scale. The maximum backscatter recorded by MAS on flight leg was 48, at an altitude of 16.6 km. Right panel: temperature (K), potential temperature (K), and FOZAN ozone mixing ratio (ppbv). The potential temperature data have been divided by two for ease of plotting.

Figure 6

Meteosat 5 infra-red cloud image, 5 hours (0116 LT, 2116 GMT) prior to the take-off of the Geophysica aircraft from the Seychelles (green islands in the centre of the plot) on 11 March 1999. Flight path is given as the black line. The aircraft first flew south-southeastward, then returned along the track north-northwestwards, overflew Mahé, reached the equator, and then returned south-southeastwards. Colours indicate cloud-top height in degrees Celsius (see legend bar). Orange arrows show the transport of recently-convected air at 100 hPa in the time intervals shown.

Figure 7

Frequency distribution of relative humidity as measured by FLASH on the flight on 11 March 1999 (solid lines) and on 19 February (dashed lines), separated for observations inside (red) and outside (blue) clouds.

Figure 8

Vertical profiles from a descent near the equator on 19 February 1999. Left panel: MAS backscatter ratio at 532 nm, FISH total water (ppmv), saturation vapour mixing ratio (ppmv). Right panel: temperature (K), potential temperature (K), and ECOC ozone mixing ratio (ppbv). The potential temperature data have been divided by two for ease of plotting.
**Figure 9**

Vertical profiles from an ascent near 11 °S on 6 March 1999. Left panel: MAS backscatter ratio at 532 nm, FISH total water (ppmv), saturation vapour mixing ratio (ppmv). Right panel: temperature (K), potential temperature (K), and ECOC ozone mixing ratio (ppbv). The potential temperature data have been divided by two for ease of plotting.

**Figure 10**

Vertical profiles from an ascent near 18 °S on 6 March 1999. Left panel: MAS backscatter ratio at 532 nm, FISH total water (ppmv), saturation vapour mixing ratio (ppmv). Right panel: temperature (K), potential temperature (K), and ECOC ozone mixing ratio (ppbv). The potential temperature data have been divided by two for ease of plotting.

**Figure 11**

A comparison of the range of water vapour profile from FISH (February-March 1999) with profiles from Darwin [Kelly et al., 1993], which took place in January and February 1987, and Panama [Kley et al., 1982], which took place in September 1980. The range bars on the FISH data show ±1 standard deviation.